

STABILITY OF SPHERICALLY SYMMETRIC WAVE MAPS.

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ABSTRACT. We study Wave Maps from \mathbf{R}^{2+1} to the hyperbolic plane \mathbf{H}^2 with smooth compactly supported initial data which are close to smooth spherically symmetric initial data with respect to some $H^{1+\mu}$, $\mu > 0$. We show that such Wave Maps don't develop singularities in finite time and stay close to the Wave Map extending the spherically symmetric data (whose existence is ensured by a theorem of Christodoulou-Tahvildar-Zadeh) with respect to all $H^{1+\delta}$, $\delta < \mu_0$ for suitable $\mu_0(\mu) > 0$. We obtain a similar result for Wave Maps whose initial data are close to geodesic ones. This strengthens a theorem of Sideris for this context.

1. INTRODUCTION

Let $\{(\mathbf{x}, \mathbf{y}) | \mathbf{y} > 0\}$ be the hyperbolic plane, equipped with metric $dg = \frac{d\mathbf{x}^2 + d\mathbf{y}^2}{\mathbf{y}^2}$, and let \mathbf{R}^{n+1} , $n \geq 1$ denote the standard Minkowski space equipped with the metric $dh = -dx_0^2 + \sum_{i=1}^n dx_i^2$. We shall also use the identifications $x_0 = t$ (time), $\partial_{x_\nu} = \partial_\nu$, $\nu = 0, 1, \dots, n$. A Wave Map from Minkowski space to \mathbf{H}^2 is a map $u : \mathbf{R}^{n+1} \rightarrow \mathbf{H}^2$ which is critical with respect to the functional

$$u \rightarrow \int_{\mathbf{R}^{n+1}} \langle \partial_\alpha u, \partial^\alpha u \rangle_g dx_0 dx_1 \dots dx_n,$$

where $\partial_\alpha u = u_*(\partial_\alpha)$, $\partial^\alpha = h^{\alpha\beta} \partial_\beta$ and Einstein's summation convention is in force. The Euler Lagrange equations of this problem read as follows:

$$\square \ln \mathbf{y} = -\frac{\partial_\nu \mathbf{x}}{\mathbf{y}} \frac{\partial^\nu \mathbf{x}}{\mathbf{y}} \quad (1)$$

$$\square \left(\frac{\mathbf{x}}{\mathbf{y}} \right) = \frac{\mathbf{x}}{\mathbf{y}} \frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2} \quad (2)$$

If $n = 2$, the fundamental Conjecture associated with this problem is the following, which flows from the intuition that the negative curvature of the target should prevent a focusing of energy in small spatial regions:

Conjecture(e. g. Klainerman [10]) Let $n = 2$. Given smooth initial data (\mathbf{x}, \mathbf{y}) , $(\partial_t \mathbf{x}, \partial_t \mathbf{y}) : \mathbf{R}^2 \times \{0\} \rightarrow \mathbf{H}^2 \times T\mathbf{H}^2$, there exists a global-in-time smooth Wave Map extending them.

This is expected to be generalizable to arbitrary targets of negative curvature and satisfying some geometric niceness conditions. We stick in this paper to the \mathbf{H}^2 model on account of its simplicity.

The difficulty in establishing the above conjecture stems from the fact that the

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problem is energy critical, i. e. the natural scale invariant Sobolev space is exactly the energy space \dot{H}^1 (the energy $\sum_{\alpha=0}^n \|\frac{\partial_\alpha \mathbf{x}}{\mathbf{y}}\|_{L_x^2}^2 + \|\frac{\partial_\alpha \mathbf{y}}{\mathbf{y}}\|_{L_x^2}^2$ is preserved under the Wave Map flow). Establishing global regularity for such problems consists customarily of showing that smooth small data imply global regularity, as well as non-concentration of energy in physical space. The latter needs to depend subtly on the geometry of the target, since a priori analytic reasons cannot rule out a rapid shift of the energy from low to high frequency modes, resulting in sudden focusing. In the case $n = 3$, one expects breakdown of solutions for large data for analytic reasons (the scale invariant Sobolev space $\dot{H}^{\frac{3}{2}}$ which in some sense controls the local well-posedness behavior is not controlled by the energy). We can formulate

Conjecture: *Let $n = 3$. There exist (large) smooth initial data $(\mathbf{x}, \mathbf{y}), (\partial_t \mathbf{x}, \partial_t \mathbf{y}) : \mathbf{R}^3 \times \{0\} \rightarrow \mathbf{H}^2 \times T\mathbf{H}^2$, which lead to breakdown in finite time.*

Breakdown solutions are known in $3 + 1$ dimensions, but only for special targets [3] not including the hyperbolic plane.

The best result known at this point pertaining to the first Conjecture is the following theorem of the author [22], [23]:

Theorem 1.1. *Let $n = 2, 3, \dots$. Then there exists $\epsilon > 0$ such that for smooth initial data $(\mathbf{x}, \mathbf{y}), (\partial_t \mathbf{x}, \partial_t \mathbf{y}) : \mathbf{R}^n \times \{0\} \rightarrow \mathbf{H}^2 \times T\mathbf{H}^2$ satisfying*

$$\int_{\mathbf{R}^2} \sum_{\nu=0}^n \left\| \frac{\partial_\nu \mathbf{x}}{\mathbf{y}} \right\|_{\dot{H}^{\frac{n}{2}}} + \left\| \frac{\partial_\nu \mathbf{y}}{\mathbf{y}} \right\|_{\dot{H}^{\frac{n}{2}}} < \epsilon,$$

there exists a smooth global-in-time Wave Map extending them.

This is similar to earlier results of Tao [37] when the target is a sphere, as well as of Tataru [41] when the target is uniformly isometrically embeddable into a Euclidean space. Similar results in dimensions $n \geq 4$ for quite general targets were achieved by Klainerman-Rodnianski [15], Shatah-Struwe [27] as well as Nahmod-Stefanov-Uhlenbeck [24] after Tao's initial breakthrough [36], [37].

Thus the preceding theorem does not yet exhibit behavior reflecting the geometric nature¹ of \mathbf{H}^2 .

What we set out to do in this paper is to try to exploit features which appear to hinge on geometric properties of this target and set it apart from positively curved targets such as the sphere S^2 . We need the following definition:

Definition 1.2. : We call a Wave Map $u : \mathbf{R}^{n+1} \rightarrow M$ 'spherically symmetric' provided $u(t, \rho x) = u(t, x) \forall x \in \mathbf{R}^2$ and $\rho : S^1 \rightarrow SO(2)$ the standard representation of S^1 on \mathbf{R}^2 .

We shall use the deep results of Christodoulou-Tahvildar-Zadeh [5] on the asymptotic behavior of spherically symmetric Wave Maps, valid for certain targets which amongst other things have no conjugate points, to conclude the following:

¹Paradoxically, the proof of this result involved extra complications over the case of target a sphere, on account of the fact that one needs to work with the derivative of the Wave Map, losing one degree of smoothness.

Theorem 1.3. (*Stability*) Let $n = 2$ and $u[0] = (u(0), \partial_t u(0)) : \mathbf{R}^2 \times \{0\} \rightarrow \mathbf{H}^2 \times T\mathbf{H}^2$ be a smooth spherically symmetric Wave Map with compactly supported initial data. Then for any $\sigma > 0$ there exists a number $\epsilon > 0$ such that for all initial data $\tilde{u}[0]$ which are ϵ -close to $u[0]$ in $H^{1+\sigma}$, there exists a smooth global in time Wave Map \tilde{u} extending $\tilde{u}[0]$. Moreover, \tilde{u} will stay close to u in the energy topology² (in a suitable sense) globally in time.

This is a type of 'large data result', although of course there is still a smallness assumption present. As far as the case of target S^2 is concerned, a surprising result of M. Struwe [34] asserts that smooth radial data lead to global Wave Maps. This suggests the important question of whether these solutions are unstable:

Question (*Instability?*). Let $u[0] : \mathbf{R}^2 \times \{0\} \rightarrow S^2 \times TS^2$ be large generic spherically symmetric initial data. Is it true that for any $\sigma > 0$, $\epsilon > 0$, there exist smooth initial data $\tilde{u}[0] : \mathbf{R}^2 \times \{0\} \rightarrow S^2 \times TS^2$ with the property that $\|u[0] - \tilde{u}[0]\|_{H^{1+\sigma}} < \epsilon$ while³ the smooth Wave Map \tilde{u} extending $\tilde{u}[0]$ locally in time breaks down after finite time? More precisely for any $\delta > 0$

$$\exists T < \infty \rightarrow \forall \delta > 0 \lim_{t \rightarrow T_-} \|\tilde{u}[t]\|_{H^{1+\delta} \times H^\delta} = \infty$$

Unfortunately, our techniques appear to have no bearing on this question. For example, we don't know what the asymptotic behavior of Struwe's solutions is.

The key ingredient to prove Theorem 1.3 is the boundedness of a range of subcritical Sobolev norms for large spherically symmetric Wave Maps:

Theorem 1.4. : There exists $\delta_0 > 0$ such that $\forall 0 \leq \delta < \delta_0$ and spherically symmetric smooth Wave Maps $u : \mathbf{R}^{2+1} \rightarrow \mathbf{H}^2$ we have

$$\sup_t \|u[t]\|_{H^{1+\delta} \times H^\delta} < \infty$$

The proof of this will follow from the asymptotic results of [5], which in turn rely on a careful analysis of conservation laws associated with (1), (2).

We shall then rely on the setup of [23], using the intrinsic derivative formulation (by differentiating (1), (2)) and passing to the Coulomb Gauge. The new difficulties by comparison with [23] concern nonlinear terms which are linear in the difference $\nabla[\tilde{u} - u]$. Working in the Coulomb Gauge, this corresponds to perturbing the flat d'Alembertian \square with a potential term V which is in some sense quadratic in the derivatives of the spherically symmetric Wave Map. We shall show that the good decay behavior of the spherically symmetric Wave Map allows us to treat these terms as source terms, instead of having to modify the linear operator. However, the fact that we cannot just work with mixed Lebesgue type norms but complicated null-frame spaces will force simultaneous localizations in physical and frequency space on us, which make the argument quite intricate. These types of estimates might be useful when working on the general large data problem. Our analysis shall have as simple corollary a generalization of a result of Sideris[29] to $n = 2$: we define a geodesic Wave Map $u(t, x)$ to be of the form $u = \gamma(v)$ where $\gamma(\cdot) : \mathbf{R} \rightarrow \mathbf{H}^2$ is a geodesic and $\square v = 0$. Then we have the following:

²Indeed, even in a certain range of subcritical spaces $H^{1+\lambda}$.

³To define this norm, use $S^2 \hookrightarrow \mathbf{R}^3$ and use standard coordinates

Theorem 1.5. *Let $u(t, x) : \mathbf{R}^{2+1} \rightarrow \mathbf{H}^2$ be a smooth geodesic Wave Map. Then there exists $\epsilon > 0$ such that for all initial data $\tilde{u}[0]$ ϵ -close to $u[0]$ in $H^{1+\sigma}$, there exists a global Wave Map extending $\tilde{u}[0]$. Also, \tilde{u} will stay close to u in the energy topology in a suitable sense.*

We also point out that due to a result of Shatah-Tahvildar-Zadeh [26] on the asymptotic behavior of equivariant Wave Maps, one expects a similar result for perturbations of large equivariant Wave Maps to hyperbolic targets.

2. A PRIORI ESTIMATES FOR SPHERICALLY SYMMETRIC WAVE MAPS.

For a Wave Map $u = (\mathbf{x}, \mathbf{y}) : \mathbf{R}^{2+1} \rightarrow \mathbf{H}^2$, we define the norm $\|u(t)\|_{H^s}$ as

$$\|(\mathbf{x}, \mathbf{y})\|_{H^s} := \sum_{\nu=0}^2 \left\| \frac{\partial_\nu \mathbf{x}}{\mathbf{y}}(t) \right\|_{H^{s-1}} + \left\| \frac{\partial_\nu \mathbf{y}}{\mathbf{y}} \right\|_{H^{s-1}}$$

We also introduce the following notation: $r = \sqrt{x_1^2 + x_2^2}$. Now let $u(t, x) = (\mathbf{x}, \mathbf{y})$ be a spherically symmetric Wave Map with compactly supported smooth initial data. Then we have

Lemma 2.1. *The image of the Wave Map belongs to a bounded subset of \mathbf{H}^2 . More precisely, we have*

$$\|\ln \mathbf{y}\|_{L_t^\infty L_x^\infty} < \infty, \left\| \frac{\mathbf{x}}{\mathbf{y}} \right\|_{L_t^\infty L_x^\infty} < \infty$$

The bounds depend (at most) on the size of the support as well as some norm $\|u[0]\|_{H^{1+\delta}}$, $\delta > 0$.

Proof : We shall rely on the following Proposition in [5]:

Proposition 2.2. (Chr.-Tah) *Under the previous assumptions, the following inequalities hold:*

$$\text{'Good derivative': } \left| \frac{\partial_v \mathbf{x}}{\mathbf{y}} \right| + \left| \frac{\partial_v \mathbf{y}}{\mathbf{y}} \right| \lesssim (t+r)^{-\frac{3}{2}}, \partial_v = \partial_t + \partial_r$$

$$\text{'Bad derivative': } \left| \frac{\partial_u \mathbf{x}}{\mathbf{y}} \right| + \left| \frac{\partial_u \mathbf{y}}{\mathbf{y}} \right| \lesssim (|t-r|+1)^{-1} (t+r)^{-\frac{1}{2}}, \partial_u = \partial_t - \partial_r$$

By local well-posedness of (1), (2) in H^s , $s > 1$, there exists a time interval $[-T, T]$ with $T = T(\|u[0]\|_{H^s})$ on which we uniformly control $\|\ln \mathbf{y}(t)\|_{H^s}$, $t \in [-T, T]$. Given an arbitrary point $(t, x) \in \mathbf{R}^{2+1}$ at distance $< \frac{T}{\sqrt{2}}$ from the forward light cone (say), connect it to a point in the strip $[-T, T] \times \mathbf{R}^2$ by means of a null-geodesic γ given by $t - r = u = \text{const}$. We have

$$\left| \frac{\partial_v \mathbf{y}}{\mathbf{y}} \right| \lesssim (t+r)^{-\frac{3}{2}}$$

on γ , whence

$$\int_{\gamma \cap \{t \geq T\}} \left| \frac{\partial_v \mathbf{y}}{\mathbf{y}} \right| dt < \infty$$

Combining this with the embedding $H^s \subset L_x^\infty$ yields the claim for $\ln \mathbf{y}$ for such points (t, x) . Given a point (t, x) at distance $\geq \frac{T}{\sqrt{2}}$ from the forward light cone,

connect it via a geodesic γ : $t + r = \text{const}$ to a point in the strip of thickness $\frac{T}{\sqrt{2}}$ around the light cone. Using

$$\left| \frac{\partial_u \mathbf{y}}{\mathbf{y}} \right| \lesssim u^{-1} v^{-\frac{1}{2}}$$

yields

$$\int_{\gamma} \left| \frac{\partial_u \mathbf{y}}{\mathbf{y}} \right| dt < \infty$$

Thus the claim follows in general for $\ln \mathbf{y}$. With this, one proceeds similarly for \mathbf{x} . \blacksquare

The following is the main result of this section:

Proposition 2.3. *Let the assumptions be as in the preceding lemma. There exists $\epsilon_0 > 0$ such that $\forall 0 < \epsilon < \epsilon_0$, there exists a global bound*

$$\sum_{\nu=0}^2 \|\partial_{\nu} \ln \mathbf{y}\|_{L_t^{\infty} H^{\epsilon}} < C_{\epsilon}, \quad \sum_{\nu=0}^2 \left\| \frac{\partial_{\nu} \mathbf{x}}{\mathbf{y}} \right\|_{L_t^{\infty} H^{\epsilon}} < C_{\epsilon},$$

where C_{ϵ} depends on the size of the support as well as the L^2 -mass of finitely many derivatives of $u[0]^4$

Proof : We shall need the following lemma, which is also due to Christodoulou-Tahvildar-Zadeh:

Lemma 2.4. *(Chr.-Tah) Introduce the vector fields $S = t\partial_t + r\partial_r$, $\Omega = r\partial_t + t\partial_r$. For a smooth spherically symmetric Wave Map $u(t, x) : \mathbf{R}^{2+1} \rightarrow \mathbf{H}^2$ and $\epsilon > 0$, we have*

$$r^{\frac{1}{2}+\epsilon} v^{1-\epsilon} [|\partial_v S \mathbf{x}| + |\partial_v S \mathbf{y}|] < \infty$$

$$r^{\frac{1}{2}+\epsilon} v^{1-\epsilon} [|\partial_v \Omega \mathbf{x}| + |\partial_v \Omega \mathbf{y}|] < \infty$$

As a consequence, we conclude that

$$|\nabla_{x,t} \partial_v \mathbf{x}| + |\nabla_{x,t} \partial_v \mathbf{y}| \lesssim v^{-(1-\epsilon)} u^{-1} r^{-\frac{1}{2}-\epsilon}$$

We shall in fact prove that the quantity

$$\begin{aligned} A(t) := \sum_{k \in \mathbf{Z}} 2^{\delta|k|} [& \|P_k(\ln \mathbf{y}(t, \cdot))\|_{\dot{H}_x^1} + \|P_k(\frac{\nabla_x \mathbf{x}}{\mathbf{y}}(t, \cdot))\|_{L_x^2} \\ & + \|P_k \partial_t(\ln \mathbf{y})(t, \cdot)\|_{L_x^2} + \|P_k(\frac{\partial_t \mathbf{x}}{\mathbf{y}})(t, \cdot)\|_{L_x^2}] \end{aligned}$$

is bounded globally in time, provided $\delta > 0$ is chosen sufficiently small. We have introduced here the Littlewood-Paley multipliers P_k , $k \in \mathbf{Z}$, which localize the (spatial) Fourier support to dyadic size $|\xi| \sim 2^k$. More precisely, let $\chi_0(\cdot) \in C_0^{\infty}(\mathbf{R}_{>0})$ have support contained in $(\frac{1}{2}, 2)$ and satisfy

$$\sum_{j \in \mathbf{Z}} \chi_0\left(\frac{x}{2^j}\right) = 1 \quad \forall x \in \mathbf{R}_{>0}.$$

⁴We are being imprecise here. All that matters to us is the global bound as stated.

Then we define (see also [31]) $P_k f$ for any $f \in L^1(\mathbf{R})$ via

$$\widehat{P_k f}(\xi) = \chi_0\left(\frac{|\xi|}{2^k}\right) \hat{f}(\xi)$$

We need to show that $A(T) \lesssim 1 + \int_0^T A(t) \phi(t) dt$ where $\phi(t)$ is integrable, in order to be able to apply Gronwall's inequality.

We first establish this for the contribution from $\ln \mathbf{y}$. Frequency-localize the equation (1), resulting in

$$\square P_k \ln \mathbf{y} = -P_k \left[\frac{\partial_\nu \mathbf{x}}{\mathbf{y}} \frac{\partial^\nu \mathbf{x}}{\mathbf{y}} \right] \quad (3)$$

We observe that $\frac{\partial_\nu \mathbf{x}}{\mathbf{y}} \frac{\partial^\nu \mathbf{y}}{\mathbf{y}} = \frac{(\partial_r - \partial_t) \mathbf{x}}{\mathbf{y}} \frac{(\partial_r + \partial_t) \mathbf{x}}{\mathbf{y}}$, and apply a Littlewood-Paley trichotomy:

$$\begin{aligned} \square P_k \ln \mathbf{y} = & -P_k \left[P_{<k-10} \frac{(\partial_r - \partial_t) \mathbf{x}}{\mathbf{y}} P_{[k-5, k+5]} \frac{(\partial_r + \partial_t) \mathbf{x}}{\mathbf{y}} \right] \\ & - P_k \left[P_{[k-10, k+10]} \frac{(\partial_r - \partial_t) \mathbf{x}}{\mathbf{y}} P_{<k+15} \frac{(\partial_r + \partial_t) \mathbf{x}}{\mathbf{y}} \right] \\ & - P_k \left[P_{>k+10} \frac{(\partial_r - \partial_t) \mathbf{x}}{\mathbf{y}} P_{>k+5} \frac{(\partial_r + \partial_t) \mathbf{x}}{\mathbf{y}} \right] \end{aligned} \quad (4)$$

We use the following terminology: $P_{<a} = \sum_{k < a} P_k$, $a, k \in \mathbf{Z}$, $P_{[a, b]} = \sum_{k \in [a, b]} P_k$ etc. We restrict ourselves to time interval $[0, T]$, and let $T \rightarrow \infty$. The general case will follow from time reversal symmetry. Applying Duhamel's formula, we see that we need to control the norm $\sum_{k \in \mathbf{Z}} 2^{\delta|k|} \|P_k(\cdot)\|_{L_t^1 L_x^2}$ of the right hand side by an expression $1 + \int_0^T A(t) \phi(t) dt$. We may restrict ourselves to a time interval $[c, \infty]$ for some $c > 0$ (depending on the initial data), on account of local-in-time well-posedness and finite propagation speed. We estimate each of the terms on the right-hand side of (4): *first assume $k \geq 0$.*

(i) *The first term:* We would like to place the 2nd input $P_{[k-5, k+5]} \frac{(\partial_r + \partial_t) \mathbf{x}}{\mathbf{y}}$ into $L_t^\infty L_x^2$ and the first input $P_{<k-10} \frac{(\partial_r - \partial_t) \mathbf{x}}{\mathbf{y}}$ into $L_t^1 L_x^\infty$. This doesn't integrate up, however. Placing the 2nd input into $L_t^\infty L_x^\infty$ will work provided t is much larger than $2^{|k|}$, but not in the opposite case: Thus we subdivide

$$\begin{aligned} & P_k \left[P_{<k-10} \frac{(\partial_r - \partial_t) \mathbf{x}}{\mathbf{y}} P_{[k-5, k+5]} \frac{(\partial_r + \partial_t) \mathbf{x}}{\mathbf{y}} \right] \\ &= \phi_{\geq 2^{\frac{k}{C}}}(t) P_k \left[P_{<k-10} \frac{(\partial_r - \partial_t) \mathbf{x}}{\mathbf{y}} P_{[k-5, k+5]} \frac{(\partial_r + \partial_t) \mathbf{x}}{\mathbf{y}} \right] \\ &+ \phi_{< 2^{\frac{k}{C}}}(t) P_k \left[P_{<k-10} \frac{(\partial_r - \partial_t) \mathbf{x}}{\mathbf{y}} P_{[k-5, k+5]} \frac{(\partial_r + \partial_t) \mathbf{x}}{\mathbf{y}} \right], \end{aligned}$$

where $\phi_{\geq a}(t)$, $\phi_{< a}(t)$ are smooth cutoffs to dilates of the regions $t \geq a$, $t < a$, adding up to 1. Also, C is a large number to be chosen. We can immediately

estimate

$$\begin{aligned}
& 2^{\delta k} \|\phi_{\geq 2^{\frac{k}{C}}}(t) P_k[P_{< k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}} P_{[k-5, k+5]} \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}}]\|_{L_t^1 L_x^2} \\
& \lesssim 2^{\delta k} \|P_{< k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}}\|_{L_t^\infty L_x^2} \|\phi_{\geq 2^{\frac{k}{C}}}(t) P_{[k-5, k+5]} \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}}\|_{L_t^1 L_x^\infty} \\
& \lesssim 2^{\delta k} 2^{-\frac{k}{2C}}
\end{aligned}$$

This can be summed over $k \geq 0$ provided $\delta < \frac{1}{2C}$.

Now we proceed to the case in which time is dominated by frequency, $t \lesssim 2^{\frac{k}{C}}$. We shall distinguish between the region separated from the light cone, where we use lemma 2.4, as well as the region very close to the light cone, where we use Proposition 2.2 as well as Hoelder's inequality: we decompose

$$\begin{aligned}
P_{[k-5, k+5]} \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}} &= P_{[k-5, k+5]}(\psi_{> \frac{t}{2}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}}) \\
&\quad + P_{[k-5, k+5]}(\psi_{\frac{t}{2} > \cdot \geq 2^{-\mu k}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}}) \\
&\quad + P_{[k-5, k+5]}(\psi_{< 2^{-\mu k}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})
\end{aligned}$$

where the smooth cutoffs $\psi_{> \frac{t}{2}}(\cdot)$, $\psi_{\frac{t}{2} > \cdot \geq 2^{-\mu k}}(\cdot)$, $\psi_{< 2^{-\mu k}}(\cdot)$ add up to 1 and localize, respectively, to dilates of the regions indicated in their subscripts. We let μ be a small positive number to be chosen. We have

$$\begin{aligned}
& 2^{\delta k} \|\phi_{< 2^{\frac{k}{C}}}(t) P_k[P_{< k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}} \\
& \quad P_{[k-5, k+5]}(\psi_{> \frac{t}{2}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})]\|_{L_t^1 L_x^2([c, T] \times \mathbf{R}^2)} \\
& \lesssim 2^{(\delta-1)k} \int_c^T \|P_{< k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}}(t)\|_{L_x^\infty} \\
& \quad \|P_{[k-5, k+5]} \nabla_x(\psi_{> \frac{t}{2}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})\|_{L_x^2} dt
\end{aligned}$$

Now using lemma 2.4 as well as Proposition 2.2

$$\|P_{[k-5, k+5]} \nabla_x(\psi_{> \frac{t}{2}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})\|_{L_x^2} \lesssim t^{-(2-\epsilon)} \sqrt{\int_0^{\frac{t}{2}} r^{-1-2\epsilon} r dr} \lesssim t^{-\frac{3}{2}}.$$

This implies, reiterating application of Proposition 2.2

$$\begin{aligned}
& 2^{\delta k} \|\phi_{< 2^{\frac{k}{C}}}(t) P_k[P_{< k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}} \\
& \quad P_{[k-5, k+5]}(\psi_{> \frac{t}{2}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})]\|_{L_t^1 L_x^2([c, T] \times \mathbf{R}^2)} \\
& \lesssim 2^{(\delta-1)k} \int_c^T t^{-\frac{3}{2}} dt,
\end{aligned}$$

which can be summed over $k \geq 0$, provided $\delta < 1$. Next, we estimate

$$\begin{aligned}
& 2^{\delta k} \|\phi_{<2^{\frac{k}{C}}}(t) P_k [P_{<k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}} \\
& \quad P_{[k-5, k+5]}(\psi_{\frac{t}{2} > \cdot, \geq 2^{-\mu k}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})]\|_{L_t^1 L_x^2([c, T] \times \mathbf{R}^2)} \\
& \lesssim 2^{(\delta-1)k} \|P_{<k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}}\|_{L_t^\infty L_x^2} \\
& \quad \|P_{[k-5, k+5]} \nabla_x (\psi_{\frac{t}{2} > \cdot, \geq 2^{-\mu k}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})\|_{L_t^1 L_x^\infty([c, T] \times \mathbf{R}^2)} \\
& \lesssim 2^{(\delta-1+\mu)k} \|P_{<k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}}\|_{L_t^\infty L_x^2} \int_c^T t^{-\frac{3}{2}} dt.
\end{aligned}$$

We can sum here over k provided $\mu + \delta < 1$. Finally, we calculate using Hoelder's inequality as well as Proposition 2.2

$$\begin{aligned}
& 2^{\delta k} \|\phi_{<2^{\frac{k}{C}}}(t) P_k [P_{<k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}} \\
& \quad P_{[k-5, k+5]}(\psi_{<2^{-\mu k}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})]\|_{L_t^1 L_x^2([c, T] \times \mathbf{R}^2)} \\
& \lesssim \min\{T, 2^{\frac{k}{C}}\} 2^{\delta k} \|P_{<k-10} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}}\|_{L_x^\infty} \|P_{[k-5, k+5]}(\psi_{<2^{-\mu k}}(|t| - |x|) \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}})\|_{L_x^2} \\
& \lesssim 2^{(\delta + \frac{1}{C} - \frac{\mu}{2})k}
\end{aligned}$$

This can be summed over $k \geq 0$ provided we have $\delta + \frac{1}{C} < \frac{\mu}{2}$. Combining with the conditions obtained earlier, namely $\delta < \frac{1}{2C}$ as well as $\delta + \mu < 1$, we get $\delta < \frac{1}{7}$.

(ii) *The 2nd term of (4)*. This term appears immediate on account of Proposition 2.2. Formally

$$\begin{aligned}
& 2^{\delta k} \|P_k [P_{[k-10, k+10]} \frac{(\partial_r - \partial_t)\mathbf{x}}{\mathbf{y}} P_{<k+15} \frac{(\partial_r + \partial_t)\mathbf{x}}{\mathbf{y}}]\|_{L_t^1 L_x^2([c, T] \times \mathbf{R}^2)} \\
& \lesssim \int_c^T A(t) (1+t)^{-\frac{3}{2}} dt
\end{aligned}$$

We have to argue more carefully here since ∂_u involves $\partial_r = \frac{x_1}{r} \partial_{x_1} + \frac{x_2}{r} \partial_{x_2}$. Decompose

$$\begin{aligned}
P_k(\frac{\partial_r \mathbf{x}}{\mathbf{y}}) &= P_k[P_{<k-10}(\frac{x_1}{r}) P_{[k-10, k+10]}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}}) \\
&+ P_{[k-10, k+10]}(\frac{x_1}{r}) P_{<k+15}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}}) + P_{>k+10}(\frac{x_1}{r}) P_{>k+5}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}})] \\
&+ \text{similar terms}
\end{aligned} \tag{5}$$

The first term in this Littlewood-Paley trichotomy is estimated exactly as before, so we treat the 2nd and third term. Let $\chi_0(\cdot) \in C_0^\infty(\mathbf{R}_{>0})$ be the cutoff used for the Littlewood-Paley localizers P_k . We note that

$$P_k(\frac{x_1}{r})(x) = 2^{2k} \int_{\mathbf{R}^2} \widehat{\chi_0}(2^k(x-y)) \frac{y_1}{|y|} dy = 2^{2k} \int_{\mathbf{R}^2} \widehat{\chi_0}(2^k y) [\frac{y_1 - x_1}{|y - x|} - \frac{x_1}{|x|}] dy$$

On account of the inequality

$$\left| \frac{y_1 - x_1}{|y - x|} - \frac{x_1}{|x|} \right| \lesssim \min\left\{ \frac{|y|}{|x|}, 1 \right\},$$

we get, using the rapid decay of $y \rightarrow \widehat{\chi}_0(2^k y)$ outside of a disc of radius $\sim 2^{-k}$:

$$|P_k(\frac{x_1}{r})(x)| \lesssim \min\left\{ \frac{2^{-k}}{|x|}, 1 \right\}.$$

We introduce another cutoff $\chi_{2^{-\frac{k}{2}}}(x)$ which smoothly localizes to a dilate of the disc $B_{2^{-\frac{k}{2}}}(\mathbf{0})$ centered at $\mathbf{0} = (0, 0)$. We then decompose

$$\begin{aligned} & P_k[P_{[k-10, k+10]}(\frac{x_1}{r})P_{<k+15}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}})P_{<k+15}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})] = \\ & P_k[\chi_{2^{-\frac{k}{2}}}(x)P_{[k-10, k+10]}(\frac{x_1}{r})P_{<k+15}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}})P_{<k+15}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})] \\ & + P_k[(1 - \chi_{2^{-\frac{k}{2}}}(x))P_{[k-10, k+10]}(\frac{x_1}{r})P_{<k+15}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}})P_{<k+15}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})] \end{aligned}$$

Using Proposition 2.2 and Hoelder's inequality, we get

$$\begin{aligned} & 2^{\delta k} \|P_k[\chi_{2^{-\frac{k}{2}}}(x)P_{[k-10, k+10]}(\frac{x_1}{r})P_{<k+15}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}})P_{<k+15}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})]\|_{L_t^1 L_x^2} \\ & \lesssim 2^{\delta k} \|\chi_{2^{-\frac{k}{2}}}(x)P_{[k-10, k+10]}(\frac{x_1}{r})\|_{L_t^\infty L_x^2} \|P_{<k+15}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}})P_{<k+15}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})\|_{L_t^1 L_x^\infty} \\ & \lesssim 2^{(\delta - \frac{1}{2})k}. \end{aligned}$$

On the other hand, using the preceding calculations as well as Proposition 2.2 we get

$$\begin{aligned} & 2^{\delta k} \|P_k[(1 - \chi_{2^{-\frac{k}{2}}}(x))P_{[k-10, k+10]}(\frac{x_1}{r})P_{<k+15}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}})P_{<k+15}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})]\|_{L_t^1 L_x^2} \\ & \lesssim 2^{\delta k} \|(1 - \chi_{2^{-\frac{k}{2}}}(x))P_{[k-10, k+10]}(\frac{x_1}{r})\|_{L_t^\infty L_x^\infty} \|P_{<k+15}(\frac{\partial_1 \mathbf{x}}{\mathbf{y}})P_{<k+15}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})\|_{L_t^1 L_x^\infty} \\ & \lesssim 2^{(\delta - \frac{1}{2})k}. \end{aligned}$$

Since we have to choose $\delta < \frac{1}{7}$, both can be summed over $k \geq 0$. The case corresponding to the third term in (5) as well as the remaining terms are handled analogously.

(iii) *The third term of (4):* We can write

$$P_k[P_{>k+10}(\frac{\partial_u \mathbf{x}}{\mathbf{y}})P_{>k+5}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})] = \sum_{l_1 > k+10, |l_1 - l_2| < 5} P_k[P_{l_1}(\frac{\partial_u \mathbf{x}}{\mathbf{y}})P_{l_2}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})]$$

Next, we estimate, using Proposition 2.2

$$\begin{aligned} & \sum_{l_1 > k+10, |l_1-l_2| < 5} 2^{\delta k} \|P_k[P_{l_1}(\frac{\partial_u \mathbf{x}}{\mathbf{y}})P_{l_2}(\frac{\partial_v \mathbf{x}}{\mathbf{y}})]\|_{L_t^1 L_x^2([c,T] \times \mathbf{R}^2)} \\ & \lesssim \sum_{l_1 > k+10, |l_1-l_2| < 5} 2^{\delta(k-l_1)} \int_c^T [2^{\delta l_1} \|P_{l_1}(\frac{\partial_u \mathbf{x}}{\mathbf{y}})(t, \cdot)\|_{L_x^2}] t^{-\frac{3}{2}} dt, \end{aligned}$$

and one can sum here over both l_1, k to obtain the upper bound $\lesssim 1 + \int_c^T A(t) t^{-\frac{3}{2}} dt$. We have to argue for $\frac{\partial_u \mathbf{x}}{\mathbf{y}}$ as in (ii). This completes the estimates for case $k \geq 0$. For the case $k < 0$, we have for $M > 6$

$$\|P_k[\frac{\partial_u \mathbf{x}}{\mathbf{y}} \frac{\partial_v \mathbf{x}}{\mathbf{y}}]\|_{L_t^1 L_x^2} \lesssim 2^{\frac{2k}{M}} \|\frac{\partial_u \mathbf{x}}{\mathbf{y}}\|_{L_t^\infty L_x^2} \|\frac{\partial_v \mathbf{x}}{\mathbf{y}}\|_{L_t^1 L_x^M} \lesssim 2^{\frac{2k}{M}},$$

We have used *Bernstein's inequality* which states that for any rectangle $R \subset \mathbf{R}^2$ and smooth cutoff χ_R supported in R we have⁵

$$\|\mathcal{F}^{-1}(\chi_R \mathcal{F}f)\|_{L_x^q} \lesssim |R|^{\frac{1}{p}-\frac{1}{q}} \|f\|_{L_x^p}, \quad p \leq q.$$

Also, the estimate for $\|\frac{\partial_v \mathbf{x}}{\mathbf{y}}\|_{L_t^1 L_x^M}$ follows from interpolating between the decay estimate for $\|\frac{\partial_v \mathbf{x}}{\mathbf{y}}(t)\|_{L_x^\infty}$ and energy conservation.

The estimates for $\frac{\partial_v \mathbf{x}}{\mathbf{y}}$ are similar: we have by the same reasoning as before

$$\sum_{k \in \mathbf{Z}} 2^{\delta k} \|P_k[\frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2}]\|_{L_t^1 L_x^2([c,T] \times \mathbf{R}^2)} \lesssim 1 + \int_c^T A(t) t^{-\frac{3}{2}} dt$$

Now as for the nonlinearity on the right hand side of (2), the small frequency case $k < 0$ follows exactly as above from the boundedness of $\frac{\mathbf{x}}{\mathbf{y}}$, see lemma 2.1. As for the large frequency case, we have the usual frequency trichotomy

$$\begin{aligned} P_k[\frac{\mathbf{x}}{\mathbf{y}} \frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2}] &= P_k[P_{[k-5, k+5]}(\frac{\mathbf{x}}{\mathbf{y}}) P_{< k-10}[\frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2}]] \\ &\quad + P_k[P_{< k+15}(\frac{\mathbf{x}}{\mathbf{y}}) P_{[k-10, k+10]}[\frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2}]] \\ &\quad + P_k[P_{> k+5}(\frac{\mathbf{x}}{\mathbf{y}}) P_{> k+10}[\frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2}]] \end{aligned}$$

We need

Lemma 2.5. *The following inequality holds:*

$$\sum_{k \geq 0} 2^{\delta k} \|P_k \nabla_x(\frac{\mathbf{x}}{\mathbf{y}})(t)\|_{L_x^2} \lesssim A(t) + 1$$

Proof : Call the left hand side $B(t)$. Note that

$$B(t) \leq A(t) + \sum_{k \geq 0} 2^{\delta k} \|P_k(\frac{\mathbf{x}}{\mathbf{y}} \frac{\nabla \mathbf{y}}{\mathbf{y}})(t)\|_{L_x^2}$$

⁵We denote the spatial Fourier transform of $f(x)$ either by \hat{f} or $\mathcal{F}f$.

We have the frequency trichotomy

$$\begin{aligned} P_k\left[\frac{\nabla \mathbf{y}}{\mathbf{y}} \frac{\mathbf{x}}{\mathbf{y}}\right] &= P_k[P_{< k-10}\left(\frac{\nabla \mathbf{y}}{\mathbf{y}}\right)P_{[k-5, k+5]}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)] \\ &+ P_k[P_{[k-10, k+10]}\left(\frac{\nabla \mathbf{y}}{\mathbf{y}}\right)P_{< k+15}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)] + P_k[P_{> k+10}\left(\frac{\nabla \mathbf{y}}{\mathbf{y}}\right)P_{> k+5}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)] \end{aligned}$$

The estimate is immediate for the 2nd term on the right hand side. As to the first, we have

$$\begin{aligned} &2^{\delta k} \|P_k[P_{< k-10}\left(\frac{\nabla \mathbf{y}}{\mathbf{y}}\right)P_{[k-5, k+5]}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)](t)\|_{L_x^2} \\ &\lesssim 2^{(\delta-1)k} \|P_{< k-10}\left(\frac{\nabla \mathbf{y}}{\mathbf{y}}\right)\|_{L_x^\infty} \|\nabla_x P_{[k-5, k+5]}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)(t)\|_{L_x^2} \\ &\lesssim 2^{(\delta-1)k} B(t) \end{aligned}$$

One can also estimate this term by $\lesssim 2^{\delta k}$ from energy conservation and lemma 2.1. The estimate for the third term in the preceding trichotomy is similar. We conclude that

$$B(t) \lesssim A(t) + \sum_{0 \leq k \leq C} 2^{\delta k} + \sum_{k > C} 2^{(\delta-1)k} B(t)$$

Choosing C large enough, one obtains the claim of the lemma. \blacksquare

Armed with this, we now have (we may assume $k \geq 10$)

$$\begin{aligned} &2^{\delta k} \|P_k[P_{[k-5, k+5]}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)P_{< k-10}\left[\frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2}\right]]\|_{L_t^1 L_x^2([c, T] \times \mathbf{R}^2)} \\ &\lesssim 2^{(\delta-1)k} \int_c^T \|\nabla_x P_{[k-5, k+5]}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)(t)\|_{L_x^2} \|P_{< k-10}\left[\frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2}\right](t)\|_{L_x^\infty} dt \end{aligned}$$

Using Proposition 2.2 as well as the preceding lemma and summing over $k \geq 10$, we bound this by $\lesssim 1 + \int_c^T A(t) t^{-\frac{3}{2}} dt$. The estimate for the third term in the frequency trichotomy preceding the last lemma is more of the same. Thus we get

$$\sum_{k \in \mathbf{Z}} 2^{\delta|k|} \|P_k\left[\frac{\mathbf{x}}{\mathbf{y}} \frac{\partial_\nu \mathbf{y} \partial^\nu \mathbf{y} + \partial_\nu \mathbf{x} \partial^\nu \mathbf{x}}{\mathbf{y}^2}\right]\|_{L_t^1 L_x^2([c, T] \times \mathbf{R}^2)} \lesssim 1 + \int_c^T A(t) t^{-\frac{3}{2}} dt$$

Using Duhamel's formula, we get

$$\sum_{k \in \mathbf{Z}} 2^{\delta|k|} [\|P_k \nabla_x \left(\frac{\mathbf{x}}{\mathbf{y}}\right)(T)\|_{L_x^2} + \|P_k \partial_t \left(\frac{\mathbf{x}}{\mathbf{y}}\right)(T)\|_{L_x^2}] \lesssim 1 + \int_c^T A(t) t^{-\frac{3}{2}} dt$$

We need to estimate $\frac{\nabla \mathbf{x}}{\mathbf{y}}(T)$, which differs from the preceding by $\frac{\mathbf{x}}{\mathbf{y}} \frac{\nabla \mathbf{y}}{\mathbf{y}}(T)$. For frequencies ≥ 0 , this is estimated as in the preceding lemma, observing that we already improved the estimate for $\|\frac{\nabla \mathbf{y}}{\mathbf{y}}(T)\|_{L_x^2}$ from the preceding estimates (i)-(iii). The only case not yet covered concerns small frequencies. However, we have for $k < 0$

$$\begin{aligned} P_k\left[\frac{\mathbf{x}}{\mathbf{y}} \frac{\nabla \mathbf{y}}{\mathbf{y}}\right] &= P_k[P_{[k-5, k+5]}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)P_{< k-10}\left(\frac{\nabla \mathbf{y}}{\mathbf{y}}\right)] \\ &+ P_k[P_{< k+15}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)P_{[k-10, k+10]}\left(\frac{\nabla \mathbf{y}}{\mathbf{y}}\right)] + P_k[P_{> k+5}\left(\frac{\mathbf{x}}{\mathbf{y}}\right)P_{> k+10}\left(\frac{\nabla \mathbf{y}}{\mathbf{y}}\right)] \end{aligned}$$

Then

$$\begin{aligned} \sum_{k < 0} 2^{\delta|k|} \|P_k[P_{[k-5, k+5]}(\frac{\mathbf{x}}{\mathbf{y}})P_{<k-10}(\frac{\nabla \mathbf{y}}{\mathbf{y}})](T)\|_{L_x^2} \\ \lesssim \sum_{k < 0} 2^{\delta|k|} \|P_{[k-5, k+5]}(\frac{\mathbf{x}}{\mathbf{y}})\|_{L_x^\infty} \|P_{<k-10}(\frac{\nabla \mathbf{y}}{\mathbf{y}})(T)\|_{L_x^2}, \end{aligned}$$

which in turn is bounded by $\lesssim 1 + \int_c^T A(t)t^{-\frac{3}{2}}dt$, as is easily⁶ verified. The estimate for the 2nd term is immediate and the estimate for the third term as follows:

$$\begin{aligned} \sum_{k \in \mathbf{Z}_{<0}} 2^{\delta|k|} \|P_k[P_{>k+5}(\frac{\mathbf{x}}{\mathbf{y}})P_{>k+10}(\frac{\nabla \mathbf{y}}{\mathbf{y}})](T)\|_{L_x^2} \\ \lesssim \sum_{k \in \mathbf{Z}_{<0}} \sum_{l_1 > k+10, |l_1-l_2| < 5} \|P_k[P_{l_2}(\frac{\mathbf{x}}{\mathbf{y}})P_{l_1}(\frac{\nabla \mathbf{y}}{\mathbf{y}})](T)\|_{L_x^2} \\ \lesssim \sum_{k \in \mathbf{Z}_{<0}} \sum_{l_1 > k+10, |l_1-l_2| < 5} 2^{\delta|k|} 2^{k-l_1} \|P_{l_2} \nabla_x(\frac{\mathbf{x}}{\mathbf{y}})(T)\|_{L_x^2} \|P_{l_1}(\frac{\nabla \mathbf{y}}{\mathbf{y}})\|_{L_x^2} \end{aligned}$$

One verifies easily from the preceding estimates that this is $\lesssim 1 + \int_c^\infty A(t)t^{-\frac{3}{2}}dt$. Putting all of these ingredients together, we obtain

$$A(t) \lesssim 1 + \int_c^T A(t)t^{-\frac{3}{2}}dt.$$

The desired upper bound now follows from Gronwall's inequality. \blacksquare

Corollary 2.6. *Let $N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})$ denote any one of the nonlinearities occurring on the right hand side of (1), (2). Then for δ as in Proposition 2.3, we have for $\delta < \frac{1}{7}$*

$$\sum_{k \in \mathbf{Z}} 2^{\delta|k|} \|P_k N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})\|_{L_t^1 L_x^2([-T, T] \times \mathbf{R}^2)} < \infty$$

This follows from the preceding proof and time reversal symmetry. In the same vein, we have the following lemma:

Lemma 2.7. *Choosing $\delta > 0$ small enough, we have the inequality*

$$\sum_{k \in \mathbf{Z}} 2^{\delta|k|} \|P_k N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})\|_{L_t^2 \dot{H}^{-\frac{1}{2}}} < \infty$$

Proof : We work with $N(\dots) = \frac{\partial_\nu \mathbf{x}}{\mathbf{y}} \frac{\partial^\nu \mathbf{x}}{\mathbf{y}}$, the other cases being similar. Divide into the cases $k \geq 0$ and $k < 0$. In the first case, estimate

$$\begin{aligned} 2^{\delta k} \|P_k N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})\|_{L_t^2 \dot{H}^{-\frac{1}{2}}} \\ \lesssim 2^{(\delta - \frac{1}{2})k} \left\| \frac{(\partial_t - \partial_r) \mathbf{x}}{\mathbf{y}} \right\|_{L_t^\infty L_x^2} \left\| \frac{(\partial_t + \partial_r) \mathbf{x}}{\mathbf{y}} \right\|_{L_t^2 L_x^\infty} \lesssim 2^{(\delta - \frac{1}{2})k} \end{aligned}$$

In the 2nd case, estimate

$$\begin{aligned} 2^{-\delta k} \|P_k N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})\|_{L_t^2 \dot{H}^{-\frac{1}{2}}} \\ \lesssim 2^{(-\delta + \frac{2}{4-} - \frac{1}{2})k} \left\| \frac{(\partial_t - \partial_r) \mathbf{x}}{\mathbf{y}} \right\|_{L_t^\infty L_x^2} \left\| \frac{(\partial_t + \partial_r) \mathbf{x}}{\mathbf{y}} \right\|_{L_t^2 L_x^{4-}} \lesssim 2^{(-\delta + \frac{2}{4-} - \frac{1}{2})k} \end{aligned}$$

⁶Use Bernstein's inequality.

We have used here that

$$\| \frac{(\partial_t + \partial_r)\mathbf{x}}{\mathbf{y}}(t) \|_{L_x^{4-}} \lesssim t^{-\frac{3}{4+}}$$

which follows from interpolating between Proposition 2.2 and energy conservation. Choosing $\delta < \frac{2}{4-} - \frac{1}{2}$ results in the claim of the lemma. \blacksquare

3. THE PERTURBATION ARGUMENT

3.1. Precise statement of theorem. Outline of the procedure. The formulation (1), (2), while good enough for the purposes of the last section, will not suffice for us here⁷. Instead, following the procedure in [23], we shall pass to the derivative formulation of the problem, and translate everything into the Coulomb Gauge. More precisely, introduce the variables $\phi_\nu^1 = -\frac{\partial_\nu \mathbf{x}}{\mathbf{y}}$, $\phi_\nu^2 = -\frac{\partial_\nu \mathbf{y}}{\mathbf{y}}$, pass to complex notation $\phi_\nu = \phi_\nu^1 + i\phi_\nu^2$, and revert to the Coulomb Gauge by introducing the variables $\psi_\nu = \phi_\nu e^{-i\Delta^{-1} \sum_{i=1,2} \partial_i \phi_i^1}$. One gets the following remarkable self-contained divergence curl system:

$$\partial_\alpha \psi_\beta - \partial_\beta \psi_\alpha = i\psi_\beta \Delta^{-1} \sum_{j=1,2} \partial_j (\psi_\alpha^1 \psi_j^2 - \psi_\alpha^2 \psi_j^1) - i\psi_\alpha \Delta^{-1} \sum_{j=1,2} \partial_j (\psi_\beta^1 \psi_j^2 - \psi_\beta^2 \psi_j^1) \quad (6)$$

$$\partial_\nu \psi^\nu = i\psi^\nu \Delta^{-1} \sum_{j=1,2} \partial_j (\psi_\nu^1 \psi_j^2 - \psi_\nu^2 \psi_j^1). \quad (7)$$

From these one easily deduces the following system of wave equations:

$$\begin{aligned} \square \psi_\alpha &= i\partial^\beta [\psi_\alpha \Delta^{-1} \sum_{j=1}^2 \partial_j [\psi_\beta^1 \psi_j^2 - \psi_\beta^2 \psi_j^1]] \\ &\quad - i\partial^\beta [\psi_\beta \Delta^{-1} \sum_{j=1}^2 \partial_j [\psi_\alpha^1 \psi_j^2 - \psi_\alpha^2 \psi_j^1]] \\ &\quad + i\partial_\alpha [\psi_\nu \Delta^{-1} \sum_{j=1}^2 \partial_j [\psi^{1\nu} \psi_j^2 - \psi^{2\nu} \psi_j^1]]. \end{aligned} \quad (8)$$

As in [23], these in conjunction with the underlying first-order system (6), (7) shall form the basis for our estimates. We can now give the precise version of Theorem 1.3:

Theorem 3.2. *Let $u : \mathbf{R}^{2+1} \rightarrow \mathbf{H}^2$ be a smooth spherically symmetric Wave Map with compactly supported (large) initial data. Let $\{\psi_\nu\}_{\nu=0}^2$ be the derivative components in the Coulomb Gauge. Then for any $\mu > 0$ there exists $\epsilon = \epsilon(u, \mu) > 0$ such that for all smooth initial data $\tilde{u}[0] = (\tilde{u}(0), \partial_t \tilde{u}(0))$ with $\|(u - \tilde{u})[0]\|_{H^{1+\mu} \times H^\mu} < \epsilon$,*

⁷It appears that the fact that we impose stronger control over \tilde{u} than just the energy (indeed stronger than a Besov norm) should allow us to work with the original coordinate formulation, see e. g. [40]. However, it appears that the bilinear null-structure in (1), (2) is not good enough to obtain the gains in time we shall need, see Proposition 3.17. Indeed, proving an equivalent of this Proposition for the bilinear expressions appears to require time decay (in the sense that the norm evaluated on the function truncated to large times decays) for norms such as $\|u\|_{\dot{X}_k^{1, \frac{1}{2}, \infty}}$, which already fails for free waves. Moreover, our proof will actually reveal that one gets an honest H^1 -stability result provided one restricts oneself to large enough times.

there exists a smooth Wave Map \tilde{u} extending $\tilde{u}[0]$. Also, \tilde{u} will stay close to u in the sense that

$$\sup_t \|(\psi_\nu - \tilde{\psi}_\nu)(t)\|_{L_x^2} \lesssim \epsilon$$

The proof of this shall consist in analyzing the wave equation satisfied by the difference $\delta\psi_\nu := \tilde{\psi}_\nu - \psi_\nu$. Subtracting the wave equations for $\tilde{\psi}_\nu$, ψ_ν , and eliminating the $\tilde{\psi}_\nu$ results in terms linear, quadratic and cubic in $\delta\psi_\nu$. As these expressions have no apparent null-structure in them, we shall revert to the device of a Hodge-type decomposition used already in [22], [23]: we shall write $\psi_\nu = R_\nu\psi + \chi_\nu$ and similarly for $\tilde{\psi}_\nu$, where we impose the condition $\sum_{i=1,2} \partial_i \chi_i = 0$. Note that this results in a similar decomposition for $\delta\psi_\nu$. One easily deduces an elliptic div-curl system for χ_ν , $\tilde{\chi}_\nu$, from which one deduces the schematic identities $\chi_\nu = \nabla^{-1}(\psi \nabla^{-1}(\psi^2))$ etc., where the operators ∇^{-1} stand for linear combinations of operators of the form $\Delta^{-1}\partial_j$. Plugging these ingredients back into the wave equations satisfied by the $\delta\psi_\nu$ and eliminating all $\tilde{\psi}_\nu$ results in trilinear null-form terms as well terms of higher degree of linearity, either linear or of higher degree in the $\delta\psi_\nu$. All of this is just like in [23]. Terms which are at least quadratic in the $\delta\psi_\nu$ can be treated just as there, using the fact that Corollary 2.6 shall allow us to retrieve all the necessary estimates about ψ_ν . The only added difficulty comes from the terms linear in $\delta\psi_\nu$. One way to think of these is as an extra driving term added to the flat operator \square . However, the very good decay estimates satisfied by the ψ_ν shall allow us to treat these terms as source terms instead. The added difficulty over [23] we encounter here has to do with the fact that we need to gain explicitly in time. This will force us to localize simultaneously in physical and frequency space. In fact, we shall use a kind of wave packet decomposition to get the necessary estimates. The next two subsections provide the technical setup. In the same vein as the preceding theorem, we have

Theorem 3.3. *Let $u : \mathbf{R}^{2+1} \rightarrow \mathbf{H}^2$ be a smooth geodesic Wave Map with compactly supported initial data. Then for any $\mu > 0$, there exists $\epsilon = \epsilon(u, \mu) > 0$ such that for all smooth initial data $\tilde{u}[0] = (\tilde{u}(0), \partial_t \tilde{u}(0))$, with $\|(u - \tilde{u})(0)\|_{H^{1+\mu} \times H^\mu} < \epsilon$, there exists a smooth Wave Map $\tilde{u} : \mathbf{R}^{2+1} \rightarrow \mathbf{H}^2$ extending $\tilde{u}[0]$. \tilde{u} will stay close to u in the sense that $\sup_t \|(\psi_\nu - \tilde{\psi}_\nu)(t)\|_{L_x^2} \lesssim \epsilon$.*

3.4. Sobolev type spaces. We commence by introducing the functional analytic framework of [40], [37], [23] which we have to rely on to run the perturbation argument. We recall the Littlewood-Paley multipliers P_k introduced in the previous section:

$$\widehat{P_k f}(\xi) = \chi_0\left(\frac{|\xi|}{2^k}\right) \hat{f}(\xi),$$

for a suitable cutoff $\chi_0(\cdot)$. These are not flexible enough, and we also introduce the multipliers Q_j which localize the space-time Fourier support to dyadic distance $\sim 2^j$ from the light cone: letting

$$\tilde{\phi}(\xi, \tau) = \int_{\mathbf{R}^{2+1}} e^{-i(\tau t + \xi \cdot x)} \phi(t, x) dt dx$$

denote the space-time Fourier transform, we let

$$\widetilde{Q_j \phi}(\tau, \xi) := \chi_0\left(\frac{||\tau| - |\xi||}{2^j}\right) \tilde{\phi}(\tau, \xi)$$

where $\chi_0(\cdot)$ is as for the P_k 's. We note that these definitions entail the identities

$$\sum_{k \in \mathbf{Z}} P_k \phi = \phi, \quad \sum_{j \in \mathbf{Z}} Q_j \phi = \phi, \quad \phi \in \mathcal{S}(\mathbf{R}^{2+1}).$$

We have the basic inhomogeneous Sobolev spaces H^s , and their homogeneous counterparts \dot{H}^s :

$$\|\phi\|_{H^s} = \sqrt{\int_{\mathbf{R}^2} (1 + |\xi|^2)^s |\hat{\phi}(\xi)|^2 d\xi}, \quad \|\phi\|_{\dot{H}^s} = \sqrt{\int_{\mathbf{R}^2} |\xi|^{2s} |\hat{\phi}(\xi)|^2 d\xi}$$

Note that the space H^s is defined as completion of $\mathcal{S}(\mathbf{R}^2)$ with respect to the first norm. Trying to do the same for \dot{H}^s leads to difficulties (one gets not necessarily locally integrable functions). We shall only work with smooth functions anyways, so we only care about $\|\cdot\|_{\dot{H}^s}$. These norms are not flexible enough, and we also need the $X^{s,\theta}$ spaces of Klainerman-Machedon as well as their ('frequency localized') homogeneous Besov analogs (again only as norms):

$$\|\phi\|_{X^{s,\theta}} := \sqrt{\int_{\mathbf{R}^{2+1}} (1 + |\xi|^2)^s (1 + |\tau| - |\xi|)^{2\theta} |\tilde{\phi}(\tau, \xi)|^2 d\tau d\xi}, \quad \theta > \frac{1}{2}$$

$$\|\phi\|_{\dot{X}_k^{a,b,c}} := 2^{ka} \left(\sum_{j \in \mathbf{Z}} [2^{bj} \|Q_j \phi\|_{L_t^2 L_x^2}]^c \right)^{\frac{1}{c}}, \quad c < \infty$$

$$\|\phi\|_{\dot{X}_k^{a,b,\infty}} := 2^{ak} \sup_{j \in \mathbf{Z}} [2^{bj} \|Q_j \phi\|_{L_t^2 L_x^2}]$$

We shall always have $b = \frac{1}{2}$. The latter norms can be assembled to 'global versions', most naturally via

$$\|\phi\|_{\dot{X}^{a,b,c}} := \sqrt{\sum_{k \in \mathbf{Z}} \|\phi_k\|_{\dot{X}_k^{a,b,c}}^2}$$

The most intuitive way to think about the $X^{s,\theta}$ etc is to view them as superpositions of 'twisted free waves', gotten by foliating space-time by cones $|\tau| - |\xi| = \lambda$. One has the representation (see [17])

$$\phi = \int_{\lambda \in \mathbf{R}} \phi_\lambda e^{it\lambda} d\lambda,$$

where $\square \phi_\lambda = 0$ and

$$\int_{\lambda} \|\phi_\lambda\|_{H^s} d\lambda \lesssim \|\phi\|_{X^{s,\theta}}.$$

At the homogeneous level, we have the embedding⁸

$$\dot{X}_k^{0, \frac{1}{2}, 1} \subset L_t^\infty L_x^2.$$

More generally, the Strichartz estimates (see e. g. [9]) imply that the following embeddings hold:

$$\dot{X}_k^{0, \frac{1}{2}, 1} \subset 2^{k(1 - \frac{1}{p} - \frac{2}{q})} L_t^p L_x^q,$$

⁸The way to think about these is in the sense of inequalities between the associated norms: $A \subset B \rightarrow \|u\|_B \lesssim \|u\|_A$.

where $\frac{1}{p} + \frac{1}{2q} \leq \frac{1}{4}$. Similar embeddings hold for the 'subcritical spaces' $X^{s,\theta}$. We shall need slightly shrunk versions of the spaces $X^{s,\theta}$ etc. which give stronger control for the 'elliptic regions' far away from the light cone. For example, we have the norms (see [17]) $\|\cdot\|_{\mathcal{X}^{s,\theta}}$, which are defined via

$$\|\phi\|_{\mathcal{X}^{s,\theta}} := \|\phi\|_{X^{s,\theta}} + \|\partial_t \phi\|_{X^{s-1,\theta}}$$

Similarly, we introduce the space \mathcal{H}^s defined as the completion of $\mathcal{S}(\mathbf{R}^2)$ under the norm

$$\|\psi\|_{\mathcal{H}^s} := \|\psi\|_{H^s} + \|\partial_t \psi\|_{H^{s-1}}$$

3.5. Tataru's null-frame spaces. This subsection also summarizes material expounded in greater detail elsewhere (e. g. [37], [22], [23]). The spaces $X^{s,\theta}$ and their homogeneous Besov counterparts are unfortunately only part of the story. This has to do with the fact that even the strongest homogeneous versions of these norms (the norms $\|\cdot\|_{X^{a,\frac{1}{2},1}}$) do not yield good algebra type estimates, due to logarithmic divergences in low frequencies. A solution to this problem is given by 'spaces' incorporating Tataru's null-frame spaces. We present here a first version of spaces that overcome this difficulty. We shall construct norms $\|\cdot\|_S$ assembled from a family of 'frequency localized' norms $\|\cdot\|_{S[k]}$:

$$\|\phi\|_S := \sqrt{\sum_{k \in \mathbf{Z}} \|P_k \phi\|_{S[k]}^2}$$

The norms $\|\cdot\|_{S[k]}$ in turn are gotten as in [23]: they are constructed to satisfy $\|\cdot\|_{X_k^{0,\frac{1}{2},\infty}} \leq \|\cdot\|_{S[k]} \lesssim \|\cdot\|_{X_k^{0,\frac{1}{2},1}}$. We arrange that the norms are invariant under the natural scaling operation associated with derivatives of Wave Maps in $2+1$ dimensions, since we shall be working at the level of the derivative. The precise definition of $S[k]$ is complicated: we first construct norms $\|\cdot\|_{S[k,\kappa]}$ associated with every integer k and cap $\kappa \subset S^1$. To do so, we introduce *null-frame coordinates* (t_ω, x_ω) , $\omega \in S^1$, on space-time, whose definition is as follows:

$$t_\omega = \frac{1}{\sqrt{2}}(1, \omega) \cdot (t, x)$$

$$x_\omega = (t, x) - \frac{t_\omega}{\sqrt{2}}(1, \omega)$$

Thus these are Cartesian coordinates with respect to a tilted reference frame, whose 'time axis' with direction $\frac{1}{\sqrt{2}}(1, \omega)$ lies along the light cone. Now we introduce the space $PW[\kappa]$ defined as the atomic Banach space whose atoms are Schwartz functions $\psi \in \mathcal{S}(\mathbf{R}^{2+1})$ satisfying

$$\inf_{\omega \in \tilde{\kappa}} \|\psi\|_{L_{t_\omega}^2 L_{x_\omega}^\infty} \leq 1,$$

where $\tilde{\kappa}$ is a slightly grown version of κ (say by a factor $\frac{11}{10}$) concentric with it. Thus for $\psi \in \mathcal{S}(\mathbf{R}^{2+1})$, we have

$$\|\psi\|_{PW[\kappa]} := \inf_{\int_{\tilde{\kappa}} \psi_\omega d\omega = \psi} \int_{\tilde{\kappa}} \|\psi_\omega\|_{L_{t_\omega}^2 L_{x_\omega}^\infty} d\omega$$

Moreover, we put for ψ as above

$$\|\psi\|_{NFA[\kappa]^*} = \sup_{\omega \notin 2\kappa} \text{dist}(\omega, \kappa) \|\psi\|_{L_{t\omega}^\infty L_{x\omega}^2}$$

Now we put

$$\|\psi\|_{S[k,\kappa]} = \|\psi\|_{L_t^\infty L_x^2} + 2^{-\frac{k}{2}} |\kappa|^{-\frac{1}{2}} \|\psi\|_{PW[\kappa]} + \|\psi\|_{NFA[\kappa]^*}$$

This definition immediately entails the following fundamental first bilinear inequality

$$\|\phi\psi\|_{L_t^2 L_x^2} \lesssim \frac{|\kappa'|^{\frac{1}{2}} 2^{\frac{k'}{2}}}{\text{dist}(\kappa, \kappa')} \|\phi\|_{S[k,\kappa]} \|\psi\|_{S[k',\kappa']}, \quad (9)$$

provided $2\kappa \cap 2\kappa' = \emptyset$. We now construct the norms $\|\psi\|_{S[k]}$ by evaluating suitably microlocalized pieces of ψ with respect to the $\|\cdot\|_{S[k,\kappa]}$, taking a suitable mean and combining this with $\|\cdot\|_{\dot{X}_k^{a,b,c}}$ type norms. The null-frame norms may be thought of as controlling the 'free wave-like' character of ψ , while the remaining norms may be thought of as controlling the 'elliptic character' of ψ .

For every integer $l < -10$, subdivide S^1 into a uniformly finitely overlapping collection K_l of caps κ of diameter 2^l . Also, for every integer λ with $-10 \geq \lambda \geq l$, we subdivide the angular sector $\{\xi \in \mathbf{R}^2 \mid \frac{\xi}{|\xi|} \in \kappa, |\xi| \sim 2^k\}$ into a uniformly finitely overlapping collection $C_{k,\kappa,\lambda}$ of slabs R of width $2^{k+\lambda}$. We introduce various localization operators associated with these regions: for each $\kappa \in K_l$, choose a smooth cutoff $a_\kappa : S^1 \rightarrow \mathbf{R}_{\geq 0}$ supported on a dilate of κ . These are to be chosen such that $\sum_{\kappa \in K_l} a_\kappa = 1$. We also introduce cutoffs $m_R(\cdot) : \mathbf{R}_{>0} \rightarrow \mathbf{R}_{\geq 0}$ such that the cutoff $m_R(|\xi|)a_\kappa(\frac{\xi}{|\xi|})$ localizes to a dilate of the slab R . Also, we require that $\sum_{R \in C_{k,\kappa,\lambda}} m_R(|\xi|) = \chi_0(\frac{|\xi|}{2^k})$. We have the associated pseudo differential operator $\tilde{P}_R\psi$:

$$\widehat{\tilde{P}_R\psi}(t, \xi) = m_R(|\xi|)a_\kappa\left(\frac{\xi}{|\xi|}\right)\hat{\psi}(\xi)$$

We also have the Ψ DO's $P_{k,\kappa}$ associated with multiplier $a_\kappa(\frac{\xi}{|\xi|})\chi_0(\frac{|\xi|}{2^k})$. Then, almost⁹ as in [23] we define

$$\begin{aligned} \|\psi\|_{S[k]} := & \|\psi\|_{L_t^\infty L^2} + \|\psi\|_{\dot{X}_k^{0,\frac{1}{2},\infty}} + \|\psi\|_{\dot{X}_k^{-\frac{1}{2},1,2}} \\ & + \sup_{\pm} \sup_{l < -10} \sup_{-10 \geq \lambda \geq l} |\lambda|^{-1} \left(\sum_{\kappa \in K_l} \sum_{R \in C_{k,\kappa,\lambda}} \|\tilde{P}_R Q_{<k+2l}^\pm \psi\|_{S[k,\pm\kappa]}^2 \right)^{\frac{1}{2}}. \end{aligned} \quad (10)$$

This norm looks very complicated, but it isn't too hard to get control over its ingredients. A fundamental inequality [23] for example states that

$$\|P_k Q_{<k+O(1)} \psi\|_{S[k]} \lesssim \|P_k \psi\|_{\dot{X}_k^{0,\frac{1}{2},1}} \quad (11)$$

⁹The original definition also contained a norm $\|P_k Q_{\geq k} \partial_t \psi\|_{L_t^M \dot{H}^{-1+\frac{1}{M}}}$. Inspection of the proof there yields that this is superfluous, though, since the norm $\|\cdot\|_{\dot{X}_k^{-\frac{1}{2},-1,2}}$ suffices for the elliptic estimates.

We also need a norm $||\cdot||_N$ with respect to which we evaluate the nonlinearities of our wave equations. We shall put

$$||\psi||_N := \sqrt{\sum_{k \in \mathbf{Z}} ||P_k \psi||_{N[k]}^2},$$

where the $N[k]$ will be constructed as atomic Banach spaces. More precisely, we let $N[k]$ be the atomic Banach space whose atoms are Schwartz functions $F \in \mathcal{S}(\mathbf{R}^{2+1})$ with spatial Fourier support contained in the region $|\xi| \sim 2^k$ and

- (1) $||F||_{L_t^1 \dot{H}^{-1}} \leq 1$ and F has modulation $< 2^{k+100}$.
- (2) F is at modulation $\sim 2^j$ and satisfies $||F||_{L_t^2 L_x^2} \leq 2^{\frac{j}{2}} 2^k$.
- (3) F satisfies $||F||_{\dot{X}_k^{-\frac{1}{2}, -1, 2}} \leq 1$.
- (4) There exists an integer $l < -10$, and Schwartz functions F_κ with Fourier support in the region

$$\{(\tau, \xi) \mid \pm \tau > 0, \mid |\tau| - |\xi| \mid \leq 2^{k-2l-100}, \frac{\xi}{|\xi|} \in \pm \kappa\}$$

with the properties

$$F = \sum_{\kappa \in K_l} F_\kappa, \left(\sum_{\kappa \in K_l} ||F_\kappa||_{NFA[\kappa]}^2 \right)^{\frac{1}{2}} \leq 2^k$$

In the last inequality, $NFA[\kappa]$ denotes the dual of $NFA[\kappa]^*$ (the completion of $\mathcal{S}(\mathbf{R}^{2+1})$ with respect to $||\cdot||_{NFA[\kappa]^*}$) used in the definition of $S[k, \kappa]$: Thus $NFA[\kappa]$ is the atomic Banach space whose atoms F satisfy

$$\frac{1}{\text{dist}(\omega, \kappa)} ||F||_{L_{t\omega}^1 L_{x\omega}^2} \leq 1$$

for some $\omega \notin 2\kappa$.

Observe that $\dot{X}_k^{-1, -\frac{1}{2}, 1} \subset N[k]$, $P_k(L_t^1 \dot{H}^{-1}) \subset N[k]$.

This definition immediately entails the fundamental 2nd bilinear inequality (again $2\kappa \cap 2\kappa' = \emptyset$)

$$||\phi\psi||_{NFA[\kappa]} \lesssim \frac{2^{\frac{k'}{2}} |\kappa'|^{\frac{1}{2}}}{\text{dist}(\kappa, \kappa')} ||\phi||_{L_t^2 L_x^2} ||\psi||_{S[k', \kappa']} \quad (12)$$

We quickly summarize here the main properties of these spaces we shall need, all proved in [23]: all functions ϕ, ψ etc. below shall be in $\mathcal{S}(\mathbf{R}^{2+1})$.

(a): *Product type estimate*

$$||P_k Q_j (P_{k_1} \phi_1 P_{k_2} \phi_2)||_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \lesssim 2^{\min\{k_1, k_2\}} 2^{\min\{\frac{j - \min\{k, k_1, k_2\}}{4+}, 0\}} \\ 2^{\min\{\frac{\max\{k_1, k_2\} - j}{2}, 0\}} ||P_{k_1} \phi_1||_{S[k_1]} ||P_{k_2} \phi_2||_{S[k_2]}$$

(b): *Bilinear null-form estimate*: let R_ν denote the Riesz type operator $R_\nu = \frac{\partial_\nu}{\sqrt{-\Delta}}$. For $0 \leq p < \frac{1}{4}$, we have¹⁰

$$\begin{aligned} & \|P_k[R_1 P_{k_1} \psi_1 R_2 P_{k_2} \psi_2 - R_2 P_{k_1} \psi_1 R_2 P_{k_2} \psi_2]\|_{\dot{X}_0^{p,-p,2}} \\ & \leq C_p 2^{\frac{\min\{k_1, k_2, k\}}{2}} \prod_{i=1,2} \|P_{k_i} \psi_i\|_{S[k_i]} \end{aligned} \quad (13)$$

$$\begin{aligned} & \|P_k[R_1 P_{k_1} \psi_1 R_\nu P_{k_2} \psi_2 - R_2 P_{k_1} \psi_1 R_\nu P_{k_2} \psi_2]\|_{\dot{X}_0^{p,-p,2}} \\ & \leq C_p |k - k_1|^2 2^{\frac{\min\{k_1, k_2, k\}}{2}} \prod_{i=1,2} \|P_{k_i} \psi_i\|_{S[k_i]} \end{aligned}$$

(c): *Trilinear null-form estimates*: these arise upon formulating the derivative wave equations in the Coulomb Gauge and applying Hodge type decompositions, as explained below: let $I = \sum_{k \in \mathbf{Z}} P_k Q_{<k+100}$:

$$\begin{aligned} & \|\partial^\beta P_0[R_\alpha P_{k_1} \psi_1 \Delta^{-1} \sum_{j=1}^2 \partial_j I[R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_\beta P_{k_3} \psi_3 R_j P_{k_2} \psi_2]] \\ & + \partial_\alpha P_0[R_\beta P_{k_1} \psi_1 \Delta^{-1} \sum_{j=1}^2 \partial_j I[R^\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R^\beta P_{k_3} \psi_3]]\|_{N[0]} \quad (14) \\ & \leq C 2^{\delta_1 \min\{-\min\{k_1, k_2, k_3\}, 0\}} \prod_i 2^{\delta_2 \min\{\max_{j \neq i} \{k_i, k_i - k_j\}, 0\}} \prod_l \|P_{k_l} \psi_l\|_{S[k_l]}, \\ & \|P_0 \partial^\beta [R_\beta P_{k_1} \psi_1 \Delta^{-1} \sum_j \partial_j I[R_\alpha P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\alpha P_{k_3} \psi_3]]\|_{N[0]} \\ & \leq C 2^{\delta_1 \min\{-\min\{k_1, k_2, k_3\}, 0\}} \prod_i 2^{\delta_2 \min\{\max_{j \neq i} \{k_i, k_i - k_j\}, 0\}} \prod_l \|P_{k_l} \psi_l\|_{S[k_l]}. \end{aligned} \quad (15)$$

Of course one may rescale these, i. e. replace P_0 by P_k , $k \in \mathbf{Z}$. Then one needs to replace $\min\{-\min\{k_1, k_2, k_3\}, 0\}$ by $\min\{-\min\{k_1 - k, k_2 - k, k_3 - k\}, 0\}$ etc.

(d): '*energy inequality*' The following relates the spaces $S[k]$ and $N[k]$:

$$\begin{aligned} \|P_k \psi\|_{S[k]([-T, T] \times \mathbf{R}^2)} & \lesssim \inf_{0 < T_0 \leq T} [\min\{2^k T_0, 1\}^{-\frac{1}{2}} \|\square P_k \psi\|_{N[k]([-T, T] \times \mathbf{R}^2)} \\ & + \sup_{t_0 \in [-T_0, T_0]} \|P_k \psi[t_0]\|_{L^2 \times \dot{H}^{-1}}]. \end{aligned} \quad (16)$$

In this inequality, one can leave out the factor $\min\{2^k T_0, 1\}^{-\frac{1}{2}}$ and replace $S[k]$ with the following stronger norm

$$\begin{aligned} \|\psi\|_{S'[k]} & := 2^{-k} \|\nabla_{x,t} \psi\|_{L_t^\infty L_x^2} + \|\nabla_{x,t} \psi\|_{\dot{X}_k^{-1, \frac{1}{2}, \infty}} \\ & + \sup_{\pm} \sup_{l < -10} \sup_{-10 \leq \lambda \leq l} |\lambda|^{-1} \left(\sum_{\kappa \in K_l} \sum_{R \in C_{k, \kappa, \lambda}} \|\tilde{P}_R Q_{<k+2l}^\pm \psi\|_{S[k, \pm \kappa]}^2 \right)^{\frac{1}{2}}, \end{aligned}$$

provided the norm $N[k]$ is replaced by $L_t^1 \dot{H}^{-1}$ or $\dot{X}_k^{-1, -\frac{1}{2}, 1}$.

(e): *Relation to Strichartz type spaces*: Let p, q satisfy $\frac{1}{p} + \frac{1}{2q} < \frac{1}{4}$. Then we have

$$\|P_0 \psi\|_{L_t^p L_x^q} \leq C_{p,q} \|P_0 \psi\|_{S[0]}$$

¹⁰Moreover, applying an operator Q_l in front, where $l \gg k$, we can include factors $2^{\frac{\min\{l-k_2, 0\}}{2}}$ on the right hand side.

(f): *Relation to improved Strichartz type estimates:* for $l < -10$, let $C_{0,l}$ be a covering of the frequency region $|\xi| \sim 1$ by uniformly finitely overlapping discs of radius $\sim 2^l$. Let $P_{0,c}$ localize the Fourier support to the disc c , such that $\sum_{c \in C_{0,l}} P_c = P_0$. Let $8 \geq p > 4$. Then we have

$$\left(\sum_{c \in C_{0,l}} \|P_{0,c}\psi\|_{L_t^p L_x^q}^2 \right)^{\frac{1}{2}} \leq C_p 2^{(\frac{3}{4+} - \frac{2}{p})l} \|\psi\|_{S[0]}$$

(g): *Bilinear inequality relating the $S[k]$, $N[k]$:* Let $F, \psi \in \mathcal{S}(\mathbf{R}^{2+1})$. Then for $j \leq \min\{k_1, k_2\} + O(1)$ we have

$$\begin{aligned} & \|P_k[P_{k_1}\psi P_{k_2}Q_j F]\|_{N[k]} \\ & \lesssim 2^{\delta(j - \min\{k_1, k_2\})} 2^{\min\{k_1 - k_2, 0\}} \|P_{k_1}\psi_1\|_{S[k_1]} \|P_{k_2}F\|_{\dot{X}_{k_2}^{0, -\frac{1}{2}, \infty}} \end{aligned}$$

If $k_1 = k_2 + O(1)$ $j \leq r + k$ for some $r \leq 0$, we have

$$\|P_k Q_{<r+k}[P_{k_1}Q_j F P_{k_2}Q_{<2k+r-k_1}\psi]\|_{N[k]} \lesssim 2^{\delta r} \|F\|_{\dot{X}_{k_1}^{0, -\frac{1}{2}, \infty}} \|P_{k_2}\psi_2\|_{S[k_2]}.$$

3.6. A modification of the spaces; Moser type estimates. In spite of the above properties, the spaces $S[k]$ don't appear flexible enough to handle Moser type estimates of the kind we shall need. More precisely, the property

$$\|\psi_1 A(\nabla^{-1}\psi_2)\|_S \lesssim C(\|\psi_1\|_S, \|\psi_2\|_S)$$

where $\psi_{1,2} \in \mathcal{S}(\mathbf{R}^2)$, $A(\cdot) \in C^\infty(\mathbf{R})$ with bounded derivatives, and ∇^{-1} schematic notation for linear combinations of operators of the form $\Delta^{-1}\partial_j$, appears violated. This is a consequence of the fact that the product estimate 3.4(a) does not allow us to recover enough exponential gains in the difference $k - k_1$ if $k_1 \gg k$, the high-high interaction case.

One way around this would be to re-engineer the way functions get subdivided into 'free wave parts' and 'elliptic parts'. Indeed, one has better product estimates than 3.4(a) for free waves, see Klainerman-Foschi¹¹ [11]. Our way here around this shall be to 'enlarge the space $S[k]$ ', shrinking the norm $\|\cdot\|_{S[k]}$ suitably. More precisely, we analyze the 'bad high-high' frequency interactions and observe that by virtue of the spherical symmetry assumption, these cases are actually favorable in some sense. Indeed, we shall be able to exploit the well-known fact (e. g. [30], [32]) that the range of admissible Strichartz estimates is significantly improved in this situation:

Theorem 3.7. *Let $\psi \in \mathcal{S}(\mathbf{R}^{2+1})$ be invariant under rotations. Then we have the inequality*

$$\|P_0\psi\|_{L_t^p L_x^q} \lesssim \|\psi[0]\|_{L_t^\infty L_x^2} + \|\Box\psi\|_{L_t^1 L_x^2}$$

provided the condition $\frac{1}{p} + \frac{1}{q} < \frac{1}{2}$ holds.

We note that theorem 3.7 implies easily the following corollary:

¹¹It appears, however, that even for free waves, a high-high-low interaction resulting in an elliptic product (Fourier support very far from the light cone) does not lead to the desired exponential gain in the frequencies.

Corollary 3.8. *The derivative components ψ_ν satisfy the estimates*

$$\|P_k \psi_\nu\|_{L_t^p L_x^q} \lesssim c_k$$

for a system of numbers $\{c_k\}_{k \in \mathbf{Z}}$ ('frequency envelope') which satisfies $\sum_{k \in \mathbf{Z}} 2^{\delta|k|} c_k < \infty$, provided the condition $\frac{1}{p} + \frac{1}{q} < \frac{1}{2}$ holds and $\delta > 0$ is sufficiently small.

Proof This follows by applying a simple frequency trichotomy to the frequency localized expression for ψ_ν in terms of $\frac{\partial_\nu \mathbf{x}}{\mathbf{y}}, \frac{\partial_\nu \mathbf{y}}{\mathbf{y}}$. The latter are controlled by application of theorem 3.7 as well as Corollary 2.6. \blacksquare

Definition 3.9. : We put

$$\|\psi\|_L := \sup_{(p,q) | \frac{1}{p} + \frac{1}{2q} \leq \frac{1}{4}} 2^{k(\frac{1}{p} + \frac{2}{q} - 1)} \sup_{l \leq 0} 2^{-\max\{\frac{1}{2+} - \frac{1}{q}, 0\}l} \left(\sum_{c \in C_{k,l}} \|P_c \psi\|_{L_t^p L_x^q}^2 \right)^{\frac{1}{2}}$$

Here $C_{k,l}$ is a finitely overlapping cover of the frequency region $|\xi| \sim 2^k$ by discs of radius 2^{k+l} , with associated Fourier localizers P_c , $c \in C_{k,l}$. Also, put $2+ = 2 + \frac{1}{1000}$, say, and let μ be a small positive number, say $\frac{1}{1000}$. We let $\mathcal{S}[k]$ be the atomic Banach space whose atoms satisfy one of the following:

- (1) *Type 1 atoms:* Fix δ_0 small, say $\delta_0 = \frac{1}{1000}$. These are functions $\psi \in \mathcal{S}(\mathbf{R}^{2+1})$ satisfying

$$\begin{aligned} \|\psi\|_{A[k]} &= \|\psi\|_L + \|\psi\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} + \|\psi\|_{\dot{X}_k^{-\frac{1}{2}, 1, 2}} + \sup_{\frac{1}{p} + \frac{1}{q} < \frac{1}{2} - \delta_0} 2^{(\frac{1}{p} + \frac{2}{q} - 1)k} \|\psi\|_{L_t^p L_x^q} \\ &+ \sup_{\pm} \sup_{l < -10} \sup_{-10 \leq \lambda \leq l} |\lambda|^{-1} \left(\sum_{\kappa \in K_l} \sum_{R \in C_{k, \kappa, \lambda}} \|\tilde{P}_R Q_{< k+2l}^\pm \psi\|_{\mathcal{S}[k, \pm \kappa]}^2 \right)^{\frac{1}{2}} \leq 1 \end{aligned}$$

- (2) *Atoms of the 2nd type:* Let $I = \sum_{k \in \mathbf{Z}} P_k Q_{< k+100}$. Then $\psi \in \mathcal{S}(\mathbf{R}^{2+1})$ is of the 2nd type provided

$$\begin{aligned} \|\psi\|_{B[k]} &:= \sup_{\frac{1}{p} + \frac{1}{q} < 1 - 10\delta_0} 2^{k[\frac{1}{p} + \frac{2}{q} - 1]} \|\psi\|_{L_t^p L_x^q} + \|\psi\|_L \\ &+ \|(1 - I)\psi\|_{\dot{X}_k^{-(\frac{1}{2} - \mu), 1 - \mu, 1}} \leq 1, \end{aligned}$$

The range of Lebesgue exponents (p, q) includes the pairs $(1+, \infty)$, $(\infty, 1+)$.

Note that any function $\psi \in \mathcal{S}(\mathbf{R}^{2+1})$ may be decomposed into two pieces $\psi = \alpha + \beta$ satisfying

$$\|\alpha\|_{A[k]} \lesssim \|\psi\|_{\mathcal{S}[k]}, \quad \|\beta\|_{B[k]} \lesssim \|\psi\|_{\mathcal{S}[k]}$$

We call α 'of first type' and β 'of 2nd type'. We let

$$\|\psi\|_{\mathcal{S}} := \left(\sum_{k \in \mathbf{Z}} \|P_k \psi\|_{\mathcal{S}[k]}^2 \right)^{\frac{1}{2}}$$

Unfortunately, these norms do not quite suffice to close all the estimates. The following theorem contains some bilinear estimates, which we were unable to build into a linear framework. These have to be proved independently. To state the theorem, we use the concept of *frequency envelope*: following Tao [36], we call a sequence of positive numbers $\{c_k\}_{k \in \mathbf{Z}}$ a frequency envelope provided $C^{-1} c_b 2^{-\sigma|a-b|} \leq c_a \leq C 2^{\sigma|a-b|} c_b$ for some $\sigma > 0$ and $C \gg 1$.

Theorem 3.10. *Let $\psi_{1,2} \in \mathcal{S}(\mathbf{R}^{2+1})$. Assume that $\|P_k \psi_{1,2}\|_{\mathcal{S}[k]} \lesssim c_k$ for a sufficiently flat frequency envelope $\{c_l\}$. Let*

$$\|\nabla^{-1} \psi_2\|_{L_t^\infty L_x^\infty} \leq C_0$$

More precisely, assume that for each $k \in \mathbf{Z}$ one may split $P_k \psi_2 = \alpha_k + \beta_k$ into functions of first and 2nd type, respectively, such that $\|\nabla^{-1} \sum_k \alpha_k\|_{L^\infty} \leq C_0$, $\|\nabla^{-1} \sum_k \beta_k\|_{L^\infty} \leq C_0$. Also assume that ψ_2 satisfies the bilinear estimates stated further below¹². Then we can conclude¹³ $\forall k \in \mathbf{Z}$

$$\|P_k[\psi_1 \nabla^{-1} \psi_2]\|_{\mathcal{S}[k]} \lesssim c_k$$

In particular, if $A(\cdot) : \mathbf{R} \rightarrow \mathbf{C}$ is real analytic with bounded derivatives of arbitrary order, and ψ_2 real valued, we can conclude that $\forall k \in \mathbf{Z}$

$$\|P_k[\psi_1 A(\nabla^{-1} \psi_2)]\|_{\mathcal{S}[k]} \lesssim c_k$$

A more precise version is as follows: for suitable $\delta > 0$,

$$\|P_k[P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2]\|_{\mathcal{S}[k]} \lesssim 2^{-\delta|k-k_1|} \|P_{k_1} \psi_1\|_{\mathcal{S}[k_1]} \|P_{k_2} \psi_2\|_{\mathcal{S}[k_2]}$$

Now assume that $\|\psi_{1,2}\|_{\mathcal{S}} + \|\psi_{1,2}\|_{\dot{B}_2^{0,1}} < C_0$. Then we can conclude

$$\|P_k[\psi_1 A(\nabla^{-1} \psi_2)]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \leq C,$$

for a suitable $C = C(C_0)$. More precisely, decomposing

$$P_k[\psi_1 A(\nabla^{-1} \psi_2)] = \alpha + \beta$$

into functions of first and 2nd type, respectively, we may assume that

$$\|R_0(\alpha + \beta)\|_{L_t^\infty L_x^2} \leq C, \|\beta\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \leq C$$

for $C = C(C_0, \sup_{k \in \mathbf{Z}} \|R_0 P_k \psi_1\|_{L_t^\infty L_x^2}, \sup_{k \in \mathbf{Z}} \|R_0 P_k \psi_2\|_{L_t^\infty L_x^2})$. Next, we have the bilinear estimates

$$\begin{aligned} & \sup_{\phi | \|\phi\|_{\mathcal{S}[k_3]} \leq 1} \sup_{k_1, 2, 3} \sup_{l < -10} 2^{-\frac{\min\{k_1, k_3\}}{2}} 2^{-\delta_1 l} \\ & \left(\sum_{c \in C_{k_3, l}} \|P_{k_1}[P_{k_2} R_0[\psi_1 A(\nabla^{-1} \psi_2)] P_c \phi]\|_{L_t^2 L_x^2}^2 \right)^{\frac{1}{2}} \lesssim C \end{aligned}$$

$$\sup_{\phi | \|\phi\|_{\mathcal{S}[k_3]} \leq 1} \sup_{k_1, 2 \in \mathbf{Z}} 2^{\frac{k_2(1-3\mu)}{2+2\mu}} \|P_{k_1}[P_{k_2} R_0[\psi_1 A(\nabla^{-1} \psi_2)] P_{k_2+O(1)} \phi]\|_{L_t^2 L_x^{1+\mu}} \leq C,$$

where $C = C(\|\psi_1\|_{\mathcal{S}}, \|\psi_2\|_{\mathcal{S}}, \|\psi_2\|_{\dot{B}_2^{0,1}})$, recall the definition in 3.4. The same estimates hold provided one replaces $A(\nabla^{-1} \psi_2)$ by $A(\nabla^{-1}(\psi_2 \nabla^{-1} \psi_3))$, where ψ_3 satisfies similar estimates as ψ_2 .

The proof of this is a long calculation deferred to an appendix.

¹²Substitute ψ_2 instead of $\psi_1 A(\nabla^{-1} \psi_2)$.

¹³The implied constants in the statements below depend on C_0 , the constant C in the bilinear estimates below as well as the decay of the frequency envelope and the constants chosen in the definition of $\mathcal{S}[k]$.

3.11. Proof of the Theorem 3.2. To show existence, it suffices to show that some subcritical norm $\sum_{\nu=0}^2 \|\delta\psi_\nu(t)\|_{H^\delta}$ is globally bounded in t . Indeed, reasoning exactly as in [23], one deduces that for every finite time interval $[-T, T]$, one has with $\tilde{u} = (\tilde{\mathbf{x}}, \tilde{\mathbf{y}})$ and $\tilde{\delta} < \delta$:

$$\sup_{t \in [-T, T]} \sum_{\nu=0}^2 [\|\tilde{\mathbf{x}}(t)\|_{H^{1+\tilde{\delta}}} + \|\tilde{\mathbf{y}}(t)\|_{H^{1+\tilde{\delta}}} + \|\partial_t \tilde{\mathbf{x}}(t)\|_{H^{\tilde{\delta}}} + \|\partial_t \tilde{\mathbf{y}}(t)\|_{H^{\tilde{\delta}}}] < \infty$$

Using the subcritical result of

Klainerman-Machedon [12], one deduces from here that there can't be breakdown of smoothness after finite time.

Global boundedness of a subcritical norm in turn shall follow from the following Bootstrap Proposition: to formulate it, we shall need time-localized versions of the spaces $S[k]$: for $\psi \in C_0^\infty([-T, T] \times \mathbf{R}^2)$ define

$$\|\psi\|_{S[k]([-T, T] \times \mathbf{R}^2)} := \inf_{\tilde{\psi} \in S(\mathbf{R}^{2+1}) | \tilde{\psi}|_{[-T, T] \times \mathbf{R}^2} = \psi} \|\tilde{\psi}\|_{S[k]}$$

We use similar definitions for $\|\cdot\|_{S([-T, T] \times \mathbf{R}^2)}$ $\|\psi\|_{S([-T, T] \times \mathbf{R}^{2+1})}$ etc. and also different time intervals $[T_1, T_2]$ etc.

Proposition 3.12. *In the situation of Theorem 3.2, let the smooth Wave Map extending $\tilde{u}[0]$ exist on the time interval $[-T, T]$. There exists $T_1 > 0$ such that for $T \geq T_1$ and every $K > 0$ sufficiently large, there exists $\epsilon > 0$ such that the following conclusion applies: Introduce the frequency envelope*

$$\tilde{c}_k := \sup_{t \in [-T_1, T_1]} \sum_{k_1 \in \mathbf{Z}} 2^{-\sigma|k-k_1|} \|P_{k_1} \delta\psi_\nu(t)\|_{L_x^2} + \epsilon c_k,$$

where $\{c_k\}$ is as in the proof of lemma 3.16, and assume

$\sup_{t \in [-T_1, T_1]} \sum_{\nu=0,1,2} \|\delta\psi_\nu(t)\|_{L_x^2} \leq \epsilon$. Then for any $T \geq T_1$ we have

$$\sup_{\nu} \|P_k \delta\psi_\nu\|_{S[k]([T_1, T] \times \mathbf{R}^2)} < K \tilde{c}_k \Rightarrow \|P_k \delta\psi_\nu\|_{S[k]([T_1, T] \times \mathbf{R}^2)} < \frac{K}{2} \tilde{c}_k.$$

A similar inequality holds by replacing T, T_1 by $-T, -T_1$.

Assuming this for now, we continue with the proof of Theorem 3.2. We claim that by local well-posedness of (1), (2), there exists $\epsilon = \epsilon(\mu) > 0$ such that (using terminology of Theorem 3.2) $\|u[0] - \tilde{u}[0]\|_{H^{1+\mu} \times H^\mu} < \epsilon$ implies that \tilde{u} extends smoothly to $[-T_1, T_1]$, where T_1 is as in the preceding Proposition. To see that this is possible, we shall apply an inequality of Klainerman-Selberg to the equation satisfied by the differences $\tilde{\mathbf{x}} - \mathbf{x}, \tilde{\mathbf{y}} - \mathbf{y}$ of the coordinate representations of the perturbed and the spherically symmetric Wave Map, $\tilde{u} = (\tilde{\mathbf{x}}, \tilde{\mathbf{y}})$ and $u = (\mathbf{x}, \mathbf{y})$. Subdivide the interval $[-T_1, T_1]$ into small subintervals I_i , for which¹⁴ $\|(\mathbf{x}, \mathbf{y})\|_{\mathcal{X}^{1+\mu, \theta}(I_i \times \mathbf{R}^2)} \lesssim 1$. This is possible by Corollary 2.6 and local well-posedness of (1), (2) in $H^{1+\mu}$. Note that

$$\square(\ln \tilde{\mathbf{y}} - \ln \mathbf{y}) = \left(\frac{\partial_\nu \tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} - \frac{\partial_\nu \mathbf{x}}{\mathbf{y}} \right) \frac{\partial^\nu \tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} + \frac{\partial^\nu \mathbf{x}}{\mathbf{y}} \left(\frac{\partial_\nu \tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} - \frac{\partial_\nu \mathbf{x}}{\mathbf{y}} \right)$$

¹⁴The $\theta > \frac{1}{2}$ is chosen in dependence of μ , see Klainerman-Selberg [17].

with a similar equation holding for $\square(\frac{\tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} - \frac{\mathbf{x}}{\mathbf{y}})$. We deduce that upon denoting $I_i = [a_i, a_{i+1}]$, we have

$$\begin{aligned} & \sum_{\nu} \left\| \frac{\partial_{\nu} \tilde{\mathbf{y}}}{\tilde{\mathbf{y}}} - \frac{\partial_{\nu} \mathbf{y}}{\mathbf{y}} \right\|_{\mathcal{X}^{\mu, \theta}(I_i \times \mathbf{R}^2)} + \left\| \partial_{\nu} \left(\frac{\tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} \right) - \partial_{\nu} \left(\frac{\mathbf{x}}{\mathbf{y}} \right) \right\|_{\mathcal{X}^{\mu, \theta}(I_i \times \mathbf{R}^2)} \\ & \lesssim \sum_{\nu} \left\| \frac{\partial_{\nu} \tilde{\mathbf{y}}}{\tilde{\mathbf{y}}} - \frac{\partial_{\nu} \mathbf{y}}{\mathbf{y}} \right\|_{H^{\mu}(a_i)} + \left\| \partial_{\nu} \left(\frac{\tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} \right) - \partial_{\nu} \left(\frac{\mathbf{x}}{\mathbf{y}} \right) \right\|_{H^{\mu}(a_i)} \\ & + |I_i|^{\epsilon(\mu)} [\|\ln \tilde{\mathbf{y}} - \ln \mathbf{y}\|_{\mathcal{X}^{1+\mu, \theta}(I_i \times \mathbf{R}^2)} + \left\| \frac{\tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} - \frac{\mathbf{x}}{\mathbf{y}} \right\|_{\mathcal{X}^{1+\mu, \theta}(I_i \times \mathbf{R}^2)}] \\ & + [|I_i|^{\epsilon(\mu)} \|\ln \tilde{\mathbf{y}} - \ln \mathbf{y}\|_{\mathcal{X}^{1+\mu, \theta}(I_i \times \mathbf{R}^2)}^3 + \left\| \frac{\tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} - \frac{\mathbf{x}}{\mathbf{y}} \right\|_{\mathcal{X}^{1+\mu, \theta}(I_i \times \mathbf{R}^2)}^3] \end{aligned}$$

We have used here the fact, due to Klainerman-Selberg, that

$$\|\phi\|_{\mathcal{X}^{s, \theta}([-T, T] \times \mathbf{R}^2)} \lesssim \|\phi[0]\|_{H^s \times H^{s-1}} + T^{\epsilon} \|\square \phi\|_{\mathcal{X}^{s-1, \theta-1}([-T, T] \times \mathbf{R}^2)}, \quad \theta > \frac{1}{2}, s > 1$$

as well as the following inequality of Klainerman-Machedon

$$\|\partial_{\nu} u_1 \partial^{\nu} u_2\|_{X^{s-1, \theta-1}} \lesssim \|u_1\|_{\mathcal{X}^{s, \theta}} \|u_2\|_{\mathcal{X}^{s, \theta}}$$

Refining the subdivision $[-T_1, T_1] = \bigcup_{i=1}^N I_i$, $N = N(u, \mu)$, if necessary, we see that

$$\begin{aligned} & \left\| \sum_{\nu} \left\| \frac{\partial_{\nu} \tilde{\mathbf{y}}}{\tilde{\mathbf{y}}} - \frac{\partial_{\nu} \mathbf{y}}{\mathbf{y}} \right\|_{H^{\mu}(a_{i+1})} + \left\| \partial_{\nu} \left(\frac{\tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} \right) - \partial_{\nu} \left(\frac{\mathbf{x}}{\mathbf{y}} \right) \right\|_{H^{\mu}(a_{i+1})} \right\| \\ & \lesssim 2 \left\| \sum_{\nu} \left\| \frac{\partial_{\nu} \tilde{\mathbf{y}}}{\tilde{\mathbf{y}}} - \frac{\partial_{\nu} \mathbf{y}}{\mathbf{y}} \right\|_{H^{\mu}(a_i)} + \left\| \partial_{\nu} \left(\frac{\tilde{\mathbf{x}}}{\tilde{\mathbf{y}}} \right) - \partial_{\nu} \left(\frac{\mathbf{x}}{\mathbf{y}} \right) \right\|_{H^{\mu}(a_i)} \right\|, \end{aligned}$$

provided the quantity on the right hand side is less than some constant c . Thus if we choose $\epsilon < \frac{c}{2N}$, we see that the Wave Map \tilde{u} satisfying $\|\tilde{u}[0] - u[0]\|_{H^{1+\mu} \times H^{\mu}} \ll \epsilon$ will exist and be smooth on the interval $[-T_1, T_1]$. It follows from the argument just given and a simple algebra type estimate that by possibly shrinking the size of $\|u[0] - \tilde{u}[0]\|_{H^{1+\mu} \times H^{\mu}}$ we can ensure that

$$\|\delta \psi_{\nu}\|_{L_t^{\infty} H^{\lambda}([-T_1, T_1] \times \mathbf{R}^2)} \leq \epsilon, \quad 0 < \lambda < \mu.$$

Now assume that the perturbed Wave Map \tilde{u} breaks down at some time $T > T_1$. We claim that $\sup_{T_1 \leq t < T} \sup_{k \in \mathbf{Z}} \tilde{c}_k^{-1} \|P_k \delta \psi_{\nu}\|_{S[k]([T_1, t] \times \mathbf{R}^2)} < \infty$. Indeed, in the opposite case, choosing K large enough (and if necessary shrinking ϵ), and using the continuity of the function $t \rightarrow \sup_k \tilde{c}_k^{-1} \|P_k \delta \psi_{\nu}\|_{S[k]([T_1, t] \times \mathbf{R}^2)}$ for $t \in [T_1, T)$, see e. g. [23], it follows that there exists T' satisfying the properties

$$\sup_{k \in \mathbf{Z}} \tilde{c}_k^{-1} \|P_k \delta \psi_{\nu}\|_{S[k]([T_1, T'] \times \mathbf{R}^2)} = K, \quad T_1 < T' < T$$

This, however, contradicts Proposition 3.12. This then implies that $\sup_{t < T} \|P_k \delta \psi_{\nu}\|_{S[k]([-t, t] \times \mathbf{R}^2)} \lesssim \tilde{c}_k$. But by definition of \tilde{c}_k , $\|\cdot\|_{S[k]}$, this implies that some subcritical norm $\|\delta \psi_{\nu}\|_{L_t^{\infty} H^{\mu}([-T, T] \times \mathbf{R}^2)} < \infty$, $\mu > 0$, which in turn implies that some $\|\tilde{u}\|_{L_t^{\infty} H^{1+\mu'}([-T, T] \times \mathbf{R}^2)} < \infty$. This in turn contradicts breakdown by the result of Klainerman-Machedon [12]. Of course, the preceding argument entails the bound $\|\delta \psi_{\nu}\|_{L_t^{\infty} L_x^2} \lesssim \epsilon$. Indeed, one obtains that some range of subcritical norms $\|\cdot\|_{H^{\delta}}$ satisfy that estimate.

3.13. Proof of Proposition 3.12. We first recall the following theorem from [23]: let $N(\{\psi_\nu\})$ be the nonlinearity on the right hand side of (8). Then we have

Theorem 3.14. [23] *Let $\psi_\nu \in C_0^\infty([-T, T] \times \mathbf{R}^2)$ solve the combined system (6), (7). Then provided*

$$\sup_\nu \|\psi_\nu\|_{S([-T, T] \times \mathbf{R}^2)} < K\epsilon$$

and ϵ is sufficiently small in relation to K , we have

$$\|N(\{\psi_\nu\})\|_{N([-T, T] \times \mathbf{R}^2)} \lesssim K^3 \epsilon^3$$

More precisely, there exists a number $\delta_1 > 0$ such that we have

$$\|P_k N(\{\psi_\nu\})\|_{N[k]([-T, T] \times \mathbf{R}^2)} \lesssim \left(\sup_\nu \sup_{k_1 \in \mathbf{Z}} 2^{-\delta_1 |k - k_1|} \|P_{k_1} \psi_\nu\|_{S[k_1]([-T, T] \times \mathbf{R}^2)} \right) K^2 \epsilon^2$$

The proof in [23] of this relied on introduction of null-form structure into the nonlinearities by means of Hodge type decompositions, as briefly outlined in subsection 3.1. Thus writing $\psi_\nu = R_\nu \psi + \chi_\nu$ and requiring $\sum_{i=1,2} \partial_i \chi_i = 0$ results in

$$\chi_\nu = i \sum_{i,j=1}^2 \partial_i \Delta^{-1} (\psi_\nu \Delta^{-1} \partial_j (\psi_i^1 \psi_j^2 - \psi_j^1 \psi_i^2) - \psi_i \Delta^{-1} \partial_j (\psi_\nu^1 \psi_j^2 - \psi_j^1 \psi_\nu^2)). \quad (17)$$

$$\psi = - \sum_{i=1,2} R_i \psi_i \quad (18)$$

One now writes the nonlinearities $N(\{\psi_\nu\})$ as sums of various terms which are gotten by substituting either gradient components $R_\nu \psi$ or elliptic components χ_ν in place of ψ_ν , substituting¹⁵ the Schwartz extensions ρ_ν for ψ_ν which satisfy

$$\|\rho_\nu\|_{S([-T, T] \times \mathbf{R}^2)} \lesssim K\epsilon,$$

and further microlocalizing constituents of the expressions thus obtained. One thereby obtains trilinear null-forms of the types recorded in 3.4(c). Substituting elliptic components χ_ν results in terms at least quintilinear in the variables ψ_ν , which are more elementary to estimate, but still appear to require null-form structure, which is obtained upon reiterating the Hodge type decomposition. One keeps going like this until the error terms obtained can be estimated without using null-structures, based only on Strichartz type estimates. Summarizing, we have

Theorem 3.15. [23] *Under the hypotheses of Theorem 3.14, we can construct a function*

$$\tilde{N}(\{\rho_\nu\}) \in \mathcal{S}(\mathbf{R}^{2+1})$$

which is expressible as a sum of terms trilinear, quadrilinear etc. up to degree 11 in the ρ_ν , and satisfies

$$\tilde{N}(\{\rho_\nu\})|_{[-T, T]} = N(\psi_\nu)$$

$$\|P_k \tilde{N}(\{\rho_\nu\})\|_{S[k]} \lesssim \left(\sup_\nu \sup_{k_1 \in \mathbf{Z}} 2^{-\delta |k - k_1|} \|P_{k_1} \rho_\nu\|_{S[k_1]} \right) \max\{\|\rho_\nu\|_S^2, \|\rho_\nu\|_S^{10}\}$$

¹⁵One reexpresses ψ, χ_ν in terms of ψ_ν via (17), (18).

We shall apply this theorem to our situation. The complication that arises here has to do with the fact that the estimates for $\{\psi_\nu\}$, the derivative components of the spherically symmetric Wave Map, are not with respect to $\|\cdot\|_S$, but rather $\|\cdot\|_{S'}$, in view of theorem 3.10. We state here

Lemma 3.16. *For any $\sigma > 0$ sufficiently small, there exists a frequency envelope $\{c_l\}_{l \in \mathbf{Z}}$ with exponent σ and $\sum_{l \in \mathbf{Z}} c_l^2 \lesssim 1$ such that $\forall T > 0$ we have*

$$\|P_k \psi_\nu\|_{S[k]([-T, T] \times \mathbf{R}^2)} \lesssim c_k$$

We can also assume $\sum_l 2^{\mu|l|} c_l \lesssim 1$ for $\mu > 0$ sufficiently small. Moreover, choosing a Schwartz extension $\widetilde{P_k \psi_\nu}$ of $P_k \psi_\nu|_{[-T, T]}$ satisfying the above estimates, we may decompose each $\widetilde{P_k \psi_\nu}$ into functions of first and 2nd type, $\widetilde{P_k \psi_\nu} = \alpha_\nu + \beta_\nu$, such that the following properties hold:

$$\|\alpha_\nu\|_{A[k]} \lesssim c_k, \|\beta_\nu\|_{B[k]} \lesssim c_k, \|\beta_\nu\|_{X_k^{0, \frac{1}{2}, 1}} \leq C, \|R_0(\alpha_\nu + \beta_\nu)\|_{L_t^\infty L_x^2} \leq C$$

where C depends on $\|\frac{\mathbf{x}}{\mathbf{y}}\|_{\dot{B}_2^{1,1}} + \|\ln \mathbf{y}\|_{\dot{B}_2^{1,1}}$. Moreover, the bilinear inequalities enunciated in theorem 3.10 hold for ψ_ν in place of $[\psi_1 A(\nabla^{-1} \psi_2)]$ there.

Proof We define

$$\begin{aligned} c_k := & \sum_{k_1 \in \mathbf{Z}} 2^{-\sigma|k-k_1|} [\|P_{k_1} N(\dots)\|_{L_t^1 L_x^2(\mathbf{R}^{1+2})} + \|P_{k_1} N(\dots)\|_{L_t^2 \dot{H}^{-\frac{1}{2}}}] \\ & + \sum_{k_1 \in \mathbf{Z}} 2^{-\sigma|k-k_1|} \|\frac{\partial_\nu \mathbf{y}}{\mathbf{y}}(0)\|_{L_x^2} + \sup_{\nu=0,1,2} \sum_{k_1 \in \mathbf{Z}} 2^{-\sigma|k-k_1|} \|P_{k_1} \partial_\nu(\frac{\mathbf{x}}{\mathbf{y}})(0)\|_{L_x^2}, \end{aligned}$$

where $N(\dots) = N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})$ runs over the nonlinearities in (1), (2). That this is indeed a frequency envelope with the desired properties follows from Corollary 2.6 as well as lemma 2.7. We need to exercise some care to get good enough control over the elliptic portions of ψ_ν . For this, truncate $N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})$ past some time $T_0 \gg \max\{2^{-k}, T\}$, and (committing abuse of notation) decompose the nonlinearity

$$P_k N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y}) = P_k Q_{<k} N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y}) + P_k Q_{\geq k} N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y}).$$

Then consider

$$\square^{-1} P_k Q_{\geq k} N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y}),$$

where the operator \square^{-1} is division by the symbol $(\tau^2 - |\xi|^2)$ on the space-time Fourier side. Clearly, from definition we have

$$\|\square^{-1} P_k Q_{\geq k} \nabla_{x,t} N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})\|_{L_t^1 L_x^2} \lesssim 2^{-k} c_k$$

Thus there exists a time $t_0 < T_0$ with the property

$$\|\square^{-1} P_k Q_{\geq k} \nabla_{x,t} N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})(t_0)\|_{L_x^2} \lesssim c_k$$

We easily check that (for C independent of k)

$$\|\square^{-1} R_0 P_k Q_{\geq k} \nabla_{x,t} N(\nabla \mathbf{x}, \nabla \mathbf{y}, \mathbf{x}, \mathbf{y})\|_{L_t^\infty L_x^2} \leq C,$$

while also (using the wave equation)

$$\|R_0 \nabla_{x,t} P_k(\frac{\mathbf{x}}{\mathbf{y}})\|_{L_t^\infty L_x^2} + \|R_0 \nabla_{x,t} P_k \ln \mathbf{y}\|_{L_t^\infty L_x^2} \leq C$$

Now construct a free wave a with the properties

$$a(t_0) = P_k(\frac{\mathbf{x}}{\mathbf{y}})(t_0) - \square^{-1}P_kQ_{\geq k}N(\nabla\mathbf{x}, \nabla\mathbf{y}, \mathbf{x}, \mathbf{y})(t_0)$$

$$\partial_t a(t_0) = P_k\partial_t(\frac{\mathbf{x}}{\mathbf{y}})(t_0) - \square^{-1}\partial_t P_kQ_{\geq k}N(\nabla\mathbf{x}, \nabla\mathbf{y}, \mathbf{x}, \mathbf{y})(t_0)$$

and similarly for $\ln \mathbf{y}$. It follows that the quantity $\frac{\mathbf{x}}{\mathbf{y}} - a - \square^{-1}P_kQ_{\geq k}N(\dots)$ satisfies

$$\square(\frac{\mathbf{x}}{\mathbf{y}} - a - \square^{-1}P_kQ_{\geq k}N(\dots)) = P_kQ_{<k}N(\nabla\mathbf{x}, \nabla\mathbf{y}, \mathbf{x}, \mathbf{y}),$$

$$\|R_0\nabla_{x,t}P_k(\frac{\mathbf{x}}{\mathbf{y}} - a - \square^{-1}P_kQ_{\geq k}N(\dots))\|_{L_t^\infty L_x^2} \leq C$$

as well as $\|\nabla_{x,t}(\frac{\mathbf{x}}{\mathbf{y}} - a - \square^{-1}P_kQ_{\geq k}N(\dots))(t_0)\|_{L_x^2} \lesssim c_k$. One also verifies that

$$\|\square^{-1}\nabla_{x,t}P_kQ_{\geq k}N(\nabla\mathbf{x}, \nabla\mathbf{y}, \mathbf{x}, \mathbf{y})\|_{\dot{X}_k^{-\frac{1}{2},1,2}} \lesssim c_k,$$

which by Sobolev's inequality also implies control over $\|\cdot\|_L$ as well as the Strichartz type norms $\|\cdot\|_{L_t^p L_x^q}$, $\frac{1}{p} + \frac{1}{q} < \frac{1}{2} - \delta_0$, of this expression. Now one solves the wave equation for $\frac{\mathbf{x}}{\mathbf{y}} - a - \square^{-1}P_kQ_{\geq k}N(\dots)$ with initial data given at time t_0 . Using 3.4(d) (which in turn relies on a truncated Duhamel's formula, see [23]), one constructs Schwartz extensions $\widetilde{P_k(\frac{\mathbf{x}}{\mathbf{y}})}$, $\widetilde{P_k \ln \mathbf{y}}$ of $P_k(\frac{\mathbf{x}}{\mathbf{y}})|_{[-T,T]}$, $P_k \ln \mathbf{y}|_{[-T,T]}$, respectively, with the properties

$$\|\nabla_{x,t}\widetilde{P_k(\frac{\mathbf{x}}{\mathbf{y}})}\|_{A[k]} \lesssim c_k, \|\nabla_{x,t}\widetilde{P_k \ln \mathbf{y}}\|_{A[k]} \lesssim c_k$$

as well as

$$\|R_0\nabla_{x,t}[\widetilde{P_k(\frac{\mathbf{x}}{\mathbf{y}})}]\|_{L_t^\infty L_x^2} \leq C, \|R_0\nabla_{x,t}\widetilde{P_k \ln \mathbf{y}}\|_{L_t^\infty L_x^2} \leq C$$

where C is independent of k . Using a partition of unity, one glues these extensions together to get Schwartz extensions $\widetilde{\frac{\mathbf{x}}{\mathbf{y}}}$, $\widetilde{\ln \mathbf{y}}$ of $\frac{\mathbf{x}}{\mathbf{y}}|_{[-T,T]}$, $\ln \mathbf{y}|_{[-T,T]}$ which satisfy

$$\|P_k(\frac{\widetilde{\mathbf{x}}}{\mathbf{y}})\|_{A[k]} \lesssim c_k$$

etc. Now one recalls that

$$\psi_\nu = (\frac{\partial_\nu \mathbf{x}}{\mathbf{y}} + i\frac{\partial_\nu \mathbf{y}}{\mathbf{y}})e^{i\sum_{j=1,2}\Delta^{-1}\partial_j(\frac{\partial_j \mathbf{x}}{\mathbf{y}})},$$

plugs in the Schwartz extensions of $\frac{\mathbf{x}}{\mathbf{y}}$ etc. and uses theorem 3.10 to obtain the desired conclusion. \blacksquare

Continuing with the proof of Proposition 3.12, our strategy now will be to analyze the wave equation satisfied by $\delta\psi_\nu = \tilde{\psi}_\nu - \psi_\nu$. Using (8) for both ψ_ν , $\tilde{\psi}_\nu$, and subtracting, we obtain a first version. We eliminate $\tilde{\psi}_\nu$ by substituting $\delta\psi_\nu + \psi_\nu$. One thereby obtains a sum of products of components $\delta\psi_\nu$, ψ_ν which are at least linear in $\delta\psi_\nu$. Proceeding as in the previous description, we decompose the $\delta\psi_\nu$, ψ_ν into gradient and elliptic parts. For the $\delta\psi_\nu$, this is obtained by applying the procedure to $\tilde{\psi}_\nu$, ψ_ν and forming the difference, resulting in

$$\delta\psi_\nu = R_\nu(\delta\psi) + \delta\chi_\nu, \delta\psi = -\sum_{i=1,2} R_i\psi_i$$

$$\begin{aligned} \delta\chi_\nu &= i \sum_{i,j=1}^2 \partial_i \Delta^{-1} (\tilde{\psi}_\nu \Delta^{-1} \partial_j (\tilde{\psi}_i^1 \tilde{\psi}_j^2 - \tilde{\psi}_j^1 \tilde{\psi}_i^2) - \tilde{\psi}_i \Delta^{-1} \partial_j (\tilde{\psi}_\nu^1 \tilde{\psi}_j^2 - \tilde{\psi}_j^1 \tilde{\psi}_\nu^2)) \\ &\quad - i \sum_{i,j=1}^2 \partial_i \Delta^{-1} (\psi_\nu \Delta^{-1} \partial_j (\psi_i^1 \psi_j^2 - \psi_j^1 \psi_i^2) - \psi_i \Delta^{-1} \partial_j (\psi_\nu^1 \psi_j^2 - \psi_j^1 \psi_\nu^2)). \end{aligned}$$

Clearly one can reexpress the latter difference as a sum of terms linear, quadratic and cubic in the $\delta\psi_\nu$, eliminating the $\tilde{\psi}_\nu = \delta\psi_\nu + \psi_\nu$. In order to demonstrate Proposition 3.12, we shall rely on the following refined

Proposition 3.17. *Let*

$$\square\delta\psi_\alpha = N_\alpha(\delta\psi_\nu, \psi_\nu)$$

on $[-T, T]$. Proceeding as above, express the nonlinearity as a sum of trilinear null-forms (substituting the gradient components for $\delta\psi_\nu, \psi_\nu$), as well as error terms 'at least quintilinear' in $\delta\psi_\nu, \psi_\nu$ (which arise upon substituting $\delta\chi_\nu, \chi_\nu$). Denote the sum of terms which are linear in $\delta\psi_\nu$ by $N_{1\alpha}(\delta\psi_\nu, \psi_\nu)$. Then for any $\zeta > 0$ there exists $T_0 > 0$, such that for any fixed smooth function $\chi(t) \in C^\infty(\mathbf{R})$ with $\text{supp}(\chi) \subset [-1, 1]^c$, $\chi|_{[-2, 2]^c} = 1$, we have

$$\begin{aligned} \|P_k(\delta\psi_\nu)\|_{S[k]([-T, T] \times \mathbf{R}^2)} &\leq K\tilde{c}_k, \quad T > \tilde{T}_0 \geq T_0 \\ \Rightarrow \|\chi(\frac{t}{\tilde{T}_0})P_k N_{1\alpha}(\delta\psi_\nu, \psi_\nu)\|_{N[k]([-T, T] \times \mathbf{R}^2)} &\lesssim \zeta K\tilde{c}_k \end{aligned}$$

Here $\{\tilde{c}_k\}$ is associated with \tilde{T}_0 as in Proposition 3.12 (substitute \tilde{T}_0 for T_1). Moreover, denoting the terms at least quadratic in $\delta\psi_\nu$ by $N_2(\delta\psi_\nu, \psi_\nu)$, and letting ϵ, \tilde{c}_k be as in the statement of Proposition 3.12, the following conclusion holds provided ϵ is small enough and $\{\tilde{c}_k\}$ 'flat enough':

$$\|P_k(\delta\psi_\nu)\|_{S[k]([-T, T] \times \mathbf{R}^2)} \leq K\tilde{c}_k \Rightarrow \|P_k N_{2\alpha}(\delta\psi_\nu, \psi_\nu)\|_{N[k]([-T, T] \times \mathbf{R}^{2+1})} \lesssim \epsilon K^2 \tilde{c}_k.$$

Deferring the proof of this for the moment, we continue with the proof of Proposition 3.12. Let $\zeta < 1/C$ for some $C \gg 1$, and construct T_0 as in Proposition 3.17; Define $T_1 := 2T_0$. Now assume we have the situation in the statement of Proposition 3.12. We intend to use the energy inequality 3.4(d). Fix $k \in \mathbf{Z}$, and consider $P_k \delta\psi_\nu$. We distinguish between the cases $T - T_1 < \frac{2^{-k}}{C}$ and the opposite. In the former case, the wave equation becomes useless, and we use the divergence-curl system directly: observe that by virtue of (6) we have for $i = 1, 2$, $T_1 < t \leq T$

$$P_k \delta\psi_i(t) - P_k \delta\psi_i(T_1) = \int_{T_1}^t \partial_i P_k \delta\psi_i dt + \int_{T_1}^t P_k [N(\psi, \delta\psi)] dt$$

In this equation, by abuse of notation, $N(\psi, \delta\psi)$ is a linear combination of terms of the schematic form $\delta\psi \nabla^{-1}(\psi^2)$, $\delta\psi \nabla^{-1}(\psi \delta\psi)$ etc. Let's put $N(\psi, \delta\psi) = \delta\psi \nabla^{-1}(\psi^2)$, the other terms being treated along the same lines (but also requiring ϵ to be small enough). We note that

$$\|P_k[\delta\psi \nabla^{-1}(\psi^2)]\|_{L_t^B L_x^2} \lesssim K 2^{(1-\frac{1}{B})k} \tilde{c}_k,$$

where B is an arbitrarily large number (the implied constants will depend on it). This follows from a simple frequency trichotomy and the bootstrap assumption.

Now using Holder's inequality, we deduce

$$\| \int_{T_1}^t P_k[N(\psi, \delta\psi)]dt \|_{L_x^2} \lesssim |t - T_1|^{1-\frac{1}{B}} 2^{(1-\frac{1}{B})k} K \tilde{c}_k \lesssim \frac{1}{C^{1-\frac{1}{B}}} K \tilde{c}_k$$

Clearly we also have

$$\| \int_{T_1}^t \partial_i P_k \delta\psi_i dt \|_{L_x^2} \lesssim |t - T_1| 2^k K \tilde{c}_k \lesssim \frac{1}{C} K \tilde{c}_k.$$

Therefore, we infer that

$$\| P_k \delta\psi_i \|_{L_t^\infty L_x^2} \lesssim (1 + \frac{K}{C^{1-\frac{1}{B}}}) \tilde{c}_k \lesssim \frac{K}{100} \tilde{c}_k$$

provided K, C are chosen¹⁶ large enough (in relation to B). Arguing similarly, one deduces as well that

$$2^{\frac{k}{2}} \| P_k \delta\psi_i \|_{L_t^2 L_x^2} + 2^{-\frac{k}{2}} \| P_k \partial_t \delta\psi_i \|_{L_t^2 L_x^2} \lesssim \frac{K}{C^{\frac{1}{2}}} \tilde{c}_k.$$

Using the fact that (see e. g. [23])

$$\begin{aligned} \| P_k \psi \|_{S[k]([-T, T] \times \mathbf{R}^2)} &\lesssim \| P_k \psi \|_{L_t^\infty L_x^2([-T, T] \times \mathbf{R}^2)} + 2^{\frac{k}{2}} \| P_k \psi \|_{L_t^2 L_x^2([-T, T] \times \mathbf{R}^2)} \\ &\quad + 2^{-\frac{k}{2}} \| P_k \partial_t \psi \|_{L_t^2 L_x^2([-T, T] \times \mathbf{R}^2)} \end{aligned}$$

and choosing K, C large enough, one deduces from this that

$$\| P_k \delta\psi_i \|_{S[k]([T_1, T] \times \mathbf{R}^2)} \lesssim \frac{K}{2} \tilde{c}_k,$$

which is the desired conclusion for $P_k \delta\psi_i$. The argument for $\delta\psi_0$ is similar using (7). Thus we see that we may assume $|T - T_1| \geq \frac{2^{-k}}{C}$. Moreover, reiterating the preceding argument, and choosing K large enough, we conclude that $\| P_k \delta\psi_\nu \|_{L_t^\infty L_x^2([T_1, T_1 + \frac{2^{-k}}{C}] \times \mathbf{R}^2)} \leq \frac{K}{100C} \tilde{c}_k$. Now revert to the old notation

$$\square \delta\psi_\nu = N(\delta\psi_\nu, \psi_\nu) = N_1(\delta\psi_\nu) + N_2(\delta\psi_\nu, \psi_\nu)$$

as in Proposition 3.17. Clearly, we have

$$\| P_k N_1(\delta\psi_\nu, \psi_\nu) \|_{N[k]([T_1, T] \times \mathbf{R}^2)} \leq \| P_k \chi(\frac{t}{T_0}) N_1(\delta\psi_\nu, \psi_\nu) \|_{N[k]([T_1, T] \times \mathbf{R}^2)}$$

$$\| P_k N_2(\delta\psi_\nu, \psi_\nu) \|_{N[k]([T_1, T] \times \mathbf{R}^2)} \lesssim \| P_k N_2(\delta\psi_\nu, \psi_\nu) \|_{N[k]([-T, T] \times \mathbf{R}^2)}$$

Using 3.4(d) as well as time translation invariance, we can now infer that

$$\| P_k \delta\psi_\nu \|_{S[k]([T_1, T] \times \mathbf{R}^2)} \lesssim \frac{K}{100} \tilde{c}_k + \zeta K \tilde{c}_k + \epsilon^2 K^2 \tilde{c}_k^2 \lesssim \frac{K}{2} \tilde{c}_k$$

provided ϵ, ζ are small enough. This yields the desired conclusion.

3.18. The proof of theorem 3.3. This is basically identical to the proof of theorem 3.2. Control over some subcritical norm $\|u\|_{L_t^\infty H^{1+\epsilon}}$ follows from standard Moser estimates instead of Corollary 2.6.

¹⁶Of course C is chosen independently of K .

4. THE PROOF OF PROPOSITION 3.17.

We have thus reduced the proof of theorem 3.2 to the verification of Proposition 3.17 in addition to the technical Moser type estimates allowing estimation of $\|P_k \psi\|_{S[k]}$. The proof of this Proposition is divided into the part dealing with expressions linear in $\delta\psi_\nu$, as well as those of higher degree of linearity. We commence by spelling out in detail the decomposition $N_\alpha(\delta\psi_\nu, \psi_\nu) = N_{1\alpha}(\delta\psi_\nu, \psi_\nu) + N_{2\alpha}(\delta\psi_\nu, \psi_\nu)$, where $\square\delta\psi_\alpha = N_\alpha(\delta\psi_\nu, \psi_\nu)$. As in [23], this decomposition requires extreme care in order to avoid too many time derivatives. Recalling (8), we define

$$N_1(\delta\psi_\nu, \psi_\nu) := \sum_{i=1}^3 A_{i\alpha}(\delta\psi_\nu, \psi_\nu) + \sum_{i=1}^5 B_{i\alpha}(\delta\psi_\nu, \psi_\nu) + \sum_{i=1}^5 C_{i\alpha}$$

where

$$\begin{aligned} A_1(\delta\psi_\nu, \psi_\nu) = & i\partial^\beta[\delta\psi_\alpha\Delta^{-1}\sum_{j=1}^2\partial_j[R_\beta\psi^1R_j\psi^2 - R_\beta\psi^2R_j\psi^1]] \\ & - i\partial^\beta[\delta\psi_\beta\Delta^{-1}\sum_{j=1}^2\partial_j[R_\alpha\psi^1R_j\psi^2 - R_\alpha\psi^2R_j\psi^1]] \\ & + i\partial_\alpha[\delta\psi_\nu\Delta^{-1}\sum_{j=1}^2\partial_j[R^\nu\psi^1R_j\psi^2 - R^\nu\psi^2R_j\psi^1]] \end{aligned} \quad (19)$$

$$\begin{aligned} A_1(\delta\psi_\nu, \psi_\nu) = & + i\partial^\beta[\psi_\alpha\Delta^{-1}\sum_{j=1}^2\partial_j[R_\beta\delta\psi^1R_j\psi^2 - R_\beta\psi^2R_j\delta\psi^1]] \\ & - i\partial^\beta[\psi_\beta\Delta^{-1}\sum_{j=1}^2\partial_j[R_\alpha\delta\psi^1R_j\psi^2 - R_\alpha\psi^2R_j\delta\psi^1]] \\ & + i\partial_\alpha[\psi_\nu\Delta^{-1}\sum_{j=1}^2\partial_j[R^\nu\delta\psi^1R_j\psi^2 - R^\nu\psi^2R_j\delta\psi^1]] \end{aligned}$$

$$\begin{aligned} A_3(\delta\psi_\nu, \psi_\nu) = & + i\partial^\beta[\psi_\alpha\Delta^{-1}\sum_{j=1}^2\partial_j[R_\beta\psi^1R_j\delta\psi^2 - R_\beta\delta\psi^2R_j\psi^1]] \\ & - i\partial^\beta[\psi_\beta\Delta^{-1}\sum_{j=1}^2\partial_j[R_\alpha\psi^1R_j\delta\psi^2 - R_\alpha\delta\psi^2R_j\psi^1]] \\ & + i\partial_\alpha[\psi_\nu\Delta^{-1}\sum_{j=1}^2\partial_j[R^\nu\psi^1R_j\delta\psi^2 - R^\nu\delta\psi^2R_j\psi^1]]. \end{aligned}$$

$$B_1, C_1(\delta\psi_\nu, \psi_\nu) = \nabla_{x,t}[\delta\psi\nabla^{-1}[\psi\nabla^{-1}[\psi\nabla^{-1}(\psi^2)]]]$$

$$B_2, C_2(\delta\psi_\nu, \psi_\nu) = \nabla_{x,t}[\psi\nabla^{-1}[\delta\psi\nabla^{-1}[\psi\nabla^{-1}(\psi^2)]]]$$

$$B_3, C_3(\delta\psi_\nu, \psi_\nu) = \nabla_{x,t}[\psi\nabla^{-1}[\psi\nabla^{-1}[\delta\psi\nabla^{-1}(\psi^2)]]]$$

$$B_{4,5}, C_{4,5}(\delta\psi_\nu, \psi_\nu) = \nabla_{x,t}[\psi \nabla^{-1}[\psi \nabla^{-1}[\psi \nabla^{-1}(\delta\psi\psi)]]$$

Of course we have used schematic notation for the B, C 's, as their fine structure won't matter. They are obtained by substituting one χ_ν instead of the corresponding entry ψ_ν in the inner square bracket expressions on the right hand side of (8), where χ_ν is the 'elliptic component' of the spherically symmetric ψ_ν in the decomposition $\psi_\nu = R_\nu\psi + \chi_\nu$. We recall $\psi_\nu = \psi_\nu^1 + i\psi_\nu^2$, $\delta\psi_\nu = \delta\psi_\nu^1 + i\delta\psi_\nu^2$, $\psi = -\sum_{j=1,2} R_j\psi_j$, $\delta\psi = -\sum_{j=1,2} R_j\delta\psi_j$. One can then define $N_{2\alpha}(\delta\psi_\nu, \psi_\nu) = N_\alpha(\delta\psi_\nu, \psi_\nu) - N_{1\alpha}(\delta\psi_\nu, \psi_\nu)$. The quintilinear terms above shall be relatively simple to estimate on account of the strong Strichartz type estimates satisfied by the ψ_ν , see theorem 3.7 as well as the definition of $\|\cdot\|_{S[k]}$. Unfortunately, the latter norm falls short of controlling $\|\cdot\|_{L_t^2 L_x^\infty}$, which appears necessary in order to grant an elementary estimation of the trilinear terms $A_{i\alpha}$. We shall instead have to revert to the inherent null-structure in these terms as was done already in [23], in addition to the more complicated ingredients in $\|\cdot\|_{\mathcal{S}}$. The main new difficulty over the estimates in [23] has to do with the fact that we need to gain explicitly in time in these estimates. This would be relatively straightforward if we were working with Lebesgue type spaces; however, we shall work with null-frame spaces of type $L_{t\omega}^2 L_{x\omega}^\infty$, which considerably complicates obtaining gains in time. The main novelty here (lemma 4.4) shall be a special type of decomposition of the spherical components ψ_ν into pieces which have well-defined physical as well as frequency localization properties. More precisely, we shall be able to physically localize ψ_ν closely to the light cone. This part will then be written as a sum of two components, the first of which can be written as a sum of pieces which propagate in a direction essentially opposite to their physical support. Thus the first component is obtained by first localizing ψ_ν to an angular sector in Fourier space, then multiplying with a physical cutoff localizing to an *opposite or identical*¹⁷ angular sector, and finally summing over all sectors. The size of the angular sectors shall essentially be dictated by the ζ in the statement of Proposition 3.17. While the first component is exactly the part which fails to decay in $L_{t\omega}^2 L_{x\omega}^\infty$ as $t \rightarrow \infty$, it does lead to improved trilinear null-form estimates due to the dual localization properties. The 2nd component in turn will decay like a standard Lebesgue norm as $t \rightarrow \infty$. The next subsection contains the core estimates. As the estimates are rather technical, we briefly explain the strategy of the proof, which is conceptually simple:

(1): First, upon localizing the nonlinearity to a time interval $t \sim 2^i$, one tries to reduce the frequencies of all functions occurring inside the nonlinearity to absolute size $<< 2^{\delta i}$, for some small $\delta > 0$. The idea here is that far apart frequencies should interact little. But this in addition to the refined control over the frequency modes of the spherically symmetric components should suffice to get control over the cases when extremely small or large frequencies are present. The tool to achieve this are the refined trilinear estimates in 3.4(c). Unfortunately, these estimates aren't quite good enough to control certain high-high interactions, which accounts for a number of extra cases that need to be considered.

(2) Having controlled the cases when the frequencies are very small or large in relation to the time interval one works on, one now tries to exploit the pointwise

¹⁷This depends on whether the space-time Fourier support is contained in the upper half-space $\tau > 0$ or lower half-space $\tau < 0$.

estimates provided by Christodoulou-Tahvildar-Zadeh, since one has gained some room to lose in the frequencies. The device here is the decomposition of the spherically symmetric components referred to in the preceding paragraph, which is a direct consequence of the pointwise decay estimates. This allows one to decompose these components into pieces that disperse quickly enough, as well as other pieces that interact very weakly. Of course one exploits the trilinear structure of the nonlinearity to make this work.

4.1. Estimating the trilinear null-forms. We use the operator $I = \sum_{k \in \mathbf{Z}} P_k Q_{<k+100}$ as before and employ the schematic decomposition

$$\nabla_{x,t}[\psi_1 \nabla^{-1}[\psi_2 \psi_3]] = \nabla_{x,t}[\psi_1 I \nabla^{-1}[\psi_2 \psi_3]] + \nabla_{x,t}[\psi_1 (1 - I) \nabla^{-1}[\psi_2 \psi_3]]$$

for each of the $A_{i\alpha}$'s. In order to make sense of this, one needs to substitute Schwartz extensions for the inputs $\delta\psi^{1,2}|_{[T_1, T]}$, $\psi^{1,2}|_{[T_1, T]}$ of the inner square brackets, in accordance with the bootstrap assumption in Proposition 3.17. In the following we shall localize the frequency localized nonlinearities $P_k N(\dots)$ to a dyadic time interval $t \sim 2^i$ and strive for an estimate of the form

$$\|\chi_i(t) P_k N(\dots)\|_{N[k]} \lesssim 2^{-\mu i} \tilde{c}_k$$

One can then sum over i large enough to obtain the estimate in Proposition 3.12.

(A): The large modulation case. Estimating the terms

$$(I) : \nabla_{x,t}[\delta\psi \Delta^{-1} \sum_{j=1,2} \partial_j (1 - I)[R_\beta \psi_2 R_j \psi_3 - R_j \psi_2 R_\beta \psi_3]]$$

$$(II) : \nabla_{x,t}[\psi \Delta^{-1} \sum_{j=1,2} \partial_j (1 - I)[R_\beta \delta\psi_2 R_j \psi_3 - R_j \delta\psi_2 R_\beta \psi_3]]$$

(I): *The first term.* We use the decomposition

$$\begin{aligned} & \chi\left(\frac{t}{\tilde{T}_0}\right) \nabla_{x,t}[\delta\psi \Delta^{-1} \sum_{j=1,2} \partial_j (1 - I)[R_\beta \psi_2 R_j \psi_3 - R_j \psi_2 R_\beta \psi_3]] \\ &= \sum_{i \geq \log_2 \tilde{T}_0} \chi_i(t) \nabla_{x,t}[\delta\psi \Delta^{-1} \sum_{j=1,2} \partial_j (1 - I)[R_\beta \psi_2 R_j \psi_3 - R_j \psi_2 R_\beta \psi_3]], \end{aligned}$$

where $\chi_i(t)$ smoothly localizes to the interval $t \sim 2^i$. Then we localize the frequencies and freeze $i \in \mathbf{Z}$, arriving at an expression

$$\chi_i(t) \nabla_{x,t} P_{k_0} [P_{k_1} \delta\psi \Delta^{-1} \sum_{j=1,2} \partial_j (1 - I) P_k [R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]]$$

We distinguish between the following cases:

(I.a): *One of the following options hold:* $i \lesssim |k_2|$, $i \lesssim |k_3|$, $i \lesssim |k_0 - k_1|$, $i \lesssim \min\{|k - k_1|, |k - k_2|\}$. This case is handled by means of lemma 3.16 as well as the following lemma, provided $P_{k_{2,3}} \psi_{2,3}$ are of the first type:

Lemma 4.2. [23] *Let $\psi_{1,2,3} \in \mathbf{R}^{2+1}$. Then, for integers $k_{1,2,3}$ and suitable $\delta_{1,2} > 0$, the following inequality holds:*

$$\begin{aligned} & \|\nabla_{x,t} P_0 [P_{k_1} \psi_1 \nabla^{-1} (1 - I) [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[0]} \\ & \lesssim 2^{\delta_1 \min\{-\min\{k_1, k_2, k_3\}, 0\}} \prod_i 2^{\delta_2 \min\{\max_{j \neq i}\{k_i, k_i - k_j\}, 0\}} \prod_l \|P_{k_l} \psi_l\|_{S[k_l]} \end{aligned}$$

Indeed, observe that if $i \lesssim \max\{|k_2|, |k_3|\}$, we obtain from lemma 3.16 that $\min\{\|P_{k_2} \psi_2\|_{S[k_2]}, \|P_{k_3} \psi_3\|_{S[k_3]}\} \lesssim 2^{-\mu^i}$. Carrying out the summations over k, k_i satisfying these assumptions, we arrive at the upper bound $\lesssim 2^{-\mu^i} K \tilde{c}_{k_0}$. Summing over $i \geq \log_2 \tilde{T}_0$ results in a small exponential gain in $T_0 \leq \tilde{T}_0$. If one of the other cases occurs, one gets an exponential gain $2^{-\min\{\delta_1, \delta_2\}i}$ from the above lemma. We are fudging a bit since we have thrown the localizer $\chi_i(t)$ in front, and this may affect the space-time Fourier support of the expression, hence its norm $\|\cdot\|_{N[k_0]}$. However, this is detrimental only if the modulation (i. e. distance of the space-time Fourier support to the light cone) is $\lesssim 2^{-i}$, and only affects those parts estimated with respect to $\|\cdot\|_{\dot{X}_{k_0}^{-1, -\frac{1}{2}, 1}}$, as null-frame spaces aren't needed yet, see the proof in [23].

Assuming $Q_{<-i+O(1)}(\text{Output})$ to be a $\dot{X}_{k_0}^{-1, -\frac{1}{2}, 1}$ -atom, we estimate

$$\begin{aligned} & \|\chi_i(t) P_{k_0} Q_{<-i+O(1)}(\text{Output})\|_{L_t^1 \dot{H}^{-1}} \lesssim \|\chi_i(t)\|_{L_t^2 L_x^\infty} \|P_{k_0} Q_{<-i+O(1)}(\text{Output})\|_{L_t^2 \dot{H}^{-1}} \\ & \lesssim \sum_{a < 0} 2^{\frac{a}{2}} 2^{\frac{-i+a}{2}} \lesssim 1 \end{aligned}$$

Thus the cutoff is irrelevant.

Now assume at least one of $P_{k_{2,3}}$ is of the 2nd type. We need the following lemma

Lemma 4.3. *Let $\psi_{2,3} \in \mathcal{S}(\mathbf{R}^{2+1})$. Assume also that $\|P_{k_{2,3}} \psi_{2,3}\|_{S[k_{2,3}]} \lesssim \frac{\tilde{c}_{k_{2,3}}}{\epsilon}$ with a frequency envelope \tilde{c}_k as in the preceding. Assume that $P_{k_2} \psi_2$ is of the 2nd type, and $P_{k_3} \psi_3$ admits a decomposition into functions of first and 2nd type as enunciated in theorem 3.10. Then we have for suitable $\delta_{1,2} > 0$*

$$\begin{aligned} & \|P_k [R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]\|_{L_t^2 L_x^2} \\ & \lesssim 2^{\frac{\min\{k_3, k\}}{2}} 2^{\delta_1 \min\{k_2 - k_3, 0\}} 2^{\delta_2 \min\{k - k_2\}} \left[\frac{\tilde{c}_{k_2}}{\epsilon} + \frac{\tilde{c}_{k_3}}{\epsilon} \right] \end{aligned}$$

Proof First assume that $P_{k_3} \psi_3$ is of the first type. Using the definition of $\|\cdot\|_{S[k]}$, we infer the desired estimate for the contributions of

$$P_k [R_\beta P_{k_2} (1 - I) \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} (1 - I) \psi_2 R_\beta P_{k_3} \psi_3]$$

and similarly for

$$P_k [R_\beta P_{k_2} \psi_2 R_j P_{k_3} (1 - I) \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} (1 - I) \psi_3]$$

Take the first expression: first consider the case $k_2 = k_3 + O(1)$. We estimate, using theorem 3.10

$$\begin{aligned} & \|P_k [R_\beta P_{k_2} (1 - I) \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} (1 - I) \psi_2 R_\beta P_{k_3} \psi_3]\|_{L_t^2 L_x^2} \\ & \lesssim 2^{2k(\frac{1}{1+\mu} - \frac{1}{2})} \|P_k [R_\beta P_{k_2} (1 - I) \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} (1 - I) \psi_2 R_\beta P_{k_3} \psi_3]\|_{L_t^2 L_x^{1+\mu}} \\ & \lesssim 2^{\frac{k}{2}} 2^{\frac{(k-k_2)(1-3\mu)}{2+2\mu}} \frac{\tilde{c}_{k_3}}{\epsilon} \end{aligned}$$

Next, in case $k_2 \ll k_3$, we estimate

$$\begin{aligned} & \|P_k[R_\beta P_{k_2}(1-I)\psi_2 R_j P_{k_3}\psi_3]\|_{L_t^2 L_x^2} \\ & \lesssim \left(\sum_{c \in C_{k_3, k_2-k_3}} \|P_k[R_\beta P_{k_2}(1-I)\psi_2 R_j P_c \psi_3]\|_{L_t^2 L_x^2}^2 \right)^{\frac{1}{2}} \lesssim 2^{\frac{k_3}{2}} 2^{\delta_1(k_2-k_3)} \frac{\tilde{c}_{k_3}}{\epsilon} \end{aligned}$$

The remaining term is estimated similarly, as is the case when $k_2 \gg k_3$. Further, if for example $k_2 = k_3 + O(1)$, we can estimate

$$\begin{aligned} & \|P_k[R_\beta P_{k_2} I \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} I \psi_2 R_\beta P_{k_3} \psi_3]\|_{L_t^2 L_x^2} \\ & \lesssim 2^{(1-\epsilon)k} \|P_{k_2} \psi_2\|_{L_t^2 L_x^{2+}} \|P_{k_3} \psi_3\|_{L_t^\infty L_x^2} \lesssim 2^{(1-\epsilon)k - (\frac{1}{2}-\epsilon)k_2} \frac{\tilde{c}_{k_2}}{\epsilon} \end{aligned}$$

The remaining cases $k = k_2 + O(1)$, $k = k_3 + O(1)$ are handled similarly. Now assume that both $P_{k_2} \psi_2$ and $P_{k_3} \psi_3$ are of 2nd type. In that case, if $k_1 = k_2 + O(1)$, estimate

$$\begin{aligned} & \|P_k[R_\beta P_{k_1} \psi_1 P_{k_2} R_j \psi_2]\|_{L_t^2 L_x^2} \lesssim 2^{(1-\epsilon)k} \|R_\beta P_{k_1} \psi_1\|_{L_t^\infty L_x^2} \|P_{k_2} \psi_2\|_{L_t^2 L_x^{2+}} \\ & \lesssim 2^{(1-\epsilon)(k-k_2)} 2^{\frac{k}{2}} \frac{\tilde{c}_{k_2}}{\epsilon} \end{aligned}$$

The remaining frequency interactions are treated similarly. ■

Returning to case (I.a) when at least one of $P_{k_2,3} \psi_{2,3}$ is of 2nd type, we claim that we have the estimate

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi_1 \nabla^{-1} (1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[k_0]} \\ & \lesssim 2^{\delta_1 \{\min_{i=2,3} \{k, k_i\} - \max_{i=2,3} \{k, k_i\}\}} 2^{-\delta_2(|k-k_1|)} \left[\frac{\tilde{c}_{k_2}}{\epsilon} + \frac{\tilde{c}_{k_3}}{\epsilon} \right] \tilde{c}_{k_0} \end{aligned}$$

One could then sum over all frequency parameters (except k_0) and obtain the required exponential gain in i under the hypotheses of case (I.a)¹⁸. To verify this estimate, we may assume $k_0 = 0$. One needs to distinguish between $k_1 \in [-10, 10]$, $k_1 > 10$, $k_1 < -10$. These are similar, so we treat the first case: we have

$$\begin{aligned} & \|\nabla_{x,t} P_0 Q_{>10} [P_{k_1} \delta \psi_1 \nabla^{-1} (1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[0]} \\ & \lesssim \|P_0 Q_{>10} [P_{k_1} \delta \psi_1 \nabla^{-1} (1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{L_t^2 L_x^2} \\ & \lesssim \|P_{k_1} \delta \psi_1\|_{L_t^\infty L_x^2} 2^{\frac{\min\{k, k_3\}}{2}} 2^{\delta_1 \min\{k_2-k_3, 0\}} 2^{\delta_2 \min\{k-k_2, 0\}} \left[\frac{\tilde{c}_{k_2}}{\epsilon} + \frac{\tilde{c}_{k_3}}{\epsilon} \right] \tilde{c}_0 \end{aligned}$$

One checks that this verifies the claim, with a lot to spare. Next, we can estimate

$$\begin{aligned} & \|\nabla_{x,t} P_0 Q_{<10} [P_{k_1} Q_{<k-100} \delta \psi_1 \\ & \quad \nabla^{-1} (1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[0]} \\ & \lesssim \|\nabla_{x,t} P_0 Q_{<10} [P_{k_1} Q_{<k-100} \delta \psi_1 \\ & \quad \nabla^{-1} (1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{\dot{X}_0^{-1, -\frac{1}{2}, 1}} \\ & \lesssim 2^{-\frac{k}{2}} 2^{\frac{\min\{k, k_3\}}{2}} 2^{\delta_1 \min\{k_2-k_3, 0\}} 2^{\delta_2 \min\{k-k_2, 0\}} \tilde{c}_0 \left[\frac{\tilde{c}_{k_2}}{\epsilon} + \frac{\tilde{c}_{k_3}}{\epsilon} \right] \end{aligned}$$

¹⁸The cutoff $\chi_i(t)$ in front is handled as before.

Again this verifies the claim. Finally, we have the estimate

$$\begin{aligned}
& \|\nabla_{x,t} P_0 Q_{<10} [P_{k_1} Q_{\geq k-100} \delta \psi_1 \\
& \quad \nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[0]} \\
& \lesssim \|\nabla_{x,t} P_0 Q_{<10} [P_{k_1} Q_{\geq k-100} \delta \psi_1 \\
& \quad \nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \|P_{k_1} Q_{\geq k-100} \delta \psi_1\|_{L_t^2 L_x^2} \\
& \quad \|\nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]\|_{L_t^2 L_x^\infty} \\
& \lesssim 2^{-\frac{k}{2}} 2^{\frac{\min\{k, k_3\}}{2}} 2^{\delta_1 \min\{k_2 - k_3, 0\}} 2^{\delta_2 \min\{k - k_2, 0\}} \tilde{c}_0 \left[\frac{\tilde{c}_{k_2}}{\epsilon} + \frac{\tilde{c}_{k_3}}{\epsilon} \right]
\end{aligned}$$

as in the preceding estimate. This concludes case (I.a).

(I.b): $i \lesssim |k|$, $i \lesssim |k_1|$, and none of the properties in (I.a) hold. Thus in this case, we have $k \lesssim -i$, $k_1 \lesssim -i$, $|k - k_1| \ll i$; we may treat the last difference as $O(1)$ ¹⁹. In this case we have to work harder to obtain the exponential gain in i , since the previous trilinear estimates won't suffice. Observe that we only need to worry about the case $\nu = 0$, though, since otherwise one can pull a derivative out of the inner square bracket expression. Also, we may easily reduce the Fourier support of $P_{k_2,3} \psi_{2,3}$ to the hyperbolic regime²⁰ (distance to light cone at most comparable to frequency). First, consider the case $|k| > (1 + \mu)i$, for some small $\mu > 0$. In that case, we have

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} \chi_i(t) [P_{k_1} \delta \psi_1 \nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[k_0]} \\
& \lesssim \|\nabla_{x,t} P_{k_0} \chi_i(t) [P_{k_1} Q_{\geq k-100} \delta \psi_1 \\
& \quad \nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[k_0]} \\
& + \|\nabla_{x,t} P_{k_0} \chi_i(t) [P_{k_1} Q_{<k-100} \delta \psi_1 \\
& \quad \nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[k_0]}
\end{aligned}$$

The first summand is further decomposed as follows:

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} \chi_i(t) [P_{k_1} Q_{\geq k-100} \delta \psi_1 \\
& \quad \nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[k_0]} \\
& \lesssim \|\nabla_{x,t} P_{k_0} Q_{>k_0} \chi_i(t) [P_{k_1} Q_{\geq k-100} \delta \psi_1 \\
& \quad \nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\
& + \|\nabla_{x,t} P_{k_0} Q_{<k_0} \chi_i(t) [P_{k_1} Q_{\geq k-100} \delta \psi_1 \\
& \quad \nabla^{-1} (1 - I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}}
\end{aligned}$$

¹⁹We do this in order to avoid carrying too many small constants around; this is legitimate since the exponential gains obtained later are independent.

²⁰We shall not include the localizers $Q_{<k_2,3}$ everywhere in order to streamline notation.

We then estimate

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} Q_{>k_0} \chi_i(t) [P_{k_1} Q_{\geq k-100} \delta \psi_1 \\
& \quad \nabla^{-1}(1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\
& \lesssim 2^{\frac{k_0}{2}} \|\chi_i(t)\|_{L_t^2} \|P_{k_1} Q_{\geq k-100} \delta \psi_1\|_{L_t^\infty L_x^2} \\
& \quad \|\nabla^{-1}(1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]\|_{L_t^\infty L_x^2} \\
& \lesssim 2^{\frac{k_0}{2} + \frac{i}{2}} \tilde{c}_{k_1} \lesssim 2^{-\frac{\mu}{2} i} \tilde{c}_{k_1}
\end{aligned}$$

Similarly, we have

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} Q_{<k_0} \chi_i(t) [P_{k_1} Q_{\geq k-100} \delta \psi_1 \\
& \quad \nabla^{-1}(1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim 2^{k_0} \|\chi_i(t)\|_{L_t^2} \|P_{k_1} Q_{\geq k-100} \delta \psi_1\|_{L_t^2 L_x^2} \\
& \quad \|\nabla^{-1}(1-I) P_k [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]\|_{L_t^\infty L_x^2},
\end{aligned}$$

and this is controlled by $2^{k_0 - \frac{k}{2} + \frac{i}{2}} \tilde{c}_{k_1} \lesssim 2^{-\frac{\mu}{2} i} \tilde{c}_{k_1}$ as desired. The remaining terms are handled similarly. Thus we now assume that $i \lesssim |k| \leq (1+\mu)|i|$. We then claim that we may replace the operator $(1-I)$ by $Q_{>\frac{k}{2}}$. Indeed, we have

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} \chi_i(t) [P_{k_1} \delta \psi_1 \\
& \quad \nabla^{-1}(1-I) P_k Q_{<\frac{k}{2}} [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[k_0]} \\
& \lesssim \sum_{k+100 < a < \frac{k}{2}} \|\nabla_{x,t} P_{k_0} \chi_i(t) [P_{k_1} Q_{<a-100} \delta \psi_1 \\
& \quad \nabla^{-1}(1-I) P_k Q_a [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[k_0]} \\
& + \sum_{k+100 < a < \frac{k}{2}} \|\nabla_{x,t} P_{k_0} \chi_i(t) [P_{k_1} Q_{\geq a-100} \delta \psi_1 \\
& \quad \nabla^{-1}(1-I) P_k Q_a [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]\|_{N[k_0]}
\end{aligned}$$

We treat the first summand, the 2nd being similar. We have

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} [\chi_i(t) [P_{k_1} Q_{<a-100} \delta \psi_1 \\
& \quad \nabla^{-1}(1-I) P_k Q_a [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]]\|_{N[k_0]} \\
& \lesssim 2^{-\frac{k_0}{2}} \|P_{k_0} Q_{a+O(1)} [\chi_i(t) [P_{k_1} Q_{<a-100} \delta \psi_1 \\
& \quad \nabla^{-1}(1-I) P_k Q_a [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]]\|_{L_t^2 L_x^2} \\
& \lesssim 2^{\frac{a+k}{2}} \tilde{c}_{k_1}
\end{aligned}$$

This is clearly acceptable. We now notice the identity

$$-2(\partial_t \psi_2 \partial_r \psi_3 - \partial_r \psi_2 \partial_t \psi_3) = (\partial_t + \partial_r) \psi_2 (\partial_t - \partial_r) \psi_3 - (\partial_t - \partial_r) \psi_2 (\partial_t + \partial_r) \psi_3$$

Applying this to our frequency localized situation, we have the identity²¹(recall that $\psi_{1,2}$ are radial)

$$\begin{aligned} & R_0 P_{k_2} Q_{<k_2} \psi_2 R_i P_{k_3} Q_{<k_3} \psi_3 - R_i P_{k_2} Q_{<k_2} \psi_2 R_0 P_{k_3} Q_{<k_3} \psi_3 \\ &= \frac{x_i}{r} [(\partial_t + \partial_r) \nabla^{-1} P_{k_2} Q_{<k_2} \psi_2 (\partial_t - \partial_r) \nabla^{-1} P_{k_3} Q_{<k_3} \psi_3 \\ & \quad - (\partial_t - \partial_r) \nabla^{-1} P_{k_2} Q_{<k_2} \psi_2 (\partial_t + \partial_r) \nabla^{-1} P_{k_3} Q_{<k_3} \psi_3] \end{aligned}$$

Now let $\phi \in C_0^\infty(\mathbf{R})$ be a smooth cutoff and use the decomposition $\psi_{2,3} = \phi_{2^{-\frac{i}{2+}}}(u) \psi_{2,3} + (1 - \phi_{2^{-\frac{i}{2+}}}(u)) \psi_{2,3}$, where $u = t - r$ and $\phi_\lambda(u) = \phi(\frac{u}{\lambda})$. Now Proposition 2.2 implies that

$$\|\chi_i(t) \phi_{2^{-\frac{i}{2+}}}(u) \psi_{2,3}\|_{L_x^2} \lesssim \sqrt{t^{-1} \times t \times 2^{-\frac{i}{2+}}} \lesssim 2^{-\frac{i}{4+}}$$

Now let $\frac{2}{4+} + \frac{1}{M} = \frac{1}{2}$ and return to the full expression. We first estimate the large-modulation contribution:

$$\begin{aligned} & \|P_{k_0} Q_{>k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} \delta \psi_1 \\ & \quad \nabla^{-1} P_k Q_{>\frac{k}{2}} [R_0 P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2] R_i P_{k_3} Q_{<k_3} \psi_3 \\ & \quad \quad - R_i P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2] R_0 P_{k_3} Q_{<k_3} \psi_3] \|_{N[k_0]} \\ & \lesssim \|P_{k_0} Q_{>k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} \delta \psi_1 \\ & \quad \quad \nabla^{-1} P_k Q_{>\frac{k}{2}} [R_0 P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2] R_i P_{k_3} Q_{<k_3} \psi_3 \\ & \quad \quad - R_i P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2] R_0 P_{k_3} Q_{<k_3} \psi_3] \|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim 2^{-\frac{k_0}{2}} 2^{(1-\epsilon)k_0 - k_1} \|\chi_i(t)\|_{L_t^{4+}} \|P_{k_1} \delta \psi_1\|_{L_t^{4+} L_x^\infty} \\ & \quad \|R_0 P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2]\|_{L_t^\infty L_x^2} \|P_{k_3} \psi_3\|_{L_t^M L_x^{2+}} \\ & \lesssim 2^{(\frac{1}{2}-\epsilon)k_0 - \frac{k_1}{4+}} \tilde{c}_{k_1} \frac{\tilde{c}_{k_3}}{\epsilon} \end{aligned}$$

Keeping our assumptions on the frequencies in mind, this is more than what we need. Next, restricting the expression to modulation $\leq 2^{k_0}$, we have

$$\begin{aligned} & \|P_{k_0} Q_{\leq k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} \delta \psi_1 \\ & \quad \nabla^{-1} P_k Q_{\geq \frac{k}{2}} [R_0 P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2] R_i P_{k_3} Q_{<k_3} \psi_3 \\ & \quad \quad - R_i P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2] R_0 P_{k_3} Q_{<k_3} \psi_3] \|_{N[k_0]} \\ & \lesssim \sum_{k_2 + O(1) > a \geq \frac{k}{2}} \|P_{k_0} Q_{\leq k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} Q_{a+O(1)} \delta \psi_1 \\ & \quad \nabla^{-1} P_k Q_a [R_0 P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2] R_i P_{k_3} Q_{<k_3} \psi_3 \\ & \quad \quad - R_i P_{k_2} Q_{<k_2} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_2] R_0 P_{k_3} Q_{<k_3} \psi_3] \|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \sum_{k_2 + O(1) > a \geq \frac{k}{2}} 2^{a-k} 2^{\frac{k}{2+}} \|\chi_i(t)\|_{L_t^{4+}} \|P_{k_1} Q_{a+O(1)} \delta \psi_1\|_{L_t^2 L_x^\infty} \\ & \quad \|\phi_{2^{-\frac{i}{2+}}}(u) \psi_2\|_{L_t^\infty L_x^2} \|P_{k_3} Q_{<k_3} \psi_3\|_{L_t^{4-} L_x^{4+}} \end{aligned}$$

²¹Recall the suppressed localizations, see previous footnote.

We can bound this by

$$2^{a-k} 2^{\frac{i}{4+}} 2^{\frac{k}{2+}} 2^{k_1} 2^{-\frac{i}{4+}} 2^{-\frac{a}{2}} \tilde{c}_{k_1}$$

Our assumptions ensure that we may sum over $\frac{k}{2} < a < k_2 + O(1)$, resulting in an exponential gain in i . An identical argument may be used when $P_{k_3} Q_{<k_3} \psi_3$ is replaced by $P_{k_3} Q_{<k_3} [\phi_{2^{-\frac{i}{2+}}}(u) \psi_3]$, so we may replace both $P_{k_{2,3}} Q_{<k_{2,3}} \psi_{2,3}$ by $P_{k_{2,3}} Q_{<k_{2,3}} [(1 - \phi_{2^{-\frac{i}{2+}}}(u)) \psi_3]$. In that case we utilize the null-form identity recorded earlier: use

$$\begin{aligned} & R_0 P_{k_2} Q_{<k_2} [(1 - \phi_{2^{-\frac{i}{2+}}}) \psi_2] R_i P_{k_3} Q_{<k_3} [(1 - \phi_{2^{-\frac{i}{2+}}}) \psi_3] \\ & \quad - R_i P_{k_2} Q_{<k_2} [(1 - \phi_{2^{-\frac{i}{2+}}}) \psi_2] R_0 P_{k_3} Q_{<k_3} [(1 - \phi_{2^{-\frac{i}{2+}}}) \psi_3] \\ &= \frac{x_i}{r} [(\partial_t + \partial_r) \nabla^{-1} P_{k_2} Q_{<k_2} [(1 - \phi_{2^{-\frac{i}{2+}}}) \psi_2] (\partial_t - \partial_r) \nabla^{-1} P_{k_3} Q_{<k_3} [(1 - \phi_{2^{-\frac{i}{2+}}}) \psi_3] \\ & \quad - \frac{x_i}{r} [(\partial_t - \partial_r) \nabla^{-1} P_{k_2} Q_{<k_2} [(1 - \phi_{2^{-\frac{i}{2+}}}) \psi_2] (\partial_t + \partial_r) \nabla^{-1} P_{k_3} Q_{<k_3} [(1 - \phi_{2^{-\frac{i}{2+}}}) \psi_3] \end{aligned}$$

We can now exploit the fact that $|\frac{k}{2}| < \frac{(1+\mu)i}{2} < i(1-\delta)$ for μ small enough, as well as T_0 large enough. Thus we may move the multiplier $\chi_i(t)$ past the Fourier multiplier $\nabla^{-1} P_k Q_{>\frac{k}{2}}$ while trading in errors exponentially decreasing²² outside of $t \sim 2^i$. In other words, under the present assumptions on the frequencies, we may write schematically

$$\begin{aligned} & \chi_i(t) \nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi_1 \nabla^{-1} P_k Q_{>\frac{k}{2}}[,]] \\ &= \chi_i(t) \nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi_1 \nabla^{-1} P_k Q_{>\frac{k}{2}}(\chi_{i1}(t) + \chi_{i2}(t))[,]], \end{aligned}$$

where $\chi_{i1}(t)$ is supported on $t \sim 2^i$ while $|\chi_{i2}(t)| \lesssim 2^{-Ni}$ for $t < 2^{i+O(1)}$, as well as $|\chi_{i2}(t)| < t^{-N}$ for $t \gg 2^i$. It is then easy to verify that this leads to acceptable terms, so we may focus on the contribution of $\chi_{i1}(t)$. We shall want to move the operator $\partial_t + \partial_r$ past the Fourier localizer $P_{k_2} Q_{<k_2}$. We write

$$\begin{aligned} & P_{k_2} Q_{<k_2} [(1 - \phi_{2^{-\frac{i}{2+}}}(u)) \psi_2] \\ &= \rho_i(r) P_{k_2} Q_{<k_2} [(1 - \phi_{2^{-\frac{i}{2+}}}(u)) \psi_2] + (1 - \rho_i(r)) P_{k_2} Q_{<k_2} [(1 - \phi_{2^{-\frac{i}{2+}}}(u)) \psi_2], \end{aligned}$$

where $\rho_i(r)$ localizes smoothly to a disc of radius $\sim 2^{i-10}$ around the origin. Now on account of the fact that $\nabla^{-1} P_{k_2}$ is given by a convolution kernel which decays rapidly outside of a disc of radius $< 2^{\frac{i}{2}}$, we see by means of Proposition 2.2 that

$$\|\chi_{i1}(t) \rho_i(r) \nabla^{-1} P_{k_2} Q_{<k_2} [(1 - \phi_{2^{-\frac{i}{2+}}}(u)) \psi_2]\|_{L_x^\infty} \lesssim 2^{-\frac{3i}{2+}}$$

One then easily concludes that the contribution of this term is negligible: indeed, plugging it into the inner bracket instead of $P_{k_2} \psi_2$ and using schematic notation, we can estimate for example

$$\begin{aligned} & \|\chi_i(t) \nabla_{x,t} P_{k_0} [P_{k_1} Q_{<k_1-100} \delta \psi_1 \nabla^{-1} P_k Q_{>\frac{k}{2}}(\chi_{i1}(t)[,])]\|_{N[k_0]} \\ & \lesssim \|\chi_i(t) \nabla_{x,t} P_{k_0} [P_{k_1} Q_{<k_1-100} \delta \psi_1 \nabla^{-1} P_k Q_{>\frac{k}{2}}(\chi_{i1}(t)[,])]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim 2^{+\frac{k_0}{2}-k} \|P_{k_1} Q_{<k_1-100} \delta \psi_1\|_{L_t^\infty L_x^2} \|P_k Q_{>\frac{k}{2}}(\chi_{i1}(t)[,])\|_{L_t^2 L_x^2}, \end{aligned}$$

²²More precisely, these errors behave like $[1 + 2^{-(1-\delta)i}(2^{i-a} - t)]^{-N}$ for $t < 2^{i-a}$ and $[1 + 2^{-(1-\delta)i}(2^{i+b} - t)]^{-N}$ for $t > 2^{i+b}$, provided $\chi_i(t)$ is supported in $[2^{i-a}, 2^{i+b}]$.

where

$$\begin{aligned}
& \|P_k Q_{>\frac{k}{2}}(\chi_{i1}(t)[,])\|_{L_t^2 L_x^2} \\
& \lesssim \|\chi_{i1}(t)(\partial_t + \partial_r)[\rho_i(r)\nabla^{-1}P_{k_2}Q_{<k_2}[(1 - \phi_{2^{-\frac{i}{2+}}}(u))\psi_2]]\|_{L_t^2 L_x^\infty} \\
& \quad \|(\partial_t - \partial_r)\nabla^{-1}P_{k_3}Q_{<k_3}[(1 - \phi_{2^{-\frac{i}{2+}}}(u))\psi_2]\|_{L_t^\infty L_x^2} \\
& \lesssim 2^{-(1-\epsilon)i}
\end{aligned}$$

Our assumption $|k| < (1 + \mu)i$ shows that putting these estimates together gives an acceptable bound. The contribution of $P_{k_1}Q_{\geq k_1-100}\delta\psi_1$ is handled similarly. Now consider the contribution of the term with $\rho_i(r)$ replaced by $(1 - \rho_i(r))$. We use the fact (see e. g. [36]) that

$$\left[\frac{x_i}{r}(1 - \rho_i(r)), P_{k_2}\right]'' = ''2^{-k_2}\nabla\left(\frac{x_i}{r}(1 - \rho_i(r))\right)$$

where the latter expression stands for a weighted average of translates of the derivatives of $\frac{x_i}{r}(1 - \rho_i(r))$. Notice that

$$\|\nabla_x\left(\frac{x_i}{r}(1 - \rho_i(r))\right)\|_{L_x^\infty} \lesssim 2^{-i},$$

hence the contribution of the commutator is treated exactly as the contribution of the term $\rho_i(r)P_{k_2}Q_{<k_2}[(1 - \phi_{2^{-\frac{i}{2+}}}(u))\psi_2]$. This finally allows us to move the operator $\partial_t + \partial_r$ past the operator $\nabla^{-1}P_{k_2}Q_{<k_2}$. Arguing as before, one can also move the operator $\chi_{i1}(t)$ past the localizer $\nabla^{-1}P_{k_2}Q_{<k_2}$, generating acceptable error terms. Using lemma 2.4, we can now estimate

$$\begin{aligned}
& \|P_{k_2}Q_{<k_2}\nabla^{-1}\chi_{i1}(t)(1 - \rho_i(r))(\partial_t + \partial_r)[(1 - \phi_{2^{-\frac{i}{2+}}}(u))\psi_2]\|_{L_t^2 L_x^2} \\
& \quad \|(\partial_t - \partial_r)\nabla^{-1}P_{k_3}Q_{<k_3}[(1 - \phi_{2^{-\frac{i}{2+}}}(u))\psi_2]\|_{L_t^2 L_x^2} \\
& \lesssim \|P_{k_2}Q_{<k_2}\nabla^{-1}\chi_{i1}(t)(1 - \rho_i(r))(\partial_t + \partial_r)[(1 - \phi_{2^{-\frac{i}{2+}}}(u))\psi_2]\|_{L_t^2 L_x^\infty} \\
& \quad \|(\partial_t - \partial_r)\nabla^{-1}P_{k_3}Q_{<k_3}[(1 - \phi_{2^{-\frac{i}{2+}}}(u))\psi_2]\|_{L_t^\infty L_x^2} \\
& \lesssim 2^{-\frac{i}{2-}}
\end{aligned}$$

Proceeding as before, one deduces from this (and an analogous estimate for the term with $(\partial_t \pm \partial_r)$ interchanged) the following estimate for the full expression:

$$\|\chi_i(t)\nabla_{x,t}P_{k_0}[P_{k_1}\delta\psi_1\nabla^{-1}P_kQ_{>\frac{k}{2}}[,]]\|_{N[k_0]} \lesssim 2^{\frac{k_0}{2}-k}2^{-\frac{i}{2-}}\tilde{c}_{k_1}$$

Since $|k| < (1 + \mu)i$ and by assumption $k_0 \leq k + O(1)$, choosing μ small enough allows us to get an exponential gain in i . This concludes case (I.b).

(I.c): *None of (I.a), (I.b) hold, and $i \lesssim |k_1|$.* This then implies $|k_{2,3}|, |k| \ll i$, and we shall treat these as $O(1)$. Note that also necessarily $i \lesssim k_1 = k_0 + O(1)$. One first reduces $P_{k_{2,3}}\psi_{2,3}$ to modulation $< 2^{\delta i}$, where δ is very small but such that $2^{\delta i} \gg \max\{|k_{2,3}|, |k|\}$. We shall treat the latter quantities as $O(1)$. To achieve this, one estimates for example

$$\begin{aligned}
& \|P_{k_0}Q_{<k_0}\nabla_{x,t}\chi_i(t)[P_{k_1}Q_{<\delta i-10}\delta\psi_1\nabla^{-1}(1 - I)[P_{k_2}Q_{>\delta i}\psi_2, P_{k_3}Q_{<\delta i-10}\psi_3]]\|_{\dot{X}_{k_0}^{-1, -\frac{1}{2}, 1}} \\
& \lesssim 2^{-\frac{\delta}{2}i}\|P_{k_1}\delta\psi_1\|_{L_t^\infty L_x^2}\|P_{k_2}Q_{>\delta i}R_\nu\psi_2\|_{L_t^2 L_x^2}\|P_{k_3}\psi_3\|_{L_t^\infty L_x^\infty} \lesssim 2^{-\frac{\delta}{2+}i}\tilde{c}_{k_1}
\end{aligned}$$

Similarly, one has

$$\begin{aligned} & \|P_{k_0} Q_{<k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} Q_{\geq \delta i - 10} \delta \psi_1 \nabla^{-1} (1 - I) [P_{k_2} Q_{>\delta i} \psi_2, P_{k_3} Q_{<\delta i - 10} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \|P_{k_1} Q_{\geq \delta i - 10} \delta \psi_1\|_{L_t^2 L_x^2} \|R_\nu P_{k_2} Q_{>\delta i} \psi_2\|_{L_t^2 L_x^\infty} \|P_{k_3} \psi_3\|_{L_t^\infty L_x^\infty} + \text{etc} \\ & \lesssim 2^{-\frac{\delta i}{2+}} \tilde{c}_{k_1} \end{aligned}$$

The estimate when $P_{k_3} Q_{<\delta i - 10} \psi_3$ gets replaced by $P_{k_3} Q_{\geq \delta i - 10} \psi_3$ is more of the same. Moreover, we have by assumption

$$\begin{aligned} & \|P_{k_0} Q_{\geq k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} \delta \psi_1 \nabla^{-1} (1 - I) [P_{k_2} Q_{>\delta i} \psi_2, P_{k_3} \psi_3]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 1}} \\ & \lesssim 2^{-\frac{k_0}{2}} \|P_{k_1} \delta \psi_1\|_{L_t^M L_x^{2+}} \|R_\nu P_{k_2} Q_{>\delta i} \psi_2\|_{L_t^\infty L_x^M} \|P_{k_3} \psi_3\|_{L_t^{2+} L_x^\infty}, \end{aligned}$$

which leads to an acceptable estimate. Now we estimate

$$\begin{aligned} & \|P_{k_0} Q_{<k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} Q_{<k-10} \delta \psi_1 \\ & \quad \nabla^{-1} (1 - I) P_k [P_{k_2} Q_{<\delta i} \psi_2, P_{k_3} Q_{<\delta i - 10} \psi_3]]\|_{\dot{X}_{k_0}^{-1, -\frac{1}{2}, 1}} \\ & \lesssim 2^{2\delta i} \|P_{k_1} \delta \psi_1\|_{L_t^\infty L_x^2} \|P_{k_2} Q_{<\delta i} \chi_i(t) \psi_2\|_{L_t^M L_x^\infty} \|P_{k_3} Q_{<\delta i - 10} \psi_3\|_{L_t^{2+} L_x^\infty} \\ & \lesssim 2^{(2\delta - \frac{1}{2+})i} \tilde{c}_{k_1} \end{aligned}$$

(I.d): *None of (I.a), (I.b), (I.c) hold.* In this case, we may treat all of $k, k_i, i = 1, 2, 3$ as $O(1)$. The exponential gain in i is again obtained as in the preceding case. This concludes the treatment of case (I).

(II): *The 2nd term.* This term is significantly simpler than the preceding one: note that if at least one of $|k|, |k_i|, i = 0, \dots, 3$, is of size at least comparable to $\log_2 T_0$, one gets an exponential gain in T_0 from lemma 3.16 in conjunction with lemma 4.2 and the calculations in (I), provided $P_{k_1} \psi$ is of first type. If it is of 2nd type, and $P_{k_3} \psi_3$ of 2nd type as well, one also argues as in (I). If $P_{k_1} \psi_1$ is of 2nd type, but $P_{k_3} \psi_3$ of first type, one uses the estimate

$$\|P_k [R_\beta P_{k_2} Q_{<k_2} \delta \psi_2 R_j P_{k_3} Q_{<k_3} \psi_3 - R_j P_{k_2} Q_{<k_2} \delta \psi_2 R_\beta P_{k_3} Q_{<k_3} \psi_3]\|_{L_t^\infty L_x^2} \lesssim 2^k$$

and (with a similar estimate when $Q_{<k_3}$ is replaced by $Q_{\geq k_3}$)

$$\begin{aligned} & \|P_k [R_\beta P_{k_2} Q_{\geq k_2} \delta \psi_2 R_j P_{k_3} Q_{<k_3} \psi_3 - R_j P_{k_2} Q_{\geq k_2} \delta \psi_2 R_\beta P_{k_3} Q_{<k_3} \psi_3]\|_{L_t^2 L_x^2} \\ & \lesssim 2^{\frac{\min\{k, k_2, k_3\}}{2}} 2^{-\delta |k_2 - k_3|} \frac{\tilde{c}_{k_2}}{\epsilon} \tilde{c}_{k_3} \end{aligned} \tag{20}$$

Using the first of these, one gets for example when $k_0 = k_1 + O(1)$

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{>k_0} \chi(\frac{t}{T_0}) [P_{k_1} \psi_1 \Delta^{-1} \sum_{j=1,2} \partial_j (1 - I) P_k [R_\beta P_{k_2} Q_{<k_2} \delta \psi_2 R_j P_{k_3} Q_{<k_3} \psi_3 \\ & \quad - R_j P_{k_2} Q_{<k_2} \delta \psi_2 R_\beta P_{k_3} Q_{<k_3} \psi_3]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim 2^{-\frac{k_0}{2}} \|P_{k_1} \psi_1\|_{L_t^2 L_x^{2+}} \|\Delta^{-1} \sum_{j=1,2} \partial_j (1 - I) P_k [R_\beta P_{k_2} Q_{<k_2} \delta \psi_2 R_j P_{k_3} Q_{<k_3} \psi_3 \\ & \quad - R_j P_{k_2} Q_{<k_2} \delta \psi_2 R_\beta P_{k_3} Q_{<k_3} \psi_3]\|_{L_t^\infty L_x^M} \end{aligned}$$

which is bounded by

$$2^{(k-k_0)(1-\epsilon)} 2^{\delta[\min\{k_2, k_3\} - \max\{k_2, k_3\}]} \tilde{c}_{k_1} \frac{\tilde{c}_{k_2}}{\epsilon} \frac{\tilde{c}_{k_3}}{\epsilon}$$

Assuming $\max\{|k|, |k_i|\}$ to be at least comparable to $\delta \log_2 T_0$ and summing over $k, k_{2,3}$ results thus in the estimate $\tilde{c}_{k_1} T_0^{-\mu}$. Using the 2nd of the above inequalities, (20), and placing $P_{k_1} \psi_1$ into $L_t^\infty L_x^2$ results in a similar estimate provided $k_0 = k_1 + O(1)$. The remaining frequency interactions $k_0 < k_1$ etc. are handled similarly, as well as the contribution when $Q_{>k_0}$ is replaced by $Q_{\leq k_0}$. Thus assume now that all the occurring frequencies $|k|, |k_i|, i = 0, \dots, 3$ are of size $< \delta \log_2 T_0$. Then we use Proposition 2.2 directly in conjunction with 3.4(b) resp. lemma 4.3, to get

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{>k_0} \chi(\frac{t}{T_0}) [P_{k_1} \psi_1 \Delta^{-1} \sum_{j=1,2} \partial_j (1-I) P_k [R_\beta P_{k_2} \delta \psi_2 R_j P_{k_3} \psi_3 \\ & \quad - R_j P_{k_2} \delta \psi_2 R_\beta P_{k_3} \psi_3]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim \|P_{k_1} \chi(\frac{t}{T_0}) \psi_1\|_{L_t^\infty L_x^\infty} \|\Delta^{-1} \sum_{j=1,2} \partial_j (1-I) P_k [R_\beta P_{k_2} \delta \psi_2 R_j P_{k_3} \psi_3 \\ & \quad - R_j P_{k_2} \delta \psi_2 R_\beta P_{k_3} \psi_3]\|_{L_t^2 L_x^2} \\ & \lesssim T_0^{-\frac{1}{2+}} [\tilde{c}_{k_3} + \tilde{c}_{k_2}] \end{aligned}$$

But using the definition of frequency envelope we have $\tilde{c}_{k_{2,3}} \lesssim T_0^{\sigma\delta} \tilde{c}_{k_1}$, so we arrive at an acceptable estimate upon summing over the admissible frequency ranges. If we replace $Q_{>k_0}$ by $Q_{\leq k_0}$, we can for example first reduce $P_{k_1} \psi_1$ to modulation $< 2\delta \log_2 T_0$, then reduce both of $P_{k_{2,3}} \psi_{2,3}$ to modulation $< 3\delta \log_2 T_0$, and finally estimate

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{\leq k_0} \chi(\frac{t}{T_0}) [P_{k_1} Q_{k-10 < \cdot < 2\delta \log_2 T_0} \psi_1 \\ & \quad \Delta^{-1} \sum_{j=1,2} \partial_j (1-I) P_k [R_\beta P_{k_2} Q_{< 3\delta \log_2 T_0} \delta \psi_2 R_j P_{k_3} Q_{< 3\delta \log_2 T_0} \psi_3 \\ & \quad - R_j P_{k_2} Q_{< 3\delta \log_2 T_0} \delta \psi_2 R_\beta P_{k_3} Q_{< 3\delta \log_2 T_0} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \|P_{k_1} Q_{k-10 < \cdot < 2\delta \log_2 T_0} \psi_1\|_{L_t^2 L_x^M} \|\nabla_{x,t} \nabla^{-1} P_{k_2} Q_{< 3\delta \log_2 T_0} \delta \psi_2\|_{L_t^M L_x^{2+}} \\ & \quad \|\nabla_{x,t} \nabla^{-1} P_{k_3} Q_{< 3\delta \log_2 T_0} \psi_3\|_{L_t^{2+} L_x^\infty} \end{aligned}$$

Arranging that $\frac{1}{2} - \frac{1}{2+} \gg \delta$ and using the usual properties of the frequency envelope easily results in the desired bound. If one replaces $P_{k_1} Q_{k-10 < \cdot < 2\delta \log_2 T_0} \psi_1$ by $P_{k_1} Q_{\leq k} \psi_1$, one can estimate the output with respect to $\|\cdot\|_{\dot{X}^{-1, -\frac{1}{2}, 1}}$ in the same manner. This finishes case (II) and thereby the large modulation case **(A)**.

(B): The small modulation case. We now study the expressions

$$\begin{aligned} (I) : & \partial^\beta [\delta \psi_\alpha \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta \psi_2 R_j \psi_3 - R_j \psi_2 R_\beta \psi_3]] \\ (II) : & \partial^\beta [\psi_\alpha \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta \delta \psi_2 R_j \psi_3 - R_j \delta \psi_2 R_\beta \psi_3]], \end{aligned}$$

as well as the analogous expressions

$$(III) : \partial^\beta [\delta \psi_\beta \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta \psi_2 R_j \psi_3 - R_j \psi_2 R_\beta \psi_3]]$$

$$(IV) : \partial^\beta [\psi_\beta \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta \delta \psi_2 R_j \psi_3 - R_j \delta \psi_2 R_\beta \psi_3]]$$

$$(V) : \partial_\alpha [\delta \psi^\nu \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\nu \psi_2 R_j \psi_3 - R_j \psi_2 R_\nu \psi_3]]$$

$$(VI) : \partial_\alpha [\psi^\nu \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\nu \delta \psi_2 R_j \psi_3 - R_j \delta \psi_2 R_\nu \psi_3]]$$

As is, these terms cannot yet be well estimated, and we need to further decompose the first input into a gradient part and elliptic error term: thus for example we write in term (I)

$$\delta \psi_\alpha = R_\alpha \delta \psi + \delta \chi_\alpha$$

Relegating the error terms involving χ_α until later, we substitute $R_\alpha \delta \psi$ for $\delta \psi_\alpha$, and similarly for the other terms (II)-(VI). We commence with the sum of first and fifth term in the list:

(I): **(I+V)**: As in the large modulation case, we shall have to consider various types of frequency interactions. We also reiterate the decomposition

$$\begin{aligned} & \chi\left(\frac{t}{T_0}\right) \partial^\beta [R_\alpha \delta \psi \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta \psi_2 R_j \psi_3 - R_j \psi_2 R_\beta \psi_3]] \\ & + \chi\left(\frac{t}{T_0}\right) \partial_\alpha [R^\nu \delta \psi \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\nu \psi_2 R_j \psi_3 - R_j \psi_2 R_\nu \psi_3]] \\ & = \sum_{i > \log_2 T_0} \chi_i(t) [\partial^\beta [R_\alpha \delta \psi \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta \psi_2 R_j \psi_3 - R_j \psi_2 R_\beta \psi_3]] \\ & + \partial_\alpha [R^\nu \delta \psi \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\nu \psi_2 R_j \psi_3 - R_j \psi_2 R_\nu \psi_3]]] \end{aligned}$$

We frequency-localize this to obtain the following expression:

$$\begin{aligned} & \chi_i(t) P_{k_0} [\partial^\beta [R_\alpha P_{k_1} \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]] \\ & + \partial_\alpha [R^\nu P_{k_1} \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\nu P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\nu P_{k_3} \psi_3]]] \end{aligned}$$

Now we subdivide into the following possibilities:

(I.a): *One of the following options hold: $i \lesssim |k_2|$, $i \lesssim |k_3|$, $i \lesssim |k_0 - k_1|$, $i \lesssim \min\{|k - k_1|, |k - k_2|\}$.* In this case, we obtain the desired estimate involving an exponential gain in i from 3.4(c) as well as lemma 3.16 if both $P_{k_2,3} \psi_{2,3}$ are of the first type. If at least one of them is of the 2nd type, this is again straightforward due to the strong estimates satisfied by these: then we have

$$\begin{aligned} & \|P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]\|_{L_t^1 L_x^\infty} \\ & \lesssim 2^{\delta[\min\{k, k_2, k_3\} - \max\{k, k_2, k_3\}]} \frac{\tilde{c}_{k_2}}{\epsilon} \frac{\tilde{c}_{k_3}}{\epsilon} \end{aligned}$$

$$\begin{aligned}
& \|P_k \Delta^{-1} \sum_{j=1,2} \partial_j I[R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]\|_{L_t^1 L_x^{2+}} \\
& \lesssim 2^{-k(1-\epsilon)} 2^{\delta[\min\{k, k_2, k_3\} - \max\{k, k_2, k_3\}]} \frac{\tilde{C}_{k_2}}{\epsilon} \frac{\tilde{C}_{k_3}}{\epsilon}
\end{aligned}$$

From this we get

$$\begin{aligned}
& \|P_{k_0} Q_{<k_0} [\partial^\beta [R_\alpha P_{k_1} \delta \psi_1 \\
& \quad P_k \Delta^{-1} \sum_{j=1,2} \partial_j I[R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim 2^{\delta[\min\{k, k_0, \dots, k_3\} - \max\{k, k_0, \dots, k_3\}]} \frac{1}{\epsilon^2} \prod_{i=1,2,3} \tilde{C}_{k_i}
\end{aligned}$$

The contribution when one has $P_{k_0} Q_{\geq k_0}$ in front is even simpler and left out. Term (V) is treated by exact analogy.

(I.b): $i \lesssim |k|$, $i \lesssim |k_1|$, and none of the properties in (I.a) hold. This implies $k, k_1 \lesssim -i$. We may and shall assume $k - k_1 = O(1)$, $k_0 - k_1 = O(1)$. Also, we may and shall assume that $|k_{2,3}| = O(1)$. We shall again treat term (I), term (V) being treated analogously. We start out by observing that we may assume $k > -i(1 + \mu)$ for any $\mu > 0$. Indeed, assume the opposite. Again considering term (I), we have

$$\begin{aligned}
& \|P_{k_0} Q_{<k_0} \chi_i(t) [\partial^\beta Q_{>k_0+100} [R_\alpha P_{k_1} \delta \psi_1 P_k \\
& \quad \Delta^{-1} \sum_{j=1,2} \partial_j I[R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \sum_{k_0+100 \leq j \leq -i+O(1)} \|P_{k_0} Q_{<k_0} \chi_i(t) [\partial^\beta Q_j [R_\alpha P_{k_1} Q_{j+O(1)} \delta \psi_1 P_k \\
& \quad \Delta^{-1} \sum_{j=1,2} \partial_j I[R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \sum_{k_0+100 \leq j \leq -i+O(1)} 2^j \|\chi_i(t)\|_{L_t^2} \|R_\alpha P_{k_1} Q_{j+O(1)} \delta \psi_1\|_{L_t^2 L_x^2} \\
& \quad \|P_k [\nabla^{-1} P_{k_2} \psi_2 P_{k_3} \psi_3]\|_{L_t^\infty L_x^\infty}
\end{aligned}$$

One checks that this is estimated by $\lesssim \tilde{C}_{k_1} 2^{k_0}$, which is acceptable. Similarly, we estimate

$$\begin{aligned}
& \|P_{k_0} Q_{<k_0} \chi_i(t) [\partial^\beta Q_{\leq k_0+100} [R_\alpha P_{k_1} \delta \psi_1 \\
& \quad P_k \Delta^{-1} \sum_{j=1,2} \partial_j I[R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \|\chi_i(t)\|_{L_t^2} \|P_{k_1} \delta \psi_1\|_{L_t^\infty L_x^\infty} \\
& \quad \|P_k \Delta^{-1} \sum_{j=1,2} \partial_j I[R_\beta P_{k_2} \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \psi_2 R_\beta P_{k_3} \psi_3]\|_{L_t^2 L_x^2},
\end{aligned}$$

which, upon using 3.4(b) as well as lemma 4.3, can be estimated by $\lesssim \tilde{C}_{k_0} 2^{\frac{k_0+i}{2}}$, which is acceptable.

The estimates when $Q_{<k_0}$ is replaced by $Q_{\geq k_0}$ are similar. Thus we now assume that $|k| < i(1 + \mu)$. Arguing as in case (A)(I.b), we may replace $P_k I$ by $P_k Q_{<\frac{k}{2}}$ while only generating acceptable error terms. We may then move the multiplier $\chi_i(t)$ past the operator $P_k \nabla^{-1} Q_{<\frac{k}{2}}$, and transform the latter back into $\nabla^{-1} P_k I$

innocuously. Proceeding as in [23], we intend to exploit the null-structure of the expression. Before being able to do so, we need to effect a few more reductions: we need to reduce the modulation of the output and first input to size $< 2^{k_0,1}$, respectively. For this, note that

$$\begin{aligned}
& \|P_{k_0} Q_{\geq k_0+100} [\partial^\beta [R_\alpha P_{k_1} \delta \psi_1 P_k \\
& \quad \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \chi_i(t) \psi_2 R_\beta P_{k_3} \psi_3]]] \|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\
& \lesssim \|P_{k_0} Q_{\geq k_0+100} [\partial^\beta [R_\alpha P_{k_1} Q_{\geq k_1} \delta \psi_1 P_k \\
& \quad \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3 - R_j P_{k_2} \chi_i(t) \psi_2 R_\beta P_{k_3} \psi_3]]] \|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\
& \lesssim 2^{-\frac{k_0}{2}} \|R_\alpha P_{k_1} Q_{\geq k_1} \delta \psi_1\|_{L_t^2 L_x^\infty} \|P_k [\nabla^{-1} P_{k_2} \psi_2 P_{k_3} R_\beta \psi_3]\|_{L_t^\infty L_x^2} \\
& \lesssim 2^k \tilde{c}_{k_1} \left[\frac{\tilde{c}_{k_2}}{\epsilon} + \frac{\tilde{c}_{k_3}}{\epsilon} \right].
\end{aligned}$$

We can also reduce $P_{k_2,3} \psi_{2,3}$ to modulation $< 2^{\max\{k_2, k_3\} + O(1)}$. For this, note that

$$\begin{aligned}
& \|P_k I [\nabla^{-1} P_{k_2} \psi_2 P_{k_3} Q_{>\max\{k_2, k_3\}+100} \psi_3] \|_{L_t^1 L_x^2} \\
& \lesssim \sum_{j>\max\{k_2, k_3\}+100} 2^k \|\nabla^{-1} P_{k_2} Q_{j+O(1)} \psi_2\|_{L_t^2 L_x^2} \|P_{k_3} Q_j \psi_3\|_{L_t^2 L_x^2} \\
& \lesssim \sum_{j>\max\{k_2, k_3\}+100} 2^k 2^{-(2-2\mu)j} 2^{(1-2\mu)\max\{k_2, k_3\}} \\
& \quad \|\nabla^{-1} P_{k_2} \psi_2\|_{\dot{X}_{k_2}^{-(\frac{1}{2}-\mu), 1-\mu, 1}} \|P_{k_3} \psi_3\|_{\dot{X}_{k_3}^{-(\frac{1}{2}-\mu), 1-\mu, 1}} \\
& \lesssim 2^{-k_2} \frac{\tilde{c}_{k_2}}{\epsilon} \frac{\tilde{c}_{k_3}}{\epsilon}
\end{aligned}$$

From this one deduces that for $a = \max\{k_2, k_3\} + 100$

$$\begin{aligned}
& \|P_{k_0} Q_{<k_0+O(1)} [\partial^\beta [R_\alpha P_{k_1} \delta \psi_1 \\
& P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [R_\beta P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} Q_{>a} \psi_3 - R_j P_{k_2} \chi_i(t) \psi_2 R_\beta P_{k_3} Q_{>a} \psi_3]]] \|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \|R_\alpha P_{k_1} Q_{<k_1+O(1)} \delta \psi_1\|_{L_t^\infty L_x^\infty} \|P_k I [\nabla^{-1} P_{k_2} \psi_2 P_{k_3} Q_{>\max\{k_2, k_3\}+100} \psi_3]\|_{L_t^1 L_x^2} \\
& \lesssim 2^{k-k_2} \tilde{c}_{k_1} \frac{\tilde{c}_{k_2}}{\epsilon} \frac{\tilde{c}_{k_3}}{\epsilon}
\end{aligned}$$

One can sum over the appropriate range of $k_{2,3}$, deducing the desired estimate. We shall always assume these reductions of modulation, but sometimes omit them to simplify notation. Now we expand the null-structure as in [22], [23]: schematically

we have

$$\begin{aligned}
& 2 \sum_{j=1}^2 \Delta^{-1} \partial_j [R_\nu f R_j g - R_j f R_\nu g] \partial^\nu h \\
&= \sum_{j=1}^2 \square [\Delta^{-1} \partial_j [\nabla^{-1} f R_j g] h] - \sum_{j=1}^2 \square \Delta^{-1} \partial_j [\nabla^{-1} f R_j g] h \\
&- \sum_{j=1}^2 \Delta^{-1} \partial_j [\nabla^{-1} f R_j g] \square h - \nabla^{-1} f \square ((\nabla^{-1} g) h) \\
&+ \nabla^{-1} f \square (\nabla^{-1} g) h + \nabla^{-1} f (\nabla^{-1} g) \square h,
\end{aligned} \tag{21}$$

$$\begin{aligned}
& \sum_{j=1}^2 \Delta^{-1} \partial_j \partial^\nu [R_\nu f R_j g - R_j f R_\nu g] h \\
&= \sum_{j=1}^2 [\Delta^{-1} \partial_j \square [\nabla^{-1} f g] h - \frac{1}{2} \square [\nabla^{-1} f \nabla^{-1} g] h \\
&+ \frac{1}{2} \square \nabla^{-1} f \nabla^{-1} g h - \frac{1}{2} \nabla^{-1} f \square g h]
\end{aligned} \tag{22}$$

The first of these identities is useful when the outer derivative ∂^β falls on the first input $R_\alpha \delta \psi_1$. The 2nd is useful provided the outer derivative lands on the inner square bracket. We shall treat each of these terms. Clearly, the terms in the 2nd expansion are almost identical to the ones in the first. We treat the first in detail, the 2nd being treated similarly.

(I.b.1) *The first term in the expansion.* This is the expression

$$P_{k_0} Q_{<k_0} \square [R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j I[\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]]$$

This is straightforward to estimate: we have

$$\begin{aligned}
& \|P_{k_0} Q_{<k_0} \square [R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j I[\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]]\|_{\dot{X}_{k_0}^{-1-\frac{1}{2},1}} \\
& \lesssim 2^{-\frac{k_0}{2}} \|R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1\|_{L_t^\infty L_x^M} \|\nabla^{-1} P_{k_2} \chi_i(t) \psi_2\|_{L_t^{2+} L_x^\infty} \|P_{k_3} \psi_3\|_{L_t^M L_x^{2+}} \lesssim 2^{\frac{k_1}{2+}} \tilde{c}_{k_1}
\end{aligned}$$

(I.b.2) *The 2nd term in the expansion.* This is the expression

$$P_{k_0} Q_{<k_0} [R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j \square I[\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]]$$

This turns out to be significantly more complicated. The reason for this is that we need to exploit the bilinear inequality 3.4(g); using Strichartz type norms here appears to result in a loss in the low frequencies, or in i . The only way we can possibly squeeze out a small gain in i is to exploit the temporal cutoff $\chi_i(t)$ applied to $P_{k_2} \psi_2$. This is a non-trivial task on account of the fact that the only way to place the inner bracket $[\cdot]$ into $L_t^2 L_x^2$ appears to involve null-frame spaces. Our main tool for this is the following lemma:

Lemma 4.4. *Let C be a sufficiently large number. The following limits hold: for any k with $|k| < \frac{\epsilon}{1000C}i$, and arbitrary $\epsilon > 0$,*

$$\lim_{i \rightarrow \infty} \|P_k Q_{[k-(1-\epsilon)i-C, k]}[\chi_i(t)\psi_\nu]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} = 0$$

More precisely, for appropriate $\mu(\epsilon) > 0$, we have

$$\|P_k Q_{[k-(1-\epsilon)i-C, k]}[\chi_i(t)\psi_\nu]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \lesssim 2^{-\mu(\epsilon)i}$$

Next, denote by $\chi_{i, \kappa}^c$ a smooth bump function which localizes to the complement in the $2^i/C$ -neighborhood of the (physical) light cone of a slab of length 2^{i-1} centered at time $t = 2^i$, angular opening 2κ with $|\kappa| \sim 2^{i(\frac{\epsilon-1}{2})}$ and distance $\lesssim 2^i/C$ from the light cone. Then with the same assumption on k ,

$$\lim_{i \rightarrow \infty} \sum_{\pm} \sup_{l \in [\frac{\epsilon-1}{2}i, -10]} 2^{-\frac{l}{2}} \left(\sum_{\tilde{\kappa} \in K_l} \|P_{k, \tilde{\kappa}} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c(t, x) P_{k, \kappa} Q_{< k+i(\epsilon-1)-C}[\chi_i(t)\psi_\nu]\|_{PW[\pm \tilde{\kappa}]}^2 \right)^{\frac{1}{2}} = 0$$

More precisely, this quantity decays like $2^{-\mu i}$ for suitable $\mu > 0$.

Proof : We first estimate

$$\|\nabla_{x,t} P_k Q_{> k-(1-\epsilon)i}[\chi_i(t) \frac{\mathbf{x}}{\mathbf{y}}]\|_{A[k]}, \|\nabla_{x,t} P_k Q_{> k-(1-\epsilon)i}[\chi_i(t) \ln \mathbf{y}]\|_{A[k]}$$

for arbitrary $k \in \mathbf{Z}$ with $|k| < \frac{\epsilon}{C}i$. Observe that

$$\square[\chi_i(t) \frac{\mathbf{x}}{\mathbf{y}}] = \chi_i''(t) \frac{\mathbf{x}}{\mathbf{y}} + \chi_i'(t) \partial_t \left(\frac{\mathbf{x}}{\mathbf{y}} \right) + \chi_i(t) \square \left[\frac{\mathbf{x}}{\mathbf{y}} \right]$$

Therefore, we obtain for $j \in [k - (1-\epsilon)i, k]$:

$$\begin{aligned} \|P_k Q_j \nabla_x \chi_i(t) \frac{\mathbf{x}}{\mathbf{y}}\|_{A[k]} &\lesssim 2^{-\frac{j}{2}} [\|P_k(\chi_i''(t) \frac{\mathbf{x}}{\mathbf{y}})\|_{L_t^2 L_x^2} + \|P_k(\chi_i'(t) \partial_t (\frac{\mathbf{x}}{\mathbf{y}}))\|_{L_t^2 L_x^2} \\ &\quad + \|P_k Q_j(\chi_i(t) \square[\frac{\mathbf{x}}{\mathbf{y}}])\|_{L_t^2 L_x^2}] \end{aligned}$$

The first two terms on the right hand side are elementary to estimate, using Holder's inequality, finite propagation speed and the energy inequality:

$$\|P_k(\chi_i''(t) \frac{\mathbf{x}}{\mathbf{y}})\|_{L_t^2 L_x^2} \lesssim \|\chi_i''(t)\|_{L_t^2} \|P_k(\frac{\mathbf{x}}{\mathbf{y}})\|_{L_t^\infty L_x^2} \lesssim 2^{\frac{j}{2}} 2^{-2i} 2^i \lesssim 2^{-\frac{j}{2}}$$

$$\|P_k(\chi_i'(t) \partial_t (\frac{\mathbf{x}}{\mathbf{y}}))\|_{L_t^2 L_x^2} \lesssim \|\chi_i'(t)\|_{L_t^2} \|\partial_t (\frac{\mathbf{x}}{\mathbf{y}})\|_{L_t^\infty L_x^2} \lesssim 2^{-\frac{j}{2}}$$

We conclude that the contribution from these terms is at most

$$\lesssim \sum_{j > k-i(1-\epsilon)} 2^{-\frac{j}{2}} 2^{-\frac{j}{2}} \lesssim 2^{-\frac{\epsilon}{2+1}i}.$$

Proceeding to the last term above, we use the null-structure in (2) as well as Proposition 2.2 to get

$$\begin{aligned} &\|P_k(\chi_i(t) \square[\frac{\mathbf{x}}{\mathbf{y}}])\|_{L_t^2 L_x^2} \\ &\lesssim \|\chi_i(t) \frac{(\partial_t + \partial_r)\mathbf{x}}{\mathbf{y}}\|_{L_t^2 L_x^\infty} \|\chi_i(t) \frac{(\partial_t - \partial_r)\mathbf{x}}{\mathbf{y}}\|_{L_t^\infty L_x^2} + \text{etc} \lesssim 2^{-i} \end{aligned}$$

This establishes the claim for $\frac{\mathbf{x}}{\mathbf{y}}$ because of 3.4(d). As far as $\partial_t P_k Q_{> k+(\epsilon-1)i}(\frac{\mathbf{x}}{\mathbf{y}})$ is concerned this only differs from the preceding as far as the estimate for

$P_k Q_{\geq k} R_0 \partial_t (\frac{x}{y})$. This is treated by using the equation for $\square(\frac{x}{y})$ and arguing as before. The estimate for $\ln y$ is similar. We now establish the 2nd inequality stated in the lemma provided ψ_ν is replaced by $\nabla_{x,t} \frac{x}{y}$, $\nabla_{x,t} \ln y$. First, let $\phi \in C^\infty(\mathbf{R}^{2+1})$ be a rotationally symmetric free wave with $\phi[0, x] = (0, g(x))$. We use the representation formula

$$\begin{aligned} P_{k,\kappa} \phi(t, x) &= c \int_{S^1} \mathcal{F}(\hat{g}(|\xi|) m_k(|\xi|) |\xi|^2) a_\kappa(\omega) (-t + x \cdot \omega) d\omega \\ &\quad - c \int_{S^1} \mathcal{F}(\hat{g}(|\xi|) m_k(|\xi|) |\xi|^2) a_\kappa(\omega) (t + x \cdot \omega) d\omega \end{aligned}$$

where (committing abuse of notation) we wrote $\hat{g}(\xi) = \hat{g}(|\xi|)$. Now assume that $\|(1 + \nabla_{|\xi|}^\alpha) \hat{g}(|\xi|)\|_{L_{|\xi|}^2} < C$. This implies that

$$\|\chi_i(\cdot) \mathcal{F}(\hat{g}(|\xi|))(\cdot)\|_{L^2} \lesssim 2^{-\alpha i}$$

Now observe that on the support of $\chi_{i,\mp\kappa}^c$, we have

$$|t \pm x \cdot \omega| \geq -|t - |x|| + |x| \pm x \cdot \omega \geq |t - |x|| - |x| \sin < x, \mp \omega >^2 \geq c 2^{i\epsilon}$$

Next, observe that the multiplier $\chi_{i,\mp\kappa}^c(t, x)$ smears out the Fourier support in the angular direction by an amount $\sim 2^{-\frac{(1+\epsilon)i}{2}}$, while $|\kappa| \sim 2^{\frac{i(\epsilon-1)}{2}}$ and $|k| < \frac{\epsilon}{C}i$. This entails that one can include an operator $P'_{k,\kappa}$ with the same properties as $P_{k,\kappa}$ in the expression and reason as follows²³:

$$\begin{aligned} &2^{-\frac{i}{2}} \|P_{k,\tilde{\kappa}} \sum_{\kappa \in K_{\frac{(\epsilon-1)}{2}i}} \chi_{i,\mp\kappa}^c(t, x) P_{k,\kappa} Q^\pm \phi\|_{PW[\pm\tilde{\kappa}]} \\ &= 2^{-\frac{i}{2}} \|P_{k,\tilde{\kappa}} \sum_{\kappa \in K_{\frac{(\epsilon-1)}{2}i}} P'_{k,\kappa} \chi_{i,\mp\kappa}^c(t, x) P_{k,\kappa} Q^\pm \phi\|_{PW[\pm\tilde{\kappa}]} \\ &\lesssim 2^{i\frac{1-\epsilon}{4}} \left(\sum_{\kappa \in K_{\frac{(\epsilon-1)}{2}i}, \kappa \cap 2\tilde{\kappa} \neq \emptyset} \|\chi_{i,\mp\kappa}^c(t, x) P_{k,\kappa} Q^\pm \phi\|_{PW[\pm\kappa]}^2 \right)^{\frac{1}{2}} \end{aligned}$$

Using Plancherel's theorem, we can further estimate this as

$$\begin{aligned} &\lesssim 2^{i\frac{1-\epsilon}{4}} \left(\sum_{\kappa \in K_{\frac{(\epsilon-1)}{2}i}, \kappa \cap 2\tilde{\kappa} \neq \emptyset} \left[\int_{\omega \in \kappa} \|\chi_{t \pm x \cdot \omega > 2^{i\epsilon}} \mathcal{F}(\hat{g}(|\xi|))(t \pm x \cdot \omega)\|_{L^2} a_\kappa(\omega) d\omega \right]^2 \right)^{\frac{1}{2}} \\ &\lesssim 2^{-\epsilon\alpha i} \left(\sum_{\kappa \in K_{\frac{(\epsilon-1)}{2}i}, \kappa \cap 2\tilde{\kappa} \neq \emptyset} \left[\int_{\omega \in \kappa} \|(1 + \nabla_{|\xi|}^\alpha) \hat{g}(|\xi|)\|_{L_{|\xi|}^2} a_\kappa(\omega) d\omega \right]^2 \right)^{\frac{1}{2}} \end{aligned}$$

Using Cauchy Schwarz' inequality, this leads to the estimate

$$\begin{aligned} &2^{-\frac{i}{2}} \sum_{\pm} \left(\sum_{\tilde{\kappa} \in K_l} \|P_{k,\tilde{\kappa}} \sum_{\kappa \in K_{\frac{(\epsilon-1)}{2}i}} \chi_{i,\mp\kappa}^c(t, x) P_{k,\kappa} Q^\pm \phi\|_{PW[\pm\tilde{\kappa}]}^2 \right)^{\frac{1}{2}} \\ &\lesssim 2^{-\alpha\epsilon i} \|(1 + \nabla_{|\xi|}^\alpha) \hat{g}(|\xi|)\|_{L_{|\xi|}^2} \end{aligned}$$

²³Let Q^\pm microlocalize to the upper or lower half-space $\tau > < 0$.

Now we proceed to the inhomogeneous situation at hand: let $S(t)$ denote the free wave propagator, i. e. $\square[S(t)(f, g)] = 0$, $S(0)(f, g) = f$, $\partial_t[S(0)(f, g)] = g$. Also, let $U(t)g = S(t)(0, f)$. Then we can write

$$\frac{\mathbf{x}}{\mathbf{y}}(t, \cdot) = S(t)\left(\frac{\mathbf{x}}{\mathbf{y}}(0, \cdot), \partial_t\left[\frac{\mathbf{x}}{\mathbf{y}}\right](0, \cdot)\right) + \int_0^t U(t-s)\square\left[\frac{\mathbf{x}}{\mathbf{y}}\right](s, \cdot)ds$$

Reasoning as above, we immediately get the desired estimate for the linear part. As concerns the inhomogeneity, we decompose the integral as

$$\int_0^t U(t-s)\square\left[\frac{\mathbf{x}}{\mathbf{y}}\right](s, \cdot)ds = \int_0^{2^{\frac{\epsilon}{C}i}} U(t-s)\square\left[\frac{\mathbf{x}}{\mathbf{y}}\right](s, \cdot)ds + \int_{2^{\frac{\epsilon}{C}i}}^t U(t-s)\square\left[\frac{\mathbf{x}}{\mathbf{y}}\right](s, \cdot)ds$$

Then we observe that from the argument given above we have for $l \in [\frac{(\epsilon-1)i}{2}, -10]$

$$\begin{aligned} & 2^{-\frac{l}{2}} \left(\sum_{\tilde{\kappa} \in K_l} \|P_{k, \tilde{\kappa}} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k, \kappa} Q^\pm \left[\int_0^{2^{\frac{\epsilon}{C}i}} U(t-s)\square\left[\frac{\mathbf{x}}{\mathbf{y}}\right](s, \cdot)ds \right] \|_{PW[\pm \tilde{\kappa}]}^2 \right)^{\frac{1}{2}} \\ & \lesssim 2^{-\epsilon\alpha i} 2^{-k} \|x^\alpha \square\left[\frac{\mathbf{x}}{\mathbf{y}}\right]\|_{L_t^1 L_x^2} \lesssim 2^{-\frac{\epsilon}{2}\alpha i}, \end{aligned}$$

provided $\frac{1}{C} < \alpha$ and we choose $\alpha > 0$ small enough, by the proof of Corollary 2.6 and finite propagation speed. Next, we can estimate by using 3.4(d):

$$\begin{aligned} & 2^{\frac{1-\epsilon}{4}i} \left(\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \|P_{k, \pm \kappa} Q_{< k + (\epsilon-1)i}^\pm \chi_i(t) \left[\int_{2^{\frac{\epsilon}{C}i}}^t U(t-s)\square\left[\frac{\mathbf{x}}{\mathbf{y}}\right](s, \cdot)ds \right] \|_{PW[\tilde{\kappa}]}^2 \right)^{\frac{1}{2}} \\ & \lesssim \|\chi_{t > 2^{\frac{\epsilon}{C}i}} \square\left[\frac{\mathbf{x}}{\mathbf{y}}\right]\|_{L_t^1 L_x^2} \lesssim 2^{-\frac{\epsilon}{C}\alpha i} \end{aligned}$$

Call the integral in the preceding ψ . From this we get control over the 2nd more complicated norm in the lemma: using Cauchy-Schwartz and the fact that $\|\cdot\|_{PW[\kappa]}$ is essentially unaffected by multiplication with bounded functions, we get

$$\begin{aligned} & 2^{-\frac{l}{2}} \left(\sum_{\tilde{\kappa} \in K_l} \|P_{k, \tilde{\kappa}} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k, \kappa} Q_{< k + (\epsilon-1)i}^\pm \psi \|_{PW[\pm \tilde{\kappa}]}^2 \right)^{\frac{1}{2}} \\ & \lesssim 2^{\frac{1-\epsilon}{4}i} \left(\sum_{\tilde{\kappa} \in K_l} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}, \kappa \subset 2\tilde{\kappa}} \|P_{k, \kappa} Q_{< k + (\epsilon-1)i}^\pm \psi \|_{PW[\pm \kappa]}^2 \right)^{\frac{1}{2}}, \end{aligned} \tag{23}$$

and we just bounded this expression. The estimate for $\ln \mathbf{y}$ is of course analogous. Let $\mu_1(\epsilon) = \frac{\epsilon}{C}\alpha$. Now we need to transfer these statements to ψ_ν . We recall the identity

$$\psi_\nu = \left(\frac{\partial_\nu \mathbf{x}}{\mathbf{y}} + i \frac{\partial_\nu \mathbf{y}}{\mathbf{y}} \right) e^{i \sum_{j=1,2} \Delta^{-1} \partial_j \left(\frac{\partial_j \mathbf{x}}{\mathbf{y}} \right)}$$

We observe as usual that $\frac{\partial_\nu \mathbf{x}}{\mathbf{y}} = \partial_\nu \left(\frac{\mathbf{x}}{\mathbf{y}} \right) - \frac{\mathbf{x}}{\mathbf{y}} \frac{\partial_\nu \mathbf{y}}{\mathbf{y}}$. For example consider the term $\frac{\mathbf{x}}{\mathbf{y}} \frac{\partial_\nu \mathbf{y}}{\mathbf{y}} e^{i \sum_{j=1,2} \Delta^{-1} \partial_j \left(\frac{\partial_j \mathbf{x}}{\mathbf{y}} \right)}$, the other terms being treated similarly. One expands the exponential in a Taylor series, which results in schematic terms of the form

$$a_k \psi \nabla^{-1} \psi_1 \nabla^{-1} \psi_2 \dots \nabla^{-1} \psi_k,$$

where $\nabla^{-1} \psi$ stands for expressions like $\frac{\mathbf{x}}{\mathbf{y}}$, $\nabla^{-1} \left(\frac{\mathbf{x}}{\mathbf{y}} \frac{\partial_i \mathbf{y}}{\mathbf{y}} \right)$, and ψ stands for either $\frac{\partial_\nu \mathbf{y}}{\mathbf{y}}$ or $\frac{\partial_\nu \mathbf{x}}{\mathbf{y}}$. Also note that the coefficients of these expressions decay faster than

exponentially. Now apply a localizer P_{k_0} , $|k_0| < \frac{\epsilon i}{1000C}$ in front. We redefine C so large that $\frac{\epsilon}{C} \ll \mu_1(\epsilon)$, the latter coming from the preceding computation. We claim that if one of the input frequencies has absolute value greater than $i\delta$ for suitable $\delta > \frac{\epsilon}{500C}$, one obtains an exponential gain in i for the norms in the statement of the lemma. This is done inductively: write the expression under consideration as

$$P_k[P_{k_1}(\psi_1 \nabla^{-1} \psi_2 \dots \nabla^{-1} \psi_{k-1}) \nabla^{-1} P_{k_2} \psi_k]$$

If $k_1 > \frac{\epsilon i}{1000C} + 5$, so is k_2 . If $P_{k_1}(\dots)$, $P_{k_2} \psi_k$ are both of first type, one estimates using 3.4(a)

$$\begin{aligned} & \|P_k[P_{k_1}(\psi_1 \nabla^{-1} \psi_2 \dots \nabla^{-1} \psi_{k-1}) \nabla^{-1} P_{k_2} \psi_k]\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \\ & \lesssim \|P_{k_1}(\psi_1 \nabla^{-1} \psi_2 \dots \nabla^{-1} \psi_{k-1})\|_{A[k_1]} \|P_{k_2} \psi_k\|_{A[k_2]} \lesssim \frac{\tilde{c}_{k_1}}{\epsilon} \frac{\tilde{c}_{k_2}}{\epsilon} \end{aligned}$$

If one of them is of 2nd type, one places this into $L_t^2 L_x^{2+}$ to control the portion of the output at modulation $< 2^{k+100}$ (the other portion being controlled by theorem 3.10.). One obtains the same bound, and our decay assumptions on the frequency envelope yield the claim, provided one shows that the 2nd more complicated norm in the statement of the lemma is controlled by $\|\cdot\|_{\dot{X}_k^{0, \frac{1}{2}, 1}}$. This follows

from (11) as well as the preceding computation (23). In case $k_2 < -\frac{\epsilon i}{1000C} - 5$, first assume $P_{k_2} \psi_2$ to be of 2nd type. If $k_2 \geq j - 10$, we estimate

$$\|P_k Q_j[P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \lesssim 2^{\frac{j}{2}} \|P_{k_1} \psi_1\|_{L_t^\infty L_x^2} \|\nabla^{-1} P_{k_2} \psi_2\|_{L_t^2 L_x^\infty} \lesssim 2^{\frac{j-k_2}{2}} \tilde{c}_{k_2}$$

If $k_2 < j - 10$, we estimate

$$\begin{aligned} & \|P_k Q_j[P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \lesssim \|P_k Q_j[P_{k_1} \psi_1 \nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \\ & \quad + \|P_k Q_j[P_{k_1} Q_{\geq j-10} \psi_1 \nabla^{-1} P_{k_2} Q_{< j-10} \psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \end{aligned}$$

Then we have

$$\begin{aligned} & \|P_k Q_j[P_{k_1} \psi_1 \nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \lesssim \|P_{k_1} \psi_1\|_{L_t^\infty L_x^2} \|\nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2\|_{L_t^2 L_x^\infty} \\ & \lesssim 2^{\frac{1}{2+}(k_2-j)} \tilde{c}_{k_2} \end{aligned}$$

Also, we have

$$\begin{aligned} & \|P_k Q_j[P_{k_1} Q_{\geq j-10} \psi_1 \nabla^{-1} P_{k_2} Q_{< j-10} \psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \\ & \lesssim 2^{\frac{j}{2}} \|P_{k_1} Q_{\geq j-10} \psi_1\|_{L_t^2 L_x^2} \|\nabla^{-1} P_{k_2} Q_{< j-10} \psi_2\|_{L_t^\infty L_x^\infty} \lesssim \tilde{c}_{k_2} \end{aligned}$$

Control over the more complicated norm in the lemma is a bit more difficult, and follows from a computation performed in the appendix, in addition to the calculation (23) performed above: one again obtains the bound $\lesssim \tilde{c}_{k_2}$, and the claim follows by summing over $k_2 < -\frac{\epsilon i}{1000C} - 5$. If, on the other hand, $P_{k_2} \psi_2$ is of first type, then one estimates

$$\|P_k[P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \lesssim \|P_{k_1} \psi_1\|_{A[k_1] + \dot{X}_{k_1}^{0, \frac{1}{2}, 1}} \|P_{k_2} \psi_2\|_{A[k_2]} \lesssim \tilde{c}_{k_2}$$

with a similar estimate (using (23)) controlling the 2nd more complicated norm in the lemma. Summing over $k_2 < -\frac{\epsilon i}{1000C} - 5$ results in the desired exponential

gain in i . Now assume $k_1 < -\frac{\epsilon i}{1000C} - 5$. In that case, the claim follows from the estimate

$$\|P_k[P_{k_1}\psi_1\nabla^{-1}P_{k_2}\psi_2]\|_{\dot{X}_k^{0,\frac{1}{2},1}} \lesssim 2^{\delta_1(k_1-k)}|k_1-k_2|\|P_{k_1}\psi_1\|_{S[k_1]}\|P_{k_2}\psi_2\|_{A[k_2]+\dot{X}_{k_2}^{0,\frac{1}{2},1}},$$

by a calculation similar to the immediately preceding. Summing over $k_1 < -\frac{\epsilon i}{1000C} - 5$ results in the desired exponential gain. Having thus shown that we may assume $|k_1| < \frac{\epsilon i}{1000C} + 5$, we apply the same procedure to the expression $P_{k_1}[\psi_1\nabla^{-1}\psi_2\dots\nabla^{-1}\psi_{k-1}]$. Of course the frequency bounds will grow the further one continues this process, but this is counteracted by the rapidly decreasing coefficients coming from the Taylor expansion²⁴. Arguing inductively, we see that it suffices to show that provided two functions $\psi_1 \in \mathcal{S}(\mathbf{R}^{2+1})$, $\psi_2 \in \mathcal{S}(\mathbf{R}^{2+1})$ satisfy the assertions in the lemma, then we get the same assertions for the expression

$$P_k\chi_i(t)[P_{k_1}\psi_1\nabla^{-1}P_{k_2}\psi_2]$$

where $|k|, |k_1|, |k_2| < \delta i$, $\delta \ll \epsilon$. First, we estimate

$$\begin{aligned} & \|P_k[P_{k_1}Q_{>k_1-(1-\epsilon)i}[\chi_i(t)\psi_1]\nabla^{-1}P_{k_2}\psi_2]\|_{\dot{X}_k^{0,\frac{1}{2},1}} \\ & \lesssim (|k| + |k_1| + |k_2|)\|P_{k_1}Q_{>k_1-(1-\epsilon)i}[\chi_i(t)\psi_1]\|_{\dot{X}_{k_1}^{0,\frac{1}{2},1}}\|\nabla^{-1}P_{k_2}\psi_2\|_{\dot{X}_{k_2}^{0,\frac{1}{2},1}+A[k_2]} \\ & \lesssim |i|2^{-\mu(\epsilon)i}. \end{aligned}$$

We have used the bound on $P_{k_2}\psi_2$ from theorem 3.10, as well as the easily verified fact that

$$\|P_{k_1}Q_{>k_1-(1-\epsilon)i}[\chi_i(t)\psi_1]\|_{\dot{X}_{k_1}^{0,\frac{1}{2},1}} \lesssim 2^{-\mu(\epsilon)i}$$

One proceeds similarly when $P_{k_2}\psi_2$ is replaced by $P_{k_2}Q_{>k_2-(1-\epsilon)i}\psi_2$. Thus it suffices to consider the expression

$$P_k\chi_i(t)[P_{k_1}Q_{<k_1-(1-\epsilon)i}\phi\nabla^{-1}P_{k_2}Q_{<k_2-(1-\epsilon)i}\psi]$$

We may always reduce the modulation of the output to size $> 2^{k-\frac{5\epsilon i}{C}}$, as follows easily from 3.4(a). Now we carry out the following decomposition

$$\begin{aligned} & P_k\chi_i(t)[P_{k_1}Q_{<k_1-(1-\epsilon)i}^\pm\phi\nabla^{-1}P_{k_2}Q_{<k_2-(1-\epsilon)i}\psi] \\ & = P_k\chi_i(t)\left[\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa}^c P_{k_1,\kappa} Q_{<k_1-(1-\epsilon)i}^\pm \phi\nabla^{-1}P_{k_2}Q_{<k_2-(1-\epsilon)i}\psi\right] \\ & \quad + P_k\chi_i(t)\left[\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa} P_{k_1,\kappa} Q_{<k_1-(1-\epsilon)i}^\pm \phi\nabla^{-1}P_{k_2}Q_{<k_2-(1-\epsilon)i}\psi\right] \\ & \quad + P_k\chi_i(t)[\phi_i(t,x)P_{k_1}Q_{<k_1-(1-\epsilon)i}^\pm\phi\nabla^{-1}P_{k_2}Q_{<k_2-(1-\epsilon)i}\psi] \end{aligned} \tag{24}$$

In the immediately preceding we let $\chi_{i,\pm\kappa}(t,x)$ localize to two slabs aligned or anti-aligned with κ of length $\sim 2^i$ and thickness $\sim 2^{\epsilon i}$ (the complement of $\chi_{i,\pm\kappa}^c$ within the $\frac{2^{\epsilon i}}{C}$ -neighborhood of the physical light cone), and we let $\phi_i(t,x)$ smoothly localize to the intersection of the complement of the $\frac{2^{\epsilon i}}{C}$ neighborhood of the light

²⁴More precisely, one needs to go to expressions of length up to $\frac{\epsilon i}{10000C}$.

cone and a box of dimensions $\sim 2^i \times 2^i \times 2^i$ (finite propagation speed and properties of the Fourier multipliers). We estimate

$$\begin{aligned} & \|P_k \chi_i(t) [\phi_i(t, x) P_{k_1} Q_{< k_1 - (1-\epsilon)i} \phi \nabla^{-1} P_{k_2} Q_{< k_2 - (1-\epsilon)i} \psi]\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \\ & \lesssim 2^{\frac{\max\{k, k_1, k_2\}}{2}} \|\phi_i(t, x) \chi_i(t) \nabla^{-1} P_{k_2} Q_{< k_2} \psi\|_{L_t^2 L_x^\infty} \|P_{k_1} Q_{< k_1 - (1-\epsilon)i} \phi\|_{L_t^\infty L_x^2} \\ & \quad + \|P_k \chi_i(t) [\phi_i(t, x) P_{k_1} Q_{< k_1 - (1-\epsilon)i} \phi \nabla^{-1} P_{k_2} Q_{[k_2 + (\epsilon-1)i, k_2]} \psi]\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \end{aligned}$$

Now the first summand on the right hand side is immediately controlled²⁵ from Proposition 2.2, while we estimate the 2nd as follows (modulo error terms of order of magnitude 2^{-Ni}):

$$\begin{aligned} & \|P_k \chi_i(t) [\phi_i(t, x) P_{k_1} Q_{< k_1 - (1-\epsilon)i} \phi \nabla^{-1} P_{k_2} Q_{[k_2 + (\epsilon-1)i, k_2]} \psi]\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \\ & \lesssim 2^{\frac{\max\{k, k_1, k_2\}}{2}} \|P_k \chi_i(t) [\phi_i(t, x) P_{k_1} Q_{< k_1 - (1-\epsilon)i} \phi \nabla^{-1} P_{k_2} Q_{[k_2 + (\epsilon-1)i, k_2]} \psi]\|_{L_t^2 L_x^2} \\ & \lesssim 2^{\frac{\max\{k, k_1, k_2\}}{2}} \|\phi_i(t, x)\|_{L_t^\infty L_x^M} \|P_{k_1} Q_{< k_1 - (1-\epsilon)i} \phi\|_{L_t^M L_x^{2+}} \\ & \quad \|\nabla^{-1} P_{k_2} Q_{[k_2 + (\epsilon-1)i, k_2]} \chi_i(t) \psi\|_{L_t^{2+} L_x^\infty} \\ & \lesssim 2^{\frac{\max\{k, k_1, k_2\}}{2}} C(M) 2^{(\frac{1}{M} - \mu)i} \end{aligned}$$

This is acceptable if we choose M large enough. Now we proceed to the other two terms in (24). Consider the first. We have for $\min\{k_1, k_2, k\} + O(1) > j > k - \frac{5\epsilon i}{C}$

$$\begin{aligned} & \|P_k Q_j \chi_i(t) [\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^{\pm} \phi \nabla^{-1} P_{k_2} Q_{< k_2 - (1-\epsilon)i}^{\pm} \psi]\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \\ & \lesssim 2^{\frac{\min\{k, k_1, k_2\} - j}{4}} \left(\sum_{\tilde{\kappa} \in K_{j - \frac{\min\{k, k_1, k_2\}}{2}}} \|P_{k_1, \tilde{\kappa}}\| \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \|\chi_{i, \mp \kappa}^c P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^{\pm} \phi\|_{PW[\pm \tilde{\kappa}]}^2 \right)^{\frac{1}{2}} \\ & \quad \left(\sum_{\tilde{\kappa} \in K_{j - \frac{\min\{k, k_1, k_2\}}{2}}} \|\nabla^{-1} P_{k_2, \tilde{\kappa}} Q_{< k_2 - (1-\epsilon)i}^{\pm} \psi\|_{NFA^*[\pm \tilde{\kappa}]}^2 \right)^{\frac{1}{2}}, \end{aligned}$$

which is bounded by $2^{-\mu i}$. One can easily sum over $j \in [k - ai, k + bi]$, obtaining the desired estimate. If $j > \min\{k_1, k_2, k\} + 100$, say, then necessarily $k_1 = k_2 + O(1) = j + O(1)$, since the inputs $\chi_{i, \mp \kappa}^c P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^{\pm} \psi_2$ etc have modulation $< C 2^{-\epsilon i} < \max\{2^k, 2^{k_{1,2}}\}$. Then one can concurrently microlocalize the inputs to caps $\kappa_{1,2} \in K_{-100}$ of distance ~ 1 , and argue just as before. The 2nd term in (24) has to be decomposed further as follows:

$$\begin{aligned} & P_k \chi_i(t) [\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^{\pm} \phi \nabla^{-1} P_{k_2} Q_{< k_2 - (1-\epsilon)i} \psi] \\ & = P_k \chi_i(t) [\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^{\pm} \phi \nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_2, \kappa} Q_{< k_2 - (1-\epsilon)i}^{\pm} \psi] \\ & \quad + P_k \chi_i(t) [\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^{\pm} \phi \nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_2, \kappa} Q_{< k_2 - (1-\epsilon)i}^{\pm} \psi] \end{aligned}$$

²⁵One may have to shrink the size of $|k|, |k_{1,2}|$ if necessary.

For the first of the immediately preceding terms, we can estimate²⁶ provided $j < \min\{k, k_1, k_2\} + O(1)$

$$\begin{aligned}
& 2^{\frac{j}{2}} \|P_k Q_j \chi_i(t) [\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^\pm \phi \\
& \quad \nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_2, \kappa} Q_{< k_2 - (1-\epsilon)i}^\pm \psi]] \|_{L_t^2 L_x^2} \\
& \lesssim (\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \|P_k Q_j [P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^\pm \phi \nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_2, \kappa} Q_{< k_2 - (1-\epsilon)i}^\pm \psi] \|_{L_t^2 L_x^2}^2)^{\frac{1}{2}}
\end{aligned}$$

Using elementary geometry and the definition of $PW[\kappa]$, $NFA^*[\kappa]$ etc., this in turn is estimated by

$$\begin{aligned}
& \lesssim 2^{-\frac{j - \min\{k, k_1, k_2\}}{4}} (\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \|P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^\pm \phi\|_{NFA^*[\pm \kappa]}^2)^{\frac{1}{2}} \\
& (\sum_{\tilde{\kappa} \in K_{\frac{j - \min\{k, k_1, k_2\}}{2}}} \|P_{k, \tilde{\kappa}} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_2, \kappa} Q_{< k_2 - (1-\epsilon)i}^\pm \psi\|_{PW[\pm \tilde{\kappa}]}^2)^{\frac{1}{2}}
\end{aligned}$$

This can again be estimated by $\lesssim 2^{-\mu i}$, as desired. The case $j > \min\{k, k_1, k_2\} + O(1)$ is treated as before. Finally, we have

$$\begin{aligned}
& \|P_k Q_j \chi_i(t) [\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_1, \kappa} Q_{< k_1 - (1-\epsilon)i}^\pm \phi \\
& \quad \nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_2, \kappa} Q_{< k_2 - (1-\epsilon)i}^\pm \psi]] \|_{L_t^2 L_x^2} \\
& \lesssim 2^{-Ni}
\end{aligned}$$

on account of the support properties of the (Fourier)multipliers. We deduce from this that the expression $P_k \chi_i(t) [P_{k_1} \phi \nabla^{-1} P_{k_2} \psi]$ also satisfies the 2nd property of the lemma: simply apply (11) in addition to (23). \blacksquare

We can now conclude case (I.b.2). The preceding proof in addition to 3.4(g) imply that

$$\begin{aligned}
& \|P_{k_0} Q_{< k_0} [R_\alpha P_{k_1} Q_{< k_0} \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]] \|_{N[k_0]} \\
& \lesssim \|P_{k_1} \delta \psi_1\|_{S[k_1]} \|P_k I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3] \|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \lesssim 2^{-\mu i} \tilde{c}_{k_1}
\end{aligned}$$

One can sum over the admissible frequency ranges here (picking up factors $O(i)$), which yields the desired result.

(I.b.3): *The third term in the expansion:*

$$P_{k_0} Q_{< k_0} [R_\alpha P_{k_1} Q_{< k_0} \square \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]]$$

²⁶Exploit the fact that $\chi_{i, \kappa}$ only smears out the Fourier support by about $2^{-\epsilon i}$, up to negligible error terms.

This is much simpler to estimate, on account of the strong Strichartz type estimates available for ψ_ν : we have

$$\begin{aligned} & \|P_{k_0} Q_{<k_0} [R_\alpha P_{k_1} Q_{<k_0} \square \delta \psi_1 P_k \triangle^{-1} \sum_{j=1,2} \partial_j I[\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \|P_{k_1} Q_{<k_0} \square \delta \psi_1\|_{L_t^2 L_x^M} \|\nabla^{-1} P_{k_2} \chi_i(t) \psi_2\|_{L_t^4 L_x^{4+}} \|P_{k_3} \psi_3\|_{L_t^4 L_x^{4+}} \lesssim 2^{-\mu i} \tilde{c}_{k_1} \end{aligned}$$

(I.b.4) *The fourth term of the expansion:*

$$P_{k_0} Q_{<k_0} [P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 I P_{k_2+O(1)} \square [\nabla^{-1} P_{k_2} Q_{<k_2} \chi_i(t) \psi_2 R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1]]$$

This is straightforward by means of 3.4(a), 3.4(g): one estimates²⁷

$$\begin{aligned} & \|P_{k_0} Q_{<k_0} [P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 P_{k_2+O(1)} I \square [\nabla^{-1} P_{k_2} Q_{<k_2} \chi_i(t) \psi_2 R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1]]\|_{N[k_0]} \\ & \lesssim |k_0 - k_3| \|P_{k_3} \nabla^{-1} \psi_3\|_{A[k_3] + \dot{X}_{k_3}^{0, \frac{1}{2}, 1}} \\ & \quad \|\nabla^{-1} P_{k_2} Q_{<k_2} \chi_i(t) \psi_2 R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1\|_{\dot{X}_{k_2}^{1, \frac{1}{2}, 1}} \\ & \lesssim |k_0 - k_3| \|P_{k_3} \nabla^{-1} \psi_3\|_{\dot{X}_{k_3}^{0, \frac{1}{2}, 1} + A[k_3]} \\ & \quad 2^{k_1} \|P_{k_1} \delta \psi_1\|_{S[k_1]} \|\nabla^{-1} P_{k_2} Q_{<k_2} \chi_i(t) \psi_2\|_{\dot{X}_{k_2}^{0, \frac{1}{2}, 1} + A[k_2]} \\ & \lesssim 2^{-\mu i} \tilde{c}_{k_1} \end{aligned}$$

(I.b.5) *The fifth term of the expansion.*

$$P_{k_0} Q_{<k_0} [P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 P_{k_3} I [\nabla^{-1} \square P_{k_2} Q_{<k_2+O(1)} \chi_i(t) \psi_2 R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1]]$$

This can be estimated by means of lemma 4.4: observe that we may throw an operator $Q_{[k_2+(\epsilon-1)i, k_2]}$ in front of $P_{k_2} Q_{<k_2+O(1)} \chi_i(t) \psi_2$, since in the opposite case we have

$$\begin{aligned} & \|P_{k_0} Q_{<k_0} [P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 P_{k_3} I [\nabla^{-1} \square P_{k_2} Q_{<k_2+(\epsilon-1)i} \chi_i(t) \psi_2 R_\alpha P_{k_1} Q_{<k_0} \delta \psi_1]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim 2^{-k_0} \|P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3\|_{L_t^{2+} L_x^\infty} \|\nabla^{-1} \square P_{k_2} Q_{<k_2+(\epsilon-1)i} \chi_i(t) \psi_2\|_{L_t^2 L_x^2} \|P_{k_1} \delta \psi_1\|_{L_t^M L_x^\infty} \\ & \lesssim 2^{-k_0} 2^{\frac{\epsilon-1}{2}i} 2^{(1-\frac{1}{M})k_1} \tilde{c}_{k_1} \lesssim 2^{-\mu i} \tilde{c}_{k_1} \end{aligned}$$

for very large M , on account of the assumptions on k_1, k_0 . Using lemma 4.4 we have

$$\begin{aligned} & \|P_{k_0} Q_{<k_0} [R_\alpha P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 P_{k_3} I [\nabla^{-1} \square P_{k_2} Q_{[k_2+(\epsilon-1)i, k_2]} \chi_i(t) \psi_2 R_j P_{k_1} \delta \psi_1]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim 2^{-k_0} \|P_{k_3} \nabla^{-1} \psi_3\|_{L_t^{2+} L_x^\infty} \|P_{k_2} Q_{[k_2+(\epsilon-1)i, k_2]} \chi_i(t) \psi_2\|_{\dot{X}_{k_2}^{0, \frac{1}{2}, 1}} \|P_{k_1} \delta \psi_1\|_{L_t^M L_x^\infty} \\ & \lesssim 2^{-\mu i} 2^{-\frac{k_1}{M}} \tilde{c}_{k_1} \end{aligned}$$

Choosing M large enough results in the desired gain in i .

(I.b.6): *The sixth term of the expansion.* This is similar to the third and hence omitted. The terms in the expansion (22) are simple variations of the same kind of reasoning and hence omitted. This concludes case (I.b).

(I.c): *None of (I.a), (I.b) hold but $i \lesssim |k_1|$.* This implies $|k_{2,3}|, |k| \ll i$, and we may treat these as $O(1)$. We also have $i \lesssim k_1$, and $k_1 = k_0 + O(1)$. Now we proceed in close analogy to the immediately preceding case. We may pull the multiplier

²⁷Recall the definition of $S[k]$ via $\|\cdot\|_{A[k]}, \|\cdot\|_{B[k]}$. Also recall that we may replace $P_{k_2} \psi_2$ by $P_{k_2} Q_{<k_2} \psi_2$ etc.

$\chi_i(t)$ past the operator $P_k Q I \nabla^{-1}$ right in front of $P_{k_2} \psi_2$. Also, we may reduce the output and first input $P_{k_1} \delta \psi_1$ to modulation $< O(1)$: for example, one estimates

$$\begin{aligned} & \|P_{k_0} Q_{\geq 0} \nabla_{x,t} [P_{k_1} Q_{< k_1} R_\alpha \delta \psi_1 \nabla^{-1} P_k I [P_{k_2} \chi_i(t) \psi_2 P_{k_3} \psi_3]]\|_{\dot{X}_{k_0}^{-1, -\frac{1}{2}, 1}} \\ & \lesssim \|P_{k_1} \delta \psi_1\|_{L_t^\infty L_x^2} \|\nabla^{-1} P_k I [P_{k_2} \chi_i(t) \psi_2 P_{k_3} \psi_3]\|_{L_t^2 L_x^\infty} \\ & \lesssim 2^{-\mu_i} \tilde{c}_{k_1} \end{aligned}$$

on account of the proof of lemma 4.4, and the fact that by theorem 3.7 we have

$$\|\nabla^{-1} P_k I [P_{k_2} \chi_i(t) \psi_2 P_{k_3} \psi_3]\|_{L_t^2 L_x^\infty} \lesssim \|\nabla^{-1} P_k I [P_{k_2} \chi_i(t) \psi_2 P_{k_3} \psi_3]\|_{\dot{X}_k^{\epsilon, \frac{1}{2^+}, 1}}$$

The estimate for when $P_{k_1} R_\alpha \delta \psi_1$ is replaced by $P_{k_1} Q_{\geq k_1} R_\alpha \psi_1$ is of course similar, since one may place this into $L_t^2 L_x^2$ and gains $2^{-\frac{k_1}{2}}$. Then we resort to the null-form identities (21), (22). We wind up having to estimate the same types of expressions as in (I.b.1-6), with our changed assumptions on the frequencies. We shall do this in a brisk manner here:

(I.c.1): we need to exert some care not to lose in k_0 : use the proof of lemma 4.4 as above to obtain

$$\begin{aligned} & \|P_{k_0} Q_{< O(1)} \square [R_\alpha P_{k_1} Q_{< O(1)} \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]]\|_{\dot{X}_{k_0}^{-1, -\frac{1}{2}, 1}} \\ & \lesssim \|P_{k_1} Q_{< O(1)} \delta \psi_1\|_{L_t^\infty L_x^2} \|P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]\|_{L_t^2 L_x^\infty} \\ & \lesssim \|P_{k_1} Q_{< O(1)} \delta \psi_1\|_{L_t^\infty L_x^2} \|P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]\|_{\dot{X}_k^{\epsilon, \frac{1}{2^+}, 1}} \\ & \lesssim 2^{-\mu_i} \tilde{c}_{k_0} \end{aligned}$$

(I.c.2): This is estimated exactly like (I.b.2).

(I.c.3): The argument is here is slightly more complicated than in (I.b.3); we argue as in (I.c.2):

$$\begin{aligned} & \|P_{k_0} Q_{< O(1)} [R_\alpha P_{k_1} Q_{< O(1)} \square \delta \psi_1 P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim 2^{-k_0} \|P_{k_1} Q_{< O(1)} \square \delta \psi_1\|_{L_t^2 L_x^2} \|P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]\|_{L_t^2 L_x^\infty} \\ & \lesssim \tilde{c}_{k_0} \|P_k \Delta^{-1} \sum_{j=1,2} \partial_j I [\nabla^{-1} P_{k_2} \chi_i(t) \psi_2 R_j P_{k_3} \psi_3]\|_{\dot{X}_k^{\epsilon, \frac{1}{2^+}, 1}} \lesssim 2^{-\mu_i} \tilde{c}_{k_0} \end{aligned}$$

(I.c.4): The argument here is more complicated than in (I.b.4) on account of the fact that we don't get the desired exponential gain from simple application of 3.4(a). Instead, we shall have to resort to lemma 4.4. We split the expression into several

manageable pieces: first, using 3.4(g), we can estimate²⁸

$$\begin{aligned}
& \|P_{k_0} Q_{<O(1)} [P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 \\
& \quad I \square [\nabla^{-1} P_{k_2} Q_{[k_2+(\epsilon-1)i, k_2+O(1)]} (\chi_i(t) \psi_2) R_j P_{k_1} R_\alpha \delta \psi_1]] \|_{N[k_0]} \\
& \lesssim \|P_{k_3} \psi_3 \nabla^{-1} \psi_3\|_{\dot{X}_{k_3}^{0, \frac{1}{2}, 1} + A[k_3]} \\
& \quad \|\nabla^{-1} P_{k_2} Q_{[k_2+(\epsilon-1)i, k_2+O(1)]} (\chi_i(t) \psi_2)\|_{\dot{X}_{k_2}^{0, \frac{1}{2}, 1}} \|P_{k_1} R_\alpha \delta \psi_1\|_{S[k_3]} \lesssim 2^{-\mu i} \tilde{c}_{k_0}
\end{aligned}$$

Next we reduce the large-frequency input $P_{k_1} \delta \psi_1$ to modulation $< 2^{-i\delta}$. For this, throw an operator $Q_{\geq -i\delta}$ in front of $P_{k_1} \delta \psi_1$, let the operator \square fall inside the inner square bracket and use the proof of lemma 4.4, placing the products $P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 \nabla^{-1} P_{k_2} Q_{<k_2+(\epsilon-1)i} (\chi_i(t) \psi_2)$ etc. into $L_t^2 L_x^\infty$ and $P_{k_1} Q_{\geq -i\delta} \delta \psi_1$ into $L_t^2 L_x^2$ and the output into $L_t^1 \dot{H}^{-1}$. Next, we claim that we may apply an operator $Q_{\geq -i\delta+C}$ in front of the inner square bracket. Indeed, if we apply an operator $Q_{<-i\delta+C}$, we estimate using 3.4(a) that

$$\begin{aligned}
& \|P_{k_0} Q_{<O(1)} [P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 \\
& \quad I \square Q_{<-i\delta+C} [\nabla^{-1} P_{k_2} Q_{<k_2+(\epsilon-1)i} (\chi_i(t) \psi_2) R_\alpha P_{k_1} Q_{<-i\delta} \delta \psi_1]] \|_{N[k_0]} \\
& \lesssim \|P_{k_3} \psi_3\|_{\dot{X}_{k_3}^{0, \frac{1}{2}, 1} + A[k]} \\
& \quad \|\square Q_{<-i\delta+C} [\nabla^{-1} P_{k_2} Q_{<k_2+(\epsilon-1)i} (\chi_i(t) \psi_2) R_j P_{k_1} Q_{<-i\delta} \delta \psi_1]\|_{\dot{X}_{k_1}^{0, -\frac{1}{2}, 1}} \\
& \lesssim 2^{\frac{-i\delta-k_2}{4+}} \tilde{c}_{k_1},
\end{aligned}$$

which is acceptable provided $|k_2|$ is small enough in relation to $i\delta$, which we may always arrange. Thus we now reduce to estimating

$$\begin{aligned}
& P_{k_0} Q_{<O(1)} [P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 \\
& \quad I \square Q_{\geq -i\delta+C} [\nabla^{-1} P_{k_2} Q_{<k_2+(\epsilon-1)i} (\chi_i(t) \psi_2) R_\alpha P_{k_1} Q_{<-i\delta} \delta \psi_1]]
\end{aligned}$$

We note that on account of lemma 4.4, we may estimate²⁹

$$\begin{aligned}
& \|P_{k_0} Q_{<O(1)} [P_{k_3} Q_{<k_3} \nabla^{-1} \psi_3 \\
& \quad I \square Q_{\geq -i\delta+C} [\nabla^{-1} \sum_{\kappa \in K \frac{\epsilon-1}{2} i} \chi_{i, \mp \kappa}^c P_{k_2, \kappa} Q_{<k_2+(\epsilon-1)i}^\pm (\chi_i(t) \psi_2) R_\alpha P_{k_1} Q_{<-i\delta} \delta \psi_1]] \|_{N[k_0]} \\
& \lesssim \|P_{k_3} \psi_3\|_{S[k_3]} \|I \square Q_{\geq -i\delta+C} [\nabla^{-1} \sum_{\kappa \in K \frac{\epsilon-1}{2} i} \chi_{i, \mp \kappa}^c P_{k_2, \kappa} Q_{<k_2+(\epsilon-1)i}^\pm (\chi_i(t) \psi_2) \\
& \quad R_\alpha P_{k_1} Q_{<-i\delta} \delta \psi_1]] \|_{\dot{X}_{k_1}^{0, -\frac{1}{2}, 1}},
\end{aligned}$$

²⁸It is easy to see that we may restrict $P_{k_1} R_\alpha \delta \psi_1$ to modulation $< O(1)$, so that we don't lose from R_α . We shall omit this to streamline the formulae.

²⁹We need to assume that $\delta \ll \epsilon$, which we may always arrange.

and one has (see the arguments in [23])

$$\begin{aligned}
& \|P_{k_1+O(1)}Q_{\geq -i\delta+C}[\nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa}^c P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm (\chi_i(t)\psi_2) \\
& \quad R_j P_{k_1} Q_{< -i\delta} \delta\psi_1]\|_{\dot{X}_{k_1}^{0, \frac{1}{2}, 1}}, \\
& \lesssim \left(\sum_{\tilde{\kappa} \in K_{\frac{-i\delta+C-k_2}{2}}} \|\nabla^{-1} P_{k_2,\tilde{\kappa}} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa}^c P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm (\chi_i(t)\psi_2)\|_{PW[\pm\tilde{\kappa}]}^2 \right)^{\frac{1}{2}} \\
& \quad 2^{\frac{-\delta i - k_2}{4}} \left(\sum_{\tilde{\kappa} \in K_{\frac{-i\delta+C-k_2}{2}}} \|P_{k_1,\tilde{\kappa}} Q_{< -i\delta} \delta\psi_1\|_{NFA^*[\pm\tilde{\kappa}]}^2 \right)^{\frac{1}{2}} \\
& \lesssim 2^{-\mu i} \tilde{c}_{k_1}
\end{aligned}$$

Also, arguing as in the proof of lemma 4.4 and using the same notation, we can estimate

$$\begin{aligned}
& \|P_{k_0}Q_{< O(1)}[P_{k_3}Q_{< k_3}\nabla^{-1}\psi_3 \\
& \quad I\Box Q_{\geq -i\delta+C}[\phi_i(t,x)\nabla^{-1}P_{k_2}Q_{< k_2+(\epsilon-1)i}(\chi_i(t)\psi_2)R_\alpha P_{k_1}Q_{< -i\delta}\delta\psi_1]]\|_{L_t^1\dot{H}^{-1}} \\
& \lesssim \|\phi_i(t,x)P_{k_3}Q_{< k_3}\nabla^{-1}\psi_3\|_{L_t^2L_x^\infty} \|\phi_i(t,x)\nabla^{-1}P_{k_2}Q_{< k_2}\chi_i(t)\psi_2\|_{L_t^2L_x^\infty} \|P_{k_1}\delta\psi_1\|_{L_t^\infty L_x^2} \\
& + \|\phi_i(t,x)R_\alpha P_{k_3}Q_{< k_3}\nabla^{-1}\psi_3\|_{L_t^2L_x^\infty} \|\phi_i(t,x)\|_{L_t^M L_x^\infty} \\
& \quad \|\nabla^{-1}P_{k_2}Q_{[k_2+(\epsilon-1)i,k_2]}(\chi_i(t)\psi_2)\|_{L_t^2+L_x^\infty} \|P_{k_1}\delta\psi_1\|_{L_t^\infty L_x^2} \\
& \lesssim 2^{-\mu i} \tilde{c}_{k_1}
\end{aligned}$$

This means we only need to estimate the following expression:

$$\begin{aligned}
& P_{k_0}Q_{< O(1)}[P_{k_3}Q_{< k_3}\nabla^{-1}\psi_3 \\
& \quad I\Box Q_{\geq -i\delta+C}[\nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm (\chi_i(t)\psi_2)R_\alpha P_{k_1}Q_{< -i\delta}\delta\psi_1]]
\end{aligned}$$

Using 3.4(g), it suffices to estimate

$$\|I\Box Q_{\geq -i\delta+C}[\nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm (\chi_i(t)\psi_2)R_\alpha P_{k_1}Q_{< -i\delta}\delta\psi_1]\|_{\dot{X}_{k_1}^{0, -\frac{1}{2}, 1}}$$

We have the identity³⁰

$$\begin{aligned}
& Q_{\geq -i\delta+C}[\nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa} \nabla^{-1} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm (\psi_2)R_\alpha P_{k_1}Q_{< -i\delta}\delta\psi_1] \\
& = \sum_{\tilde{\kappa}_{1,2} \in K_{\frac{-\delta i - k_2}{2}}, \text{dist}(\pm\kappa_1, \pm\kappa_2) \sim 2^{\frac{-i\delta - k_2}{2}}} \\
& Q_{\geq -i\delta+C}[P_{k_2,\tilde{\kappa}_1} \nabla^{-1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa} \nabla^{-1} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm (\psi_2)R_\alpha P_{k_1,\tilde{\kappa}_2} Q_{< -i\delta}\delta\psi_1]
\end{aligned}$$

One can move the multiplier $\chi_{i,\kappa}$ to the outside of this expression while generating error terms rapidly decaying in i and hence negligible for the argument. Hence

³⁰Again this requires that $\delta \ll \epsilon$.

exploiting orthogonality we may now estimate

$$\begin{aligned} & \|Q_{\geq -i\delta+C}[\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\pm\kappa} \nabla^{-1} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm (\psi_2) R_\alpha P_{k_1} Q_{< -i\delta}^\pm \delta\psi_1]\|_{L_t^2 L_x^2} \\ & \lesssim (\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} [\sum_{\tilde{\kappa}_1, 2 \in K_{\frac{-\delta i - k_2}{2}}, \text{dist}(\pm\tilde{\kappa}_1, \pm\tilde{\kappa}_2) \sim 2^{\frac{-i\delta - k_2}{2}}, \kappa \cap \tilde{\kappa}_1 \neq \emptyset} \\ & \|Q_{\geq -i\delta+C}[P_{k_2,\pm\tilde{\kappa}_1} \nabla^{-1} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm \psi_2 R_\alpha P_{k_1,\tilde{\kappa}_2} Q_{< -i\delta}^\pm \delta\psi_1]\|_{L_t^2 L_x^2}]^2)^{\frac{1}{2}} \end{aligned}$$

On the other hand, using (9), we can easily estimate that

$$\begin{aligned} & \|Q_{\geq -i\delta+C}[P_{k_2,\tilde{\kappa}_1} \nabla^{-1} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm \psi_2 R_\alpha P_{k_1,\tilde{\kappa}_2} Q_{< -i\delta}^\pm \delta\psi_1]\|_{L_t^2 L_x^2} \\ & \lesssim 2^{-\frac{-i\delta - k_2}{4}} \|\nabla^{-1} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm \psi_2\|_{PW[\pm\kappa]} \|P_{k_1,\tilde{\kappa}_2} Q_{< -i\delta}^\pm \delta\psi_1\|_{NFA[\pm\tilde{\kappa}_2]} \end{aligned}$$

Plugging this back into the preceding inequality and using the definition of $S[k, \kappa]$, we arrive at the following:

$$\begin{aligned} & \|Q_{\geq -i\delta+C}[\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i,\mp\kappa} \nabla^{-1} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm (\psi_2) R_\alpha P_{k_1} Q_{< -i\delta}^\pm \delta\psi_1]\|_{L_t^2 L_x^2} \\ & \lesssim 2^{\frac{\epsilon-1}{4}i} 2^{-\frac{-i\delta - k_2}{4}} (\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \|\nabla^{-1} P_{k_2,\kappa} Q_{< k_2+(\epsilon-1)i}^\pm \psi_2\|_{S[k_2, \pm\kappa]}^2)^{\frac{1}{2}} \\ & \quad (\sum_{\tilde{\kappa} \in K_{\frac{-\delta i - k_2}{2}}} \|P_{k_1,\tilde{\kappa}} Q_{< -i\delta}^\pm \delta\psi_1\|_{S[k_1, \pm\tilde{\kappa}]}^2)^{\frac{1}{2}} \end{aligned}$$

It follows easily that this expression is $\lesssim 2^{-\mu^i} \tilde{c}_{k_1}$, which is as desired.

(I.c.5): This is much simpler since a derivative has been moved from a high-frequency term to a low-frequency term. We use finite propagation speed and Hoelder's inequality to conclude

$$\begin{aligned} & \|P_{k_0} Q_{< O(1)} [P_{k_3} Q_{< k_3} \nabla^{-1} \psi_3 \nabla^{-1} \square [P_{k_2} Q_{< k_2} \psi_2] R_\alpha P_{k_1} \delta\psi_1]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim 2^{-k_0} 2^{\delta i} \|P_{k_3} Q_{< k_3} \nabla^{-1} \psi_3\|_{L_t^{2+} L_x^\infty} \|\nabla^{-1} \square [P_{k_2} Q_{< k_2} \psi_2]\|_{L_t^2 L_x^\infty} \|R_\alpha P_{k_1} \delta\psi_1\|_{L_t^M L_x^{2+}} \\ & \lesssim 2^{\delta i - k_0} \tilde{c}_{k_1} \end{aligned}$$

One can choose $\delta = \delta(M)$ arbitrarily small so one gets the desired exponential gain in i .

(I.c.6): This is the expression

$$\nabla^{-1} P_{k_2} Q_{< k_2+O(1)} \psi_2 \nabla^{-1} P_{k_3} \psi_3 \square P_{k_1} Q_{< O(1)} \delta\psi_1$$

One estimates this using 3.4(g) as well as lemma 4.4:

$$\begin{aligned} & \|P_{k_0} Q_{< O(1)} [\nabla^{-1} P_{k_2} Q_{< k_2+O(1)} \psi_2 \nabla^{-1} P_{k_3} \psi_3 \square P_{k_1} Q_{< O(1)} \delta\psi_1]\|_{N[k_0]} \\ & \lesssim 2^{k-k_1} \sum_{k < \max\{k_2, k_3\} + O(1)} \|P_k (\nabla^{-1} P_{k_2} Q_{< k_2+O(1)} \psi_2 \nabla^{-1} P_{k_3} \psi_3)\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \\ & \quad \|\square P_{k_1} Q_{< O(1)} \delta\psi_1\|_{\dot{X}_{k_1}^{0, -\frac{1}{2}, 1}} \\ & \lesssim 2^{-\mu^i} \tilde{c}_{k_1} \end{aligned}$$

The terms on the right-hand side of (22) are treated analogously. This concludes case (I.c).

(I.d): *None of (I.a), (I.b), (I.c) hold.* In that case we may assume $|k_{1,2,3}| = O(1)$, $|k| = O(1)$. One proceeds exactly as in the immediately preceding case (I.c), but has to argue differently for case (I.d.5): In that case, we split the expression into two manageable pieces:

$$\begin{aligned}
& \|P_{k_0}[P_{k_3}Q_{<k_3}\nabla^{-1}\psi_3\nabla^{-1}\square[P_{k_2}Q_{[k_2+(\epsilon-1)i,k_2]}\chi_i(t)\psi_2]R_\alpha P_{k_1}Q_{<O(1)}\delta\psi_1]\|_{N[k_0]} \\
& \lesssim \sum_{k < \max\{k_3, k_1\} + O(1)} 2^{k-k_2} \|R_\alpha P_{k_3}Q_{<k_3}\nabla^{-1}\psi_3 R_\alpha P_{k_1}Q_{<k_1}\delta\psi_1\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \\
& \quad \|\square[P_{k_2}Q_{[k_2+(\epsilon-1)i,k_2]}\chi_i(t)\psi_2]\|_{\dot{X}_{k_2}^{0, -\frac{1}{2}, 1}} \\
& \lesssim 2^{-\mu_i} \tilde{C}_{k_1}
\end{aligned}$$

We have used here 3.4(g), 3.4(a), as well as lemma 4.4. Next, once $P_{k_2}\psi_2$ is reduced to very low modulation, we can estimate

$$\begin{aligned}
& \|P_{k_0}[P_{k_3}Q_{<k_3}\nabla^{-1}\psi_3\nabla^{-1}\square[P_{k_2}Q_{<k_2+(\epsilon-1)i}\chi_i(t)\psi_2]R_\alpha P_{k_1}Q_{<O(1)}\delta\psi_1]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \|P_{k_3}\nabla^{-1}\psi_3\|_{L_t^{2+} L_x^\infty} \|\nabla^{-1}\square[P_{k_2}Q_{<k_2+(\epsilon-1)i}\chi_i(t)\psi_2]\|_{L_t^2 L_x^2} \|P_{k_1}\delta\psi_1\|_{L_t^M L_x^{2+}} \\
& \lesssim 2^{-\mu_i} \tilde{C}_{k_1}
\end{aligned}$$

This concludes the treatment of case (I).

(II): *The term (III).* This is treated analogously to case (I). One invokes the identity

$$\begin{aligned}
& \partial^\nu [R_\nu h \sum_{j=1,2} \Delta^{-1} \partial_j [R_i f R_j g - R_j f R_i g]] \\
& = \frac{1}{2} \square \nabla^{-1} h \sum_{j=1,2} \Delta^{-1} \partial_j [R_i f R_j g - R_j f R_i g] \\
& + \frac{1}{2} \square [\nabla^{-1} h \sum_{j=1,2} \Delta^{-1} \partial_j [R_i f R_j g - R_j f R_i g]] \\
& - \frac{1}{2} [\nabla^{-1} h \sum_{j=1,2} \Delta^{-1} \partial_j \square [R_i f R_j g - R_j f R_i g]]
\end{aligned}$$

(III): *The sum of terms (II+VI).* We frequency localize the expression and its inputs as in the preceding. If both $P_{k_1}\psi_1$, $P_{k_3}\psi_3$ are of the first type, we may assume that $|k_j| < i$, $|k| < i$, due to lemma 3.16 as well as the trilinear estimates 3.4(c). Now assume at least one of them is of the 2nd kind. Then the estimate is straightforward on account of the strong Lebesgue type estimates available: for example, assume that $P_{k_1}\psi_1$ is of the 2nd type. We need to estimate the schematically written expression³¹

$$\nabla_{x,t} P_{k_0} [P_{k_1} R_\nu \psi_1 \nabla^{-1} P_k I [P_{k_2} \delta\psi_2 P_{k_3} \psi_3]]$$

In this case one has the estimate

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} [P_{k_1} R_\nu \psi_1 \nabla^{-1} P_k I [P_{k_2} \delta\psi_2 P_{k_3} \psi_3]]\|_{N[k_0]} \\
& \lesssim 2^{-\delta_1 |k_1 - k_0|} 2^{\delta_2 [\min\{k, k_{2,3}\} - \max\{k, k_{2,3}\}]} \tilde{C}_{k_0}
\end{aligned}$$

This is a simple exercise with the exception of the case when $\nu = 0$ and $P_{k_1} R_0 \psi_1$ has large modulation, i. e. we replace this expression by $P_{k_1} Q_{>k_1+100} R_0 \psi_1$. Assume

³¹It is to be kept in mind that the inner square bracket stands for a null-form of type $Q_{\nu j}$.

that $k_1 \gg k$, the other situations being similar or simpler. Then we have

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} [P_{k_1} Q_{>k_1+100} R_0 \psi_1 \nabla^{-1} P_k I [P_{k_2} \delta \psi_2 P_{k_3} \psi_3]]\|_{N[k_0]} \\
& \lesssim 2^{-\frac{k_1}{2}} \|P_{k_0} [P_{k_1} Q_{>k_1+100} R_0 \psi_1 \nabla^{-1} P_k I [P_{k_2} \delta \psi_2 P_{k_3} \psi_3]]\|_{L_t^2 L_x^2} \\
& \lesssim 2^{-\frac{k_1}{2}} \|P_{k_0} [P_{k_1} Q_{>k_1+100} R_0 \psi_1]\|_{L_t^\infty L_x^2} \|\nabla^{-1} P_k I [P_{k_2} \delta \psi_2 P_{k_3} \psi_3]\|_{L_t^2 L_x^\infty} \\
& \lesssim 2^{\frac{k-k_1}{2}} \epsilon 2^{\frac{\min\{k,k_2,k_3\}-\max\{k,k_2,k_3\}}{2}} \frac{\tilde{c}_{k_2}}{\epsilon} \frac{\tilde{c}_{k_3}}{\epsilon}.
\end{aligned}$$

Thus in that case, too, we may assume that $\max_{j=0,\dots,3} \{|k|, |k_j|\} < i$. Next, we claim that we may also reduce the modulations of all inputs to size < 1 . Indeed, for example assume that the first input $P_{k_1} \psi_1$ has modulation > 1 . Then we may estimate (using schematic notation)

$$\begin{aligned}
& \|P_{k_0} Q_{<k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} Q_{>0} R_\nu \psi_1 \nabla^{-1} P_k I [P_{k_2} \delta \psi_2, P_{k_3} \psi_3]]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \|P_{k_1} Q_{[0, \max\{k_2, k_3\}]} \psi_1\|_{\dot{X}_{k_1}^{-\frac{1}{2}, 1, 2}} \|P_{k_2} \delta \psi_2\|_{L_t^M L_x^{2+}} \|P_{k_3} \chi_i(t) \psi_3\|_{L_t^{2+} L_x^\infty} \lesssim 2^{-\mu i} \tilde{c}_{k_0}
\end{aligned}$$

If, on the other hand, we replace $P_{k_0} Q_{<k_0}$ by $P_{k_0} Q_{\geq k_0}$, we have to be careful when $P_{k_1} \psi_1$ is of the 2nd kind. Then we first reduce both $P_{k_2} \delta \psi_2$, $P_{k_3} \psi_3$ to modulation $< 2^{\delta i + O(1)}$ (which is straightforward), and estimate

$$\begin{aligned}
& \|P_{k_0} Q_{\geq k_0} \nabla_{x,t} \chi_i(t) [P_{k_1} Q_{>0} R_\nu \psi_1 \nabla^{-1} P_k I [P_{k_2} Q_{<\delta i} \delta \psi_2, P_{k_3} Q_{<\delta i} \psi_3]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\
& \lesssim 2^{-k} \|P_{k_1} Q_{>0} R_\nu \psi_1\|_{L_t^\infty L_x^2} \|\nabla_{x,t} \nabla^{-1} P_{k_2} Q_{<\delta i} \delta \psi_2\|_{L_t^M L_x^{2+}} \\
& \quad \|\chi_i(t) P_{k_3} Q_{<\delta i} \nabla_{x,t} \nabla^{-1} \psi_3\|_{L_t^{2+} L_x^\infty}
\end{aligned}$$

Assuming (as we may) that all the absolute frequencies $|k|$, $|k_i|$ as well as δi are much less than $(\frac{1}{2} - \frac{1}{2+})i$ and summing over these frequency ranges, one obtains from this and the definition of frequency envelope the upper bound $\lesssim \tilde{c}_{k_0} 2^{-\mu i}$, which is as desired. The remaining cases are handled similarly. Now one reverts to the expansions (21), (22). The only term requiring different treatment than in case (I) is the 2nd term of the expansion. We record this with the appropriate microlocalizations as follows:

$$P_{k_0} \chi_i(t) [P_{k_1} Q_{<0} \psi_1 \sum_{j=1,2} \Delta^{-1} \partial_j \square I P_k [\nabla^{-1} P_{k_2} \delta \psi_2 P_{k_3} \psi_3]]$$

As before, we may innocuously move the localizer $\chi_i(t)$ right in front of the inner square bracket $[\cdot]$. Then we decompose this term as a sum of manageable expressions, keeping in mind lemma 4.4, as well as the estimates in 3.4(a)-3.4(g): first, we have

$$\begin{aligned}
& \|P_{k_0} \chi_i(t) [P_{k_1} Q_{[k_1+(\epsilon-1)i, k_1]} \psi_1 \sum_{j=1,2} \Delta^{-1} \partial_j \square I P_k [\nabla^{-1} P_{k_2} \delta \psi_2 P_{k_3} \psi_3]]\|_{N[k_0]} \\
& \lesssim \|P_{k_1} Q_{[k_1+(\epsilon-1)i, k_1]} \chi_i(t) \psi_1\|_{\dot{X}_{k_1}^{0, \frac{1}{2}, 1}} \|I [\nabla^{-1} P_{k_2} \delta \psi_2 P_{k_3} \psi_3]\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \lesssim 2^{-\mu i} \tilde{c}_{k_0}
\end{aligned}$$

Similarly, one estimates

$$\begin{aligned}
& \|P_{k_0} \chi_i(t) [P_{k_1} Q_{<(\epsilon-1)i+k_1} \psi_1 \\
& \quad \sum_{j=1,2} \Delta^{-1} \partial_j \square I P_k [\nabla^{-1} P_{k_2} \delta \psi_2 P_{k_3} Q_{[k_3+(\epsilon-1)i, k_3]} \psi_3]]\|_{N[k_0]} \\
& \lesssim 2^{-\mu i} \tilde{c}_{k_0}
\end{aligned}$$

It is further easy to see that one may reduce the output to modulation $< 2^{-i\delta}$ for some very small³² $\delta > 0$ (in other words, throw an operator $Q_{<-i\delta}$ in front of it), and also reduce the inner square bracket $\square IP_k[\dots]$ to modulation $> 2^{-i\delta+C}$. Further, one easily reduces the input $P_{k_2}\delta\psi_2$ to modulation $< 2^{-i\delta}$. One can then rewrite the expression as follows:

$$\begin{aligned}
& P_{k_0} Q_{<-i\delta} [\chi_i(t) [P_{k_1} Q_{<(\epsilon-1)i+k_1} \psi_1 \\
& \quad \sum_{j=1,2} \Delta^{-1} \partial_j \square IP_k Q_{>-i\delta+C} [\nabla^{-1} P_{k_2} Q_{<-i\delta} \delta\psi_2 P_{k_3} Q_{<k_3+(\epsilon-1)i} \psi_3]]] \\
& = \sum_{\pm, \pm} \sum_{\substack{\kappa_1, 2 \in K_{\frac{-i\delta - \min\{k_1, k\}}{2}}, \text{dist}(\pm\kappa_1, \pm\kappa_2) \sim 2^{\frac{-i\delta - \min\{k_1, k\}}{2}}}} \\
& P_{k_0, \kappa_1} Q_{<-i\delta}^{\pm} \chi_i(t) [P_{k_1, \kappa_2} Q_{<(\epsilon-1)i+k_1}^{\pm} \psi_1 \\
& \quad \sum_{j=1,2} \Delta^{-1} \partial_j \square IP_k Q_{>-i\delta+C} [\nabla^{-1} P_{k_2} Q_{<-i\delta} \delta\psi_2 P_{k_3} Q_{<k_3+(\epsilon-1)i} \psi_3]]]
\end{aligned} \tag{25}$$

Now we have achieved the kind of situation in which the 2nd part of lemma 4.4 becomes useful: indeed, we can estimate

$$\begin{aligned}
& \|P_{k_0} Q_{<-i\delta} \chi_i(t) [\sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_1, \kappa} Q_{<(\epsilon-1)i+k_1}^{\pm} \psi_1 \\
& \quad \sum_{j=1,2} \Delta^{-1} \partial_j \square IP_k Q_{>-i\delta+C} [\nabla^{-1} P_{k_2} Q_{<-i\delta} \delta\psi_2 P_{k_3} Q_{<k_3+(\epsilon-1)i} \psi_3]] \|_{N[k_0]} \\
& \lesssim (\sum_{\kappa_1 \in K_{\frac{-i\delta - \min\{k_1, k\}}{2}}} \|P_{k_1, \kappa_1} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_1, \kappa} Q_{<(\epsilon-1)i+k_1}^{\pm} \psi_1 \|_{PW[\kappa_1]}^2)^{\frac{1}{2}} \\
& \quad \|P_{k_2} Q_{<-i\delta} \delta\psi_2 \|_{S[k_2]} \|P_{k_3} \psi_3 \|_{A[k_3] + \dot{X}_{k_3}^{0, \frac{1}{2}, 1}}
\end{aligned}$$

This furnishes the desired estimate $\lesssim 2^{-\mu i} \tilde{c}_{k_0}$. One similarly treats the contribution when $P_{k_3} \psi_3$ is replaced by $\sum_{\pm} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_3, \kappa} Q_{<k_3+(\epsilon-1)i}^{\pm} \psi_3$. Thus we now replace both $P_{k_1, 3} \psi_{1, 3}$ by

$$\sum_{\pm} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_1, 3} Q_{<k_1, 3+(\epsilon-1)i}^{\pm} \psi_{1, 3},$$

respectively. Observe that on account of the rapid decay properties of the kernels of the multipliers $Q_{>-i\delta}$ etc., we can rewrite our term (up to errors of order of magnitude 2^{-Ni}) as

$$\begin{aligned}
& \sum_{\pm, \pm} \sum_{\kappa_1, 2 \in K_{\frac{\epsilon-1}{2}i}, \text{dist}(\pm\kappa_1, \pm\kappa_2) \lesssim 2^{\frac{\epsilon-1}{2}i}} P_{k_0, \kappa_2} Q_{<-i\delta}^{\pm} \chi_i(t) [\chi_{i, \mp \kappa_1} P_{k_1, \kappa_1} Q_{<(\epsilon-1)i+k_1}^{\pm} \psi_1 \\
& \quad \sum_{j=1,2} \Delta^{-1} \partial_j \square IP_k Q_{>-i\delta+C} [\nabla^{-1} P_{k_2} Q_{<-i\delta} \delta\psi_2 \chi_{i, \mp \kappa_2} P_{k_3, \kappa_2} Q_{<k_3+(\epsilon-1)i}^{\pm} \psi_3]]
\end{aligned}$$

³²As usual, we choose $\delta \ll \epsilon$.

We can now invoke (12) as well as (25) in order to conclude that the preceding expression is bounded with respect to $\|\cdot\|_{N[k_0]}$ by

$$\begin{aligned} &\lesssim 2^{\delta i} 2^{\frac{\min\{k_i, k\}}{2}} \left(\sum_{\pm} \sum_{\kappa_1 \in K_{\frac{\epsilon-1}{2}i}} \|\chi_{i, \mp \kappa_1} P_{k_1, \kappa_1} Q_{<(\epsilon-1)i+k_1}^{\pm} \psi_1\|_{PW[\pm \kappa_1]}^2 \right)^{\frac{1}{2}} \\ &\quad \left(\sum_{\pm} \sum_{\kappa_3 \in K_{\frac{\epsilon-1}{2}i}} \|\chi_{i, \mp \kappa_3} P_{k_1, \kappa_3} Q_{<(\epsilon-1)i+k_1}^{\pm} \psi_1\|_{PW[\pm \kappa_3]}^2 \right)^{\frac{1}{2}} \|P_{k_2} \delta \psi_2\|_{S[k_2]} \end{aligned}$$

On the other hand, it is easily seen that

$$\left(\sum_{\kappa_1 \in K_{\frac{\epsilon-1}{2}i}} \|\chi_{i, \mp \kappa_1} P_{k_1, \kappa_1} Q_{<(\epsilon-1)i+k_1}^{\pm} \psi_1\|_{PW[\pm \kappa_1]}^2 \right)^{\frac{1}{2}} \lesssim 2^{\frac{\epsilon-1}{2}i} \|P_{k_1} \psi_1\|_{A[k_1] + \dot{X}_{k_1}^{0, \frac{1}{2}, 1}}.$$

This implies that the preceding expression may be bounded by $\lesssim 2^{-\mu i} \tilde{c}_{k_0}$, as desired. This concludes case (III).

(IV): *The term (IV)*. This is similar to (II) and the preceding case (III). Details are omitted. We are now done with establishing Proposition 3.17 for the trilinear null-forms linear in the perturbation $\delta \psi$.

(C): Estimating the trilinear null-forms at least quadratic in $\delta \psi$. The claim here follows directly from 3.4(c) in conjunction with lemma 3.16, provided any function $P_k \psi_\nu$ present is of first type. The other case is treated just like for the expressions linear in $\delta \psi$. This completes the trilinear estimates.

4.5. The quintilinear and higher order terms linear in the perturbation. These turn out to be fairly simple to estimate on account of the favorable Strichartz type estimates available for the radial components ψ_ν . We recall that these terms have the following schematic structure:

$$\mathbf{A}(\delta \psi_\nu, \psi_\nu) = \nabla_{x,t} [\delta \psi \nabla^{-1} [R_\nu \psi \nabla^{-1} [\psi \nabla^{-1} (\psi^2)]]]$$

$$\mathbf{B}(\delta \psi_\nu, \psi_\nu) = \nabla_{x,t} [\psi \nabla^{-1} [R_\nu \delta \psi \nabla^{-1} [\psi \nabla^{-1} (\psi^2)]]]$$

$$\mathbf{C}(\delta \psi_\nu, \psi_\nu) = \nabla_{x,t} [\psi \nabla^{-1} [R_\nu \psi \nabla^{-1} [\delta \psi \nabla^{-1} (\psi^2)]]]$$

$$\mathbf{D}(\delta \psi_\nu, \psi_\nu) = \nabla_{x,t} [\psi \nabla^{-1} [R_\nu \psi \nabla^{-1} [\psi \nabla^{-1} (\delta \psi \psi)]]]$$

$$\mathbf{E}(\delta \psi_\nu, \psi_\nu) = \nabla_{x,t} [\nabla^{-1} [\psi \nabla^{-1} (\psi^2)] \nabla^{-1} I Q_{\nu j} (\delta \psi, \psi)] \text{ etc}$$

$$\mathbf{F}(\delta \psi_\nu, \psi_\nu) = \nabla_{x,t} [\delta \psi \nabla^{-1} [\nabla^{-1} (\psi \nabla^{-1} (\psi^2)) \nabla^{-1} [\psi \nabla^{-1} (\psi^2)]]] \text{ etc}$$

We have left out the other terms like **E**, **F**, in which $\delta \psi$ gets shifted to different positions. We treat here the first quintilinear term in the list, the other terms being tedious reiterations of similar computations.

(A): We microlocalize this term as follows:

$$\nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} R_\nu \psi \nabla^{-1} P_{k_4} [P_{k_5} \psi \nabla^{-1} P_{k_6} (P_{k_7} \psi P_{k_8} \psi)]]]]$$

We first dispose of the case when $\nu = 0$ $R_\nu \psi$ has elliptic microsupport³³, i. e. is replaced by $R_0 P_{k_3} Q_{\geq k_3+100} \psi$. Recall that the expression

$$\nabla^{-1} P_{k_4} [P_{k_5} \psi \nabla^{-1} P_{k_6} (P_{k_7} \psi P_{k_8} \psi)]$$

arises upon substituting the elliptic part for an input of the form $P_{k_4} \psi_j$, where $j \neq 0$. We then rewrite this expression as $P_{k_4} \psi_j - R_j \sum_{i=1,2} R_i \psi_i$, and are led to a term of the schematic form

$$\nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} Q_{\geq k_3+100} R_0 \psi P_{k_4} \psi]]$$

We claim the estimate (under our bootstrap assumption)

$$\begin{aligned} & \nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} Q_{\geq k_3+100} R_0 \psi P_{k_4} \psi]] \\ & \lesssim 2^{-\delta_1 |k_1 - k_0|} 2^{\delta_2 [\min_{i=2,3,4} \{k_i\} - \max_{i=2,3,4} \{k_i\}]} \frac{\tilde{c}_{k_4}}{\epsilon} \tilde{c}_{k_0} \end{aligned}$$

As usual, we may assume $k_0 = 0$. We first treat the case when $k_2 < -10$. Then consider the cases

(a): $k_2 \in [k_3 - 10, k_3 + 10]$ whence $k_4 \leq k_3 + O(1)$. Either $P_{k_4} \psi$ has modulation at least comparable to that of $P_{k_3} Q_{\geq k_3} \psi$, or else either $P_{k_1} \delta \psi$ or the output has modulation at least comparable. These are all similar, so we treat the first possibility. Write this contribution as

$$\sum_{a \geq k_3+100} \nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} Q_a R_0 \psi P_{k_4} Q_{\geq a+O(1)} \psi]]$$

Then we can estimate

$$\begin{aligned} & \left\| \sum_{a \geq k_3+100} \nabla_{x,t} P_{k_0} Q_{<k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} Q_a R_0 \psi P_{k_4} Q_{\geq a+O(1)} \psi]] \right\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \sum_{a \geq k_3+100} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_3} Q_a R_0 \psi\|_{L_t^2 L_x^2} \|P_{k_4} Q_{\geq a+O(1)} \psi\|_{L_t^2 L_x^\infty} \\ & \lesssim \sum_{a \geq k_3+100} 2^{-\frac{k_3}{2}} 2^{(1-\epsilon)(k_4-a)} 2^{\frac{k_4}{2}} \tilde{c}_{k_1} \frac{\tilde{c}_{k_3}}{\epsilon} \frac{\tilde{c}_{k_4}}{\epsilon}, \end{aligned}$$

irrespective of whether $P_{k_3,4} \psi$ is of the first or 2nd type. This is an acceptable bound. Similarly, we have

$$\begin{aligned} & \left\| \sum_{a \geq k_3+100} \nabla_{x,t} P_{k_0} Q_{\geq k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} Q_a R_0 \psi P_{k_4} Q_{\geq a+O(1)} \psi]] \right\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim \sum_{a \geq k_3+100} 2^{-\frac{k_0}{2}} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_3} Q_a R_0 \psi\|_{L_t^\infty L_x^2} \|P_{k_4} Q_{\geq a+O(1)} \psi\|_{L_t^2 L_x^\infty} \\ & \lesssim \sum_{a \geq k_3+100} 2^{\frac{k_4-k_0}{2}} 2^{(1-\epsilon)(k_4-a)} \frac{\tilde{c}_{k_4}}{\epsilon} \tilde{c}_{k_0}, \end{aligned}$$

again an acceptable bound.

(b): $k_2 < k_3 - 10$, whence $k_4 = k_3 + O(1)$. One proceeds as in the preceding case. Again treating the case when $P_{k_4} \psi_4$ is at modulation at least comparable to that

³³When $\nu \neq 0$, this distinction is irrelevant.

of $P_{k_3}\psi_3$, we can estimate

$$\begin{aligned} & \left\| \sum_{a \geq k_3+100} \nabla_{x,t} P_{k_0} Q_{<k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} Q_a R_0 \psi P_{k_4} Q_{\geq a+O(1)} \psi]] \right\|_{L_t^1 \dot{H}^{-1}} \\ & 2^{k_2} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_3} Q_a R_0 \psi\|_{L_t^2 L_x^2} \|P_{k_4} Q_{\geq a+O(1)} \psi\|_{L_t^2 L_x^2} \\ & \lesssim \sum_{a \geq k_3+100} 2^{k_2-k_3} 2^{(1-\epsilon)(k_3-a)} \frac{\tilde{c}_{k_3}}{\epsilon} \frac{\tilde{c}_{k_4}}{\epsilon} \tilde{c}_{k_0}, \end{aligned}$$

which leads to an acceptable estimate. The contribution of $P_{k_0} Q_{\geq k_0}$ is treated similarly.

(c): The case $k_3 \leq k_2 - 10$. Rewrite the term as

$$\begin{aligned} \nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi_1 \nabla^{-1} P_{k_4} \psi P_{k_3} R_0 (1-I) \psi] = \\ \sum_{a > k_3+100} \nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi_1 \nabla^{-1} P_{k_4} \psi P_{k_3} R_0 Q_a \psi] \end{aligned}$$

Then one estimates

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{<k_0} [Q_{\geq a-10} [P_{k_1} \delta \psi_1 \nabla^{-1} P_{k_4} \psi] P_{k_3} R_0 Q_a \psi]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \|Q_{\geq a-10} [P_{k_1} \delta \psi_1 \nabla^{-1} P_{k_4} \psi]\|_{L_t^2 L_x^2} \|P_{k_3} R_0 Q_a \psi\|_{L_t^2 L_x^\infty} \\ & \lesssim 2^{-\frac{a}{2}} 2^{\frac{\min\{a-k_4, 0\}}{4+}} \|P_{k_1} \delta \psi_1\|_{S[k_1]} \|P_{k_4} \psi_4\|_{\dot{X}_{k_4}^{0, \frac{1}{2}, 1} + A[k_4]} \\ & \quad 2^{\frac{k_3}{2}} 2^{\mu(a-k_3)} \|P_{k_3} R_0 Q_a \psi\|_{\dot{X}_{k_4}^{-(\frac{1}{2}-\mu), 1-\mu, 1}} \\ & \lesssim 2^{\min\{\frac{a-k_4}{4+}, 0\}} 2^{\frac{k_3-a}{2+}} \tilde{c}_{k_1} \frac{\tilde{c}_{k_4}}{\epsilon} \end{aligned}$$

One can sum over $a > k_3 + 100$ to obtain the desired estimate. Similarly, we have

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{<k_0} [Q_{<a-10} [P_{k_1} \delta \psi_1 \nabla^{-1} P_{k_4} \psi] P_{k_3} R_0 Q_a \psi]\|_{\dot{X}_{k_0}^{-1, -\frac{1}{2}, 1}} \\ & \lesssim 2^{-\frac{a}{2}} \|Q_{<a-10} [P_{k_1} \delta \psi_1 \nabla^{-1} P_{k_4} \psi]\|_{L_t^\infty L_x^2} \|P_{k_3} R_0 Q_a \psi\|_{L_t^2 L_x^\infty} \\ & \lesssim 2^{\min\{\frac{a-k_4}{4+}, 0\}} 2^{\frac{k_3-a}{2+}} \frac{\tilde{c}_{k_3}}{\epsilon} \tilde{c}_{k_1} \end{aligned}$$

Next, as to the large modulation contribution of the output, we estimate using theorem 3.10

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{\geq k_0} [P_{k_1} \delta \psi_1 \nabla^{-1} P_{k_4} \psi P_{k_3} R_0 Q_a \psi]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim 2^{-\frac{k_0}{2}} \|P_{k_1} \delta \psi_1\|_{L_t^\infty L_x^2} \|\nabla^{-1} P_{k_4} \psi P_{k_3} R_0 Q_a \psi\|_{L_t^2 L_x^\infty} \\ & \lesssim 2^{\delta(k_3-k_4)} 2^{\frac{k_4-k_0}{2}} \tilde{c}_{k_1} \frac{\tilde{c}_{k_4}}{\epsilon} \end{aligned}$$

Now consider the case $k_2 > 10$, whence $k_1 = k_2 + O(1)$. Again this only requires trivial modifications: for example, one estimates in case $k_3 = k_4 + O(1)$

$$\begin{aligned} & \left\| \sum_{a \geq k_3 + 100} \nabla_{x,t} P_{k_0} Q_{<k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} Q_a R_0 \psi P_{k_4} Q_{\geq a+O(1)} \psi]] \right\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \sum_{a \geq k_3 + 100} 2^{-k_2} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_3} Q_a R_0 \psi\|_{L_t^2 L_x^2} \|P_{k_4} Q_{\geq a+O(1)} \psi\|_{L_t^2 L_x^2} \\ & \lesssim \sum_{a \geq k_3 + 100} 2^{-k_3} 2^{(1-\epsilon)(k_3-a)} \tilde{c}_{k_1} \frac{\tilde{c}_{k_4}}{\epsilon} \lesssim 2^{-(1-\epsilon)k_1} 2^{(1-\epsilon)(k_1-k_3)} \tilde{c}_{k_0} \frac{\tilde{c}_{k_4}}{\epsilon}, \end{aligned}$$

and this is again an acceptable estimate. The remaining cases are more of the same. We conclude from this³⁴ that the estimate

$$\|\chi_i(t) P_{k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} (1-I) R_0 \psi P_{k_4} \psi]]\|_{N[k_0]} \lesssim 2^{-\mu_i} \tilde{c}_{k_0}$$

holds, provided we have $i \lesssim \max_{i=2,3,4} \{k_i\}$. Now assume that $\max\{k_{2,3,4}\} \ll i$. We shall treat this quantity as $O(1)$. Also, assume $i \lesssim k_1$. We also omit the operator $\nabla^{-1} P_{k_2}$ for simplicity's sake, and estimate

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{>k_0} [P_{k_1} \delta \psi P_{k_4} \psi R_0 (1-I) \psi]\|_{N[k_0]} \\ & \lesssim 2^{-\frac{k_0}{2}} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_4} \psi R_0 (1-I) \psi\|_{L_t^2 L_x^2} \lesssim 2^{-\frac{k_0}{2}} 2^{\frac{k_4}{2}} \tilde{c}_{k_1} \frac{\tilde{c}_{k_4}}{\epsilon} \end{aligned}$$

This yields the desired exponential gain in i . The contribution of the hyperbolic part is unfortunately a bit more complicated. First, observe that if $a > \delta i$, one estimates

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{\leq k_0} [Q_{<a-10} [P_{k_1} \delta \psi P_{k_4} \psi] R_0 Q_a \psi]\|_{N[k_0]} \\ & \lesssim 2^{-\frac{a}{2}} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_4} \psi\|_{L_t^\infty L_x^\infty} \|R_0 Q_a P_{k_3} \psi\|_{L_t^2 L_x^\infty} \lesssim 2^{-\frac{a}{2+}} \tilde{c}_{k_1} \frac{\tilde{c}_{k_4}}{\epsilon}, \end{aligned}$$

which upon summing over $a > \delta i$ results in the desired exponential gain in i . The case when $Q_{<a-10}$ is replaced by $Q_{\geq a-10}$ is similar (place the output into $L_t^1 \dot{H}^{-1}$). Now consider

$$\nabla_{x,t} P_{k_0} Q_{\leq k_0} \chi_i(t) [P_{k_1} \delta \psi P_{k_4} \psi] R_0 (1-I) P_{k_3} Q_{<\delta i} \psi]$$

We have

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{\leq k_0} \chi_i(t) [Q_{<a-10} [P_{k_1} \delta \psi P_{k_4} \psi] R_0 (1-I) P_{k_3} Q_a \psi]\|_{N[k_0]} \\ & \lesssim 2^{-\frac{a}{2}} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|\chi_i(t) P_{k_4} \psi\|_{L_t^\infty L_x^\infty} \|R_0 (1-I) P_{k_3} Q_a \psi\|_{L_t^2 L_x^2} \\ & \lesssim 2^{-\frac{i}{2+}} \tilde{c}_{k_1} \frac{\tilde{c}_{k_3}}{\epsilon} \end{aligned}$$

Summing over $a < \delta i$ yields an acceptable estimate. Now consider for $k_3 + 100 < a < \delta i$

$$\nabla_{x,t} P_{k_0} Q_{\leq k_0} \chi_i(t) [Q_{\geq a-10} [P_{k_1} \delta \psi P_{k_4} \psi] R_0 (1-I) P_{k_3} Q_a \psi]$$

³⁴Applying the multiplier $\chi_i(t)$ doesn't affect the estimate for reasons explained earlier.

We first reduce $P_{k_1}\delta\psi_1$ to modulation $< 2^{a-100}$, which is straightforward. We estimate

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} Q_{\leq k_0} [Q_{\geq a-10} [P_{k_1} Q_{< a-100} \delta\psi \\
& \quad P_{k_4} Q_{> k+(\epsilon-1)i} (\chi_i(t)\psi)] R_0(1-I) P_{k_3} Q_a \psi]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim 2^{-\frac{a}{2}} \|P_{k_1} Q_{< a-100} \delta\psi P_{k_4} Q_{> k+(\epsilon-1)i} (\chi_i(t)\psi)\|_{\dot{X}_{k_1}^{0, \frac{1}{2}, \infty}} \|R_0(1-I) P_{k_3} Q_a \psi\|_{L_t^2 L_x^2} \\
& \lesssim 2^{-\frac{a}{2}} \|P_{k_1} Q_{< a-100} \delta\psi\|_{S[k_1]} \|P_{k_4} Q_{> k+(\epsilon-1)i} (\chi_i(t)\psi)\|_{\dot{X}_{k_4}^{0, \frac{1}{2}, 1}} \|R_0(1-I) P_{k_3} Q_a \psi\|_{L_t^2 L_x^2} \\
& \lesssim 2^{-\frac{a}{2+}} 2^{-\mu^i} \tilde{c}_{k_1},
\end{aligned}$$

using lemma 4.4. Next, borrowing notation from the proof of lemma 4.4, estimate

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} Q_{\leq k_0} [Q_{\geq a-10} [P_{k_1} Q_{< a-100} \delta\psi \phi_i(t, x) \\
& \quad P_{k_4} Q_{< k+(\epsilon-1)i} (\chi_i(t)\psi)] R_0(1-I) P_{k_3} Q_a \psi]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \|P_{k_1} Q_{< a-100} \delta\psi\|_{L_t^\infty L_x^2} \|\phi_i(t, x) P_{k_4} \psi\|_{L_t^2 L_x^\infty} \|R_0(1-I) P_{k_3} Q_a \psi\|_{L_t^2 L_x^2} \\
& \quad + \|Q_{\geq a-10} [P_{k_1} Q_{< a-100} \delta\psi \phi_i(t, x) Q_{> k+(\epsilon-1)i} \psi]\|_{L_t^2 L_x^2} \|R_0(1-I) P_{k_3} Q_a \psi\|_{L_t^2 L_x^2}
\end{aligned}$$

The 2nd summand may be estimated as in the immediately preceding since we may assume that $|a|, |k|$ etc are $<< \epsilon i$. As to the first summand, we estimate it using

$$\|\phi_i(t) P_{k_4} \psi\|_{L_t^2 L_x^\infty} \lesssim 2^{-\mu^i},$$

from which the desired estimate follows. Now one decomposes

$$\begin{aligned}
(1 - \phi_i(t, x)) P_{k_4} Q_{< k+(\epsilon-1)i} (\chi_i(t)\psi) &= \sum_{\pm} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_4, \kappa} Q_{< k+(\epsilon-1)i}^\pm (\chi_i(t)\psi) \\
&+ \sum_{\pm} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa} P_{k_4, \kappa} Q_{< k+(\epsilon-1)i}^\pm (\chi_i(t)\psi), \tag{26}
\end{aligned}$$

and proceeds as in the trilinear estimates: plugging in the first summand on the right, we have

$$\begin{aligned}
& \|Q_{\geq a-10} [P_{k_1} Q_{< a-100} \delta\psi \sum_{\pm} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_4, \kappa} Q_{< k+(\epsilon-1)i}^\pm (\chi_i(t)\psi)]\|_{\dot{X}_{k_1}^{0, \frac{1}{2}, \infty}} \\
& \lesssim \left(\sum_{\tilde{\kappa}_1 \in K_{\frac{a-k_4}{2}}} \|P_{k_1, \tilde{\kappa}_1} Q_{< a-100}^\pm \delta\psi\|_{NFA^*[\pm \kappa_1]}^2 \right)^{\frac{1}{2}} \\
& \quad \left(\sum_{\tilde{\kappa}_2 \in K_{\frac{a-k_4}{2}}} \|P_{k_4, \tilde{\kappa}_2} \sum_{\kappa \in K_{\frac{\epsilon-1}{2}i}} \chi_{i, \mp \kappa}^c P_{k_4, \kappa} Q_{< k+(\epsilon-1)i}^\pm (\chi_i(t)\psi)\|_{PW[\pm \tilde{\kappa}_4]}^2 \right)^{\frac{1}{2}}
\end{aligned}$$

This can be estimated by $\lesssim 2^{-\mu^i} \tilde{c}_{k_1}$, as desired. Plugging in the 2nd part of (26) is handled as in the trilinear estimates, exploiting the orthogonality of these pieces. This gives the desired estimate. The case when $k_1 << i$ is more elementary. One may always assume that $R_0\psi_3$ lives at modulation $< \delta i$ (argue as before), whence one may estimate

$$\begin{aligned}
& \|\nabla_{x,t} P_{k_0} Q_{\leq k_0} \chi_i(t) [P_{k_1} \delta\psi P_{k_4} \psi] R_0(1-I) P_{k_3} Q_{< \delta i} \psi]\|_{L_t^1 \dot{H}^{-1}} \\
& \lesssim \|P_{k_1} \delta\psi\|_{L_t^M L_x^{2+}} \|P_{k_4} \chi_i(t)\psi\|_{L_t^2 L_x^\infty} \|R_0(1-I) P_{k_3} Q_{< \delta i} \psi\|_{L_t^2 L_x^M}
\end{aligned}$$

Choosing $\frac{1}{2} - \frac{1}{2+} \gg \delta$ results in the desired estimate $\lesssim 2^{-\mu_i} \tilde{c}_{k_1}$. Also, we have

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{>k_0} \chi_i(t) [[P_{k_1} \delta \psi P_{k_4} \psi] R_0 (1-I) P_{k_3} Q_{<\delta i} \psi]\|_{X_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|\chi_i(t) P_{k_4} \psi\|_{L_t^\infty L_x^\infty} \|R_0 (1-I) P_{k_3} Q_{<\delta i} \psi\|_{L_t^2 L_x^2} \lesssim 2^{-\frac{i}{2+}} \tilde{c}_{k_1} \end{aligned}$$

which is as desired. This finally concludes estimating the contribution when $P_{k_3} R_\nu \psi$ is replaced by $P_{k_3} (1-I) R_\nu \psi$. Thus we may replace $P_{k_3} R_\nu \psi$ by $P_{k_3} I R_\nu \psi$, whence we may discard $I R_\nu$ for all intents and purposes. We revert to the original formulation of this term given by $\mathbf{A}(\delta \psi, \psi)$, and decompose the innermost bracket (ψ^2) into a $Q_{\nu j}$ -type null-form as well as error terms at least quadrilinear by means of our standard Hodge-type decompositions, i. e. we write (using schematic notation)

$$\begin{aligned} (\psi^2) &= R_\nu \psi^1 R_j \psi^2 - R_j \psi^1 R_\nu \psi^2 + \nabla^{-1}(\psi \nabla^{-1}(\psi^2)) R_\nu \psi \\ &\quad + \nabla^{-1}(\psi \nabla^{-1}(\psi^2)) \nabla^{-1}(\psi \nabla^{-1}(\psi^2)). \end{aligned}$$

Substituting the $Q_{\nu j}$ -form for now, we claim that we have the estimate

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} \psi \nabla^{-1} P_{k_4} [P_{k_5} \psi \nabla^{-1} P_{k_6} Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]]\|_{N[k_0]} \\ & \lesssim \tilde{c}_{k_1} \tilde{c}_{k_3} \tilde{c}_{k_5} (\tilde{c}_{k_7} + \tilde{c}_{k_8}) \\ & \quad 2^{\delta_1 [\min\{k_7, k_8\} - \max\{k_7, k_8\}]} 2^{\delta_2 [\min_{i=2, \dots, 6} \{k_i\} - \max_{i=2, \dots, 6} \{k_i\}]} \end{aligned}$$

To prove this, one needs to analyze the possible frequency interactions. By scaling invariance, we may assume $k_0 = 0$. The first step consists in reducing $P_{k_6} Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)$ to hyperbolic microsupport, dealing with the contribution of $(1-I) P_{k_6} Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)$. The argument for this is given in the appendix³⁵ of [23]. Then we deal with the following cases in schematic fashion:
(A.a): $k_2 < -10$, $k_5 = k_6 + O(1)$, $k_3 = k_4 + O(1)$. We combine 3.4(b) as well as lemma 4.3 to conclude that

$$\|P_{k_6} Q_{\nu j} I(P_{k_7} \psi P_{k_8} \psi)\|_{L_t^2 L_x^2} \lesssim 2^{\frac{\min\{k_6, 7, 8\}}{2}} (\tilde{c}_{k_7} + \tilde{c}_{k_8}).$$

Next, we have

$$\begin{aligned} & \|P_{k_4} [P_{k_5} \psi \nabla^{-1} P_{k_6} A]\|_{L_t^{\frac{4}{3}} L_x^2} \lesssim 2^{-k_6} \left(\sum_{c \in C_{k_5, k_4 - k_5}} \|P_c \psi\|_{L_t^4 L_x^\infty}^2 \right)^{\frac{1}{2}} \|P_{k_6} A\|_{L_t^2 L_x^2} \\ & \lesssim 2^{\frac{3k_5}{4}} 2^{\frac{k_4 - k_5}{2}} 2^{-k_6} \|P_{k_5} \psi\|_{S[k_5]} \|P_{k_6} A\|_{L_t^2 L_x^2} \end{aligned}$$

Finally, we have

$$\begin{aligned} & \|P_{k_2} [P_{k_3} \psi \nabla^{-1} P_{k_4} A]\|_{L_t^1 L_x^2} \lesssim 2^{-k_4} \left(\sum_{c \in C_{k_3, k_2 - k_3}} \|P_c \psi\|_{L_t^4 L_x^\infty}^2 \right)^{\frac{1}{2}} \|P_{k_4} A\|_{L_t^{\frac{4}{3}} L_x^2} \\ & \lesssim 2^{\frac{3k_3}{4}} 2^{\frac{k_2 - k_3}{2}} 2^{-k_4} \|P_{k_3} \psi\|_{S[k_3]} \|P_{k_4} A\|_{L_t^{\frac{4}{3}} L_x^2} \end{aligned}$$

³⁵If one of the inputs of $Q_{\nu j}(\cdot)$ is of 2nd type, one use lemma 4.3 instead of 3.4(b).

We have used here that $(\sum_{c \in C_{k_4, k_2-k_4}} \|P_c \psi_4\|_{L_t^{\frac{4}{3}} L_x^2}^2)^{\frac{1}{2}} \lesssim \|P_{k_4} \psi_4\|_{L_t^{\frac{4}{3}} L_x^2}$. Combining these inequalities, we easily see that

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{<k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} \psi \nabla^{-1} P_{k_4} [P_{k_5} \psi \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_2} [P_{k_3} \psi \nabla^{-1} P_{k_4} [P_{k_5} \psi \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]\|_{L_t^1 L_x^2} \\ & \lesssim 2^{\frac{3k_3}{4} + \frac{k_2-k_3}{2} - k_4 + \frac{3k_5}{4} + \frac{k_4-k_5}{2} + \frac{\min\{k_6, 7, 8\}}{2} - k_6} \\ & \quad \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_3} \psi\|_{S[k_3]} \|P_{k_5} \psi\|_{S[k_5]} (\tilde{c}_{k_7} + \tilde{c}_{k_8}). \end{aligned}$$

Moreover, we can estimate

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{\geq k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [P_{k_3} \psi \nabla^{-1} P_{k_4} [P_{k_5} \psi \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim 2^{-\frac{k_0}{2}} 2^{k_2} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_2} [P_{k_3} \psi \nabla^{-1} P_{k_4} [P_{k_5} \psi \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]\|_{L_t^2 L_x^1} \\ & \lesssim 2^{k_2 - \frac{k_0}{2}} 2^{\frac{\min\{k_6, 7, 8\}}{2} - k_6} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_3} \psi\|_{L_t^\infty L_x^2} \|P_{k_5} \psi\|_{L_t^\infty L_x^2} (\tilde{c}_{k_7} + \tilde{c}_{k_8}), \end{aligned}$$

One easily verifies that these estimates verify the claim in the case under consideration.

(A.b): $k_2 + O(1) \leq k_6 < k_5$. In this case, we rewrite the term under consideration as

$$\nabla_{x,t} P_{k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [(P_{k_3} \psi \nabla^{-1} P_{k_5} \psi) \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]].$$

Note that necessarily $k_3 = k_5 + O(1)$. We have

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{<k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [(P_{k_3} \psi \nabla^{-1} P_{k_5} \psi) \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim 2^{\frac{4k_2}{M} - k_6} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{\leq k_6 + O(1)} (P_{k_3} \psi \nabla^{-1} P_{k_5} \psi)\|_{L_t^2 L_x^{\frac{M}{2}}} \\ & \quad \|P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)\|_{L_t^2 L_x^2} \end{aligned}$$

Then we estimate

$$\begin{aligned} & \|P_{\leq k_6 + O(1)} (P_{k_3} \psi \nabla^{-1} P_{k_5} \psi)\|_{L_t^2 L_x^{\frac{M}{2}}} \lesssim \sum_{a < k_6 + O(1)} \|P_a (P_{k_3} \psi \nabla^{-1} P_{k_5} \psi)\|_{L_t^2 L_x^{\frac{M}{2}}} \\ & \lesssim \sum_{a < k_6 + O(1)} \left(\sum_{c \in C_{k_3, a-k_3}} \|P_c \psi\|_{L_t^4 L_x^M}^2 \right)^{\frac{1}{2}} \left(\sum_{c \in C_{k_5, a-k_5}} \|P_c \psi\|_{L_t^4 L_x^M}^2 \right)^{\frac{1}{2}} \\ & \lesssim \sum_{a < k_6 + O(1)} 2^{\frac{a-k_3}{2+}} 2^{\frac{a-k_5}{2+}} 2^{(\frac{3}{4} - \frac{2}{M})(k_3+k_5)} 2^{-k_3} \|P_{k_3} \psi\|_{S[k_3]} \|P_{k_5} \psi\|_{S[k_5]}, \end{aligned}$$

where $2+ = 2 + (M)$. Putting this together with the above estimate for $I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)$ easily results in the claim for this case as well. Replacing the operator $Q_{<k_0}$ by $Q_{\geq k_0}$, we have

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{\geq k_0} [P_{k_1} \delta \psi \nabla^{-1} P_{k_2} [(P_{k_3} \psi \nabla^{-1} P_{k_5} \psi) \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim 2^{k_2} \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_2} [(P_{k_3} \psi \nabla^{-1} P_{k_5} \psi) \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]\|_{L_t^2 L_x^1} \\ & \lesssim 2^{\frac{4(k_2-k_5)}{M}} 2^{\frac{1}{2+}(k_6-k_5)} 2^{\frac{\min\{k_6, 7, 8\}}{2} - k_6} \\ & \quad \|P_{k_1} \delta \psi\|_{L_t^\infty L_x^2} \|P_{k_3} \psi\|_{L_t^\infty L_x^2} \|P_{k_5} \psi\|_{L_t^\infty L_x^2} (\tilde{c}_{k_7} + \tilde{c}_{k_8}) \end{aligned}$$

This again verifies the claim.

(A.c): $k_6 \leq k_2 + O(1) < k_5$. Consider the case when the output is reduced to

modulation $< 2^{k_0}$, the opposite case being treated similarly. We rewrite the term as above and estimate

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{<k_0} [P_{k_1} \delta\psi \nabla^{-1} P_{k_2} [(P_{k_3} \psi \nabla^{-1} P_{k_5} \psi) \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim 2^{(\frac{2}{M}-1)k_2} \|P_{k_1} \delta\psi\|_{L_t^\infty L_x^2} \|(P_{k_3} \psi \nabla^{-1} P_{k_5} \psi)\|_{L_t^2 L_x^{\frac{M}{2}}} \|\nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)\|_{L_t^2 L_x^\infty} \\ & \lesssim 2^{\frac{\min\{k_7, 8\} - \max\{k_7, k_8\}}{2}} 2^{\frac{k_6 - k_2}{2}} 2^{\frac{k_2 - k_3}{2+}} \tilde{c}_{k_1} \tilde{c}_{k_3} \tilde{c}_{k_5} (\tilde{c}_{k_7} + \tilde{c}_{k_8}) \end{aligned}$$

This again yields the desired claim.

(A.d): $k_5 \leq k_6 + O(1)$. This case is treated similarly and left out. The remaining frequency interactions are also simple variations of this kind of reasoning and left out. This establishes the claim for term (A)($\delta\psi, \psi$). We immediately deduce that if we reduce the expression to dyadic time $\sim 2^i$, i. e. we apply a multiplier $\chi_i(t)$ in front, we may reduce almost all (logarithmic) frequencies to norm $<< i$, i. e. we may assume $|k_2|, \dots, |k_8| << i$. We shall treat these frequencies as $O(1)$. In that case, though, we can argue rather simply: observe that

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{<k_0} \chi_i(t) [P_{k_1} \delta\psi \nabla^{-1} P_{k_2} [(P_{k_3} \psi \nabla^{-1} P_{k_5} \psi) \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]]\|_{L_t^1 \dot{H}^{-1}} \\ & \lesssim \|P_{k_1} \delta\psi\|_{L_t^\infty L_x^2} \|P_{k_3} \psi\|_{L_t^{2+} L_x^\infty} \|P_{k_5} \psi\|_{L_t^M L_x^{2+}} \|\nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)\|_{L_t^2 L_x^2} \\ & \lesssim 2^{-\mu i} \tilde{c}_{k_0}. \end{aligned}$$

We have used 3.4(c) as well as lemma 4.3. Similarly, we can estimate

$$\begin{aligned} & \|\nabla_{x,t} P_{k_0} Q_{\geq k_0} \chi_i(t) [P_{k_1} \delta\psi \nabla^{-1} P_{k_2} [(P_{k_3} \psi \nabla^{-1} P_{k_5} \psi) \\ & \quad \nabla^{-1} P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)]]]\|_{\dot{X}_{k_0}^{-\frac{1}{2}, -1, 2}} \\ & \lesssim 2^{-\frac{k_0}{2}} \|P_{k_1} \delta\psi\|_{L_t^\infty L_x^2} \|\chi_i(t) \psi_3\|_{L_t^\infty L_x^\infty} \|P_{k_5} \psi\|_{L_t^\infty L_x^\infty} \|P_{k_6} I Q_{\nu j} (P_{k_7} \psi P_{k_8} \psi)\|_{L_t^2 L_x^2} \\ & \lesssim 2^{-\mu i} \tilde{c}_{k_0} \end{aligned}$$

We still need to consider the case when (ψ^2) gets replaced by the error terms $R_\nu \psi \nabla^{-1} (\psi \nabla^{-1} (\psi^2))$ etc. But this is straightforward: first, if $\nu = 0$, one reduces $R_0 \psi$ to $I R_0 \psi$ arguing as before; the latter is morally equivalent to ψ . One then winds up with the following schematic expression:

$$\nabla_{x,t} [\delta\psi \nabla^{-1} [\psi \nabla^{-1} [\psi \nabla^{-1} [\psi \nabla^{-1} [\psi^2]]]]]$$

One can place two of the inputs ψ into $L_t^{2+} L_x^\infty$ and two others into $L_t^\infty L_x^2$, $L_t^M L_x^{2+}$, respectively³⁶. Details are tedious reiterations of previously given arguments. This concludes the estimates for term A($\delta\psi, \psi$). The estimates for the remaining terms B($\delta\psi, \psi$) etc. are more of the same. One invokes the weaker form of the improved Strichartz norms available for the spaces $S[k]$ available by 3.4(f), which is good enough when combined with the stronger version available for the ψ_ν .

4.6. The higher order error terms at least quadratic in $\delta\psi$. Completing the proof of Proposition 3.17. We need to estimate expressions like A($\delta\psi, \psi$) as in the preceding section in which at least one additional ψ has been replaced by $\delta\psi$. Of course we no longer need to gain in time, but we need to obtain an estimate like in theorem 3.15 upon frequency localizing the expression and its inputs. This will eventually allow us to sum over all frequency interactions. The degree of difficulty of these terms varies depending on how many inputs $\delta\psi$ are present, due to the

³⁶One avoids losses for high-high interactions by interpolating with the improved $L_t^4 L_x^\infty$ -norms.

weaker nature of the estimates satisfied by these. However, we shall try to treat all of these terms in a uniform manner, as these estimates are not interesting in and of themselves. Proceeding as in [23], one introduces a null-structure into these terms by splitting all inputs into gradient and elliptic parts, just as before. Substituting the elliptic parts results in error terms of higher order, while the gradient parts contribute to the null-structure. Of course, as before one may apply this splitting to both the inputs $\delta\psi$ as well as ψ . Carrying out this process a finite number of times results in the decomposition described in theorem 3.15. This theorem then ensures that all the resulting terms upon frequency localizing will satisfy the desired estimates, provided all inputs $P_{k_i}\psi$ are of the first type. Our concern here is whether the same kind of estimates hold if they are of the 2nd type. Thus consider for example a typical quintilinear term of the form

$$\nabla_{x,t}[P_{k_1}\delta\psi_1P_{k_2}\nabla^{-1}[P_{k_3}R_\nu\delta\psi\nabla^{-1}[P_{k_4}\psi P_{k_5}\psi]]]$$

The first step here consists in reducing $R_\nu\delta\psi$ to the 'hyperbolic version' $R_\nu I\delta\psi$, arguing as in the preceding subsection. Next, one applies the usual Hodge-type decomposition to the innermost square bracket $[\psi^2]$, replacing this by the sum of a $Q_{\nu j}$ -type null-form as well as error terms at least quadrilinear. One treats the resulting quintilinear null-form just as in [23], resulting in the desired estimate, provided both inputs $P_{k_{4,5}}\psi_{4,5}$ are of the first type. Now assume at least one is of the 2nd type. We first consider the expression

$$\nabla_{x,t}[P_{k_1}\delta\psi_1P_{k_2}\nabla^{-1}[P_{k_3}R_\nu I\delta\psi\nabla^{-1}P_{k_4}[P_{k_5}\psi\nabla^{-1}(1-I)Q_{\nu j}[P_{k_6}\psi P_{k_7}\psi]]]]$$

If both $P_{k_{6,7}}\psi$ are of the first type, one argues here as in the appendix of [23], resulting in the desired estimate. If at least one of these inputs is of the 2nd type, one substitutes lemma 4.3 instead of 3.4(b) in that same argument. Thus we can now replace $Q_{\nu j}(P_{k_6}\psi, P_{k_7}\psi)$ by $IQ_{\nu j}(P_{k_6}\psi, P_{k_7}\psi)$. If both $P_{k_{6,7}}\psi$ are of first type, one again argues as in [23] (these quintilinear null-forms are part of the expansion in theorem 3.15). If one of $P_{k_{6,7}}$ is of 2nd type, the estimate becomes quite simple due to the strong estimates available: observe that then

$$\|P_{k_5}\nabla^{-1}Q_{\nu j}I[P_{k_6}\psi, P_{k_7}\psi]\|_{L_t^1L_x^{2+}} \lesssim 2^{(1-\delta_1)\max\{k_i\}}2^{\delta_2\min\{k_i\}}\tilde{c}_{k_6}\tilde{c}_{k_7}$$

Repeating the usual frequency trichotomies, one gets from here that

$$\begin{aligned} & \|\nabla_{x,t}P_{k_0}Q_{<k_0}[P_{k_1}\delta\psi_1P_{k_2}\nabla^{-1}[P_{k_3}R_\nu I\delta\psi \\ & \quad \nabla^{-1}P_{k_4}[P_{k_5}\psi\nabla^{-1}(1-I)Q_{\nu j}[P_{k_6}\psi P_{k_7}\psi]]]\|_{L_t^1\dot{H}^{-1}} \\ & \lesssim 2^{-\delta_1|k_1-k_0|}2^{\delta_2[\min_{i=2,\dots,7}\{k_i\}-\max_{i=2,\dots,7}\{k_i\}]} \tilde{c}_{k_0} \\ & \|\nabla_{x,t}P_{k_0}Q_{\geq k_0}[P_{k_1}\delta\psi_1P_{k_2}\nabla^{-1}[P_{k_3}R_\nu I\delta\psi \\ & \quad \nabla^{-1}P_{k_4}[P_{k_5}\psi\nabla^{-1}(1-I)Q_{\nu j}[P_{k_6}\psi P_{k_7}\psi]]]\|_{\dot{X}_{k_0}^{-\frac{1}{2},-1,2}} \\ & \lesssim 2^{-\delta_1|k_1-k_0|}2^{\delta_2[\min_{i=2,\dots,7}\{k_i\}-\max_{i=2,\dots,7}\{k_i\}]} \tilde{c}_{k_0} \end{aligned}$$

The remaining error terms are treated analogously.

5. APPENDIX: PROOF OF THEOREM 3.10.

We first check the algebra type estimate. Thus let $\psi_{1,2} \in \mathcal{S}(\mathbf{R}^{2+1})$; we need to estimate $\|P_k[P_{k_1}\psi_1\nabla^{-1}P_{k_2}\psi_2]\|_{\mathcal{S}[k]}$. Of course we may assume that $k = 0$. We

decompose

$$P_0[P_{k_1}\psi_1\nabla^{-1}P_{k_2}\psi_2] = P_0Q_{<100}[P_{k_1}\psi_1\nabla^{-1}P_{k_2}\psi_2] + P_0Q_{\geq 100}[P_{k_1}\psi_1\nabla^{-1}P_{k_2}\psi_2]$$

We first consider the large modulation case, i. e. the 2nd summand on the right. Commence with the case $k_1 > 10$. We freeze the modulation of the output to dyadic size $\sim 2^l$, and further decompose into the following cases:

$$\begin{aligned} P_0Q_l[P_{k_1}\psi_1\nabla^{-1}P_{k_2}\psi_2] &= P_0Q_l[P_{k_1}Q_{\geq l-10}\psi_1\nabla^{-1}P_{k_2}\psi_2] \\ &+ P_0Q_l[P_{k_1}Q_{<l-10}\psi_1\nabla^{-1}P_{k_2}Q_{\geq l-10}\psi_2] + P_0Q_l[P_{k_1}Q_{<l-10}\psi_1\nabla^{-1}P_{k_2}Q_{<l-10}\psi_2] \end{aligned}$$

We treat each of the summands on the right: for the first, observe that irrespective of whether $P_{k_1}\psi_1$ is of first or 2nd type,

$$\begin{aligned} &\|P_0Q_l[P_{k_1}Q_{\geq l-10}\psi_1\nabla^{-1}P_{k_2}\psi_2]\|_{\dot{X}_0^{-(\frac{1}{2}-\mu),1-\mu,1}} \\ &\lesssim 2^{(1-\mu)l} \|P_{k_1}Q_{\geq l-10}\psi_1\|_{L_t^2L_x^2} \|\nabla^{-1}P_{k_2}\psi_2\|_{L_t^\infty L_x^2} \end{aligned}$$

When $l \geq k_1$, one estimates this by

$$\lesssim \sum_{a \geq l-10} 2^{(1-\mu)(l-a)} \|P_{k_1}Q_a\psi_1\|_{\dot{X}_{k_1}^{0,\frac{1}{2},\infty}} \|\nabla^{-1}P_{k_2}\psi_2\|_{L_t^\infty L_x^2}$$

If, on the other hand, $l < k_1$, one estimates this by

$$\lesssim 2^{(\frac{1}{2}-\mu)l} \|P_{k_1}\psi_1\|_{\dot{X}_{k_1}^{0,\frac{1}{2},\infty}} \|\nabla^{-1}P_{k_2}\psi_2\|_{L_t^\infty L_x^2}$$

if $P_{k_1}\psi_1$ is of first type, and by

$$\lesssim 2^{(1-\mu)l} \|P_{k_1}\psi_1\|_{L_t^2L_x^{2+}} \|P_{k_2}\nabla^{-1}\psi_2\|_{L_t^\infty L_x^2}$$

if it is of 2nd type. Moreover, one estimates

$$\|P_0[P_{k_1}\psi_1\nabla^{-1}\psi_2]\|_{L_t^\infty L_x^{1+}} \lesssim 2^{-k_1} \|P_{k_1}\psi_1\|_{L_t^\infty L_x^2} \|P_{k_2}\psi_2\|_{L_t^\infty L_x^2}$$

Next, one estimates the output with respect to $\|\cdot\|_{L_t^{1+}L_x^\infty}$ by interpolating between a crude estimate for $L_t^{1+}L_x^\infty$ gotten by placing the inputs into $L_t^{2+}L_x^\infty$ and a refined estimate for $\|\cdot\|_{L_t^2L_x^\infty}$ by using improved $L_t^4L_x^\infty$ -Strichartz norms for the inputs, resulting again in a small exponential gain in k_1 . This yields the desired bound upon summing over $k_1 = k_2 + O(1) > O(1)$, showing that this contribution to the output is of 2nd type. Now consider the third summand in the above trichotomy: we have $k_1 = k_2 + O(1) = l + O(1)$ in this case. First assume both $P_{k_1,2}\psi_{1,2}$ are of first type. We can decompose

$$\begin{aligned} P_0Q_l[P_{k_1}Q_{<l-10}\psi_1\nabla^{-1}P_{k_2}Q_{<l-10}\psi_2] &= \sum_{\pm} P_0Q_l[P_{k_1}Q_{<l-10}^\pm\psi_1\nabla^{-1}P_{k_2}Q_{<l-10}^\pm\psi_2] \\ &= \sum_{\pm} \sum_{\kappa_1,2 \in K_{-k_1}, \text{dist}(\kappa_1,\kappa_2) \sim 1} P_0Q_l[P_{k_1,\kappa_1}Q_{<l-10}^\pm\psi_1\nabla^{-1}P_{k_2,\kappa_2}Q_{<l-10}^\pm\psi_2] \end{aligned}$$

Now one estimates, using the definition of $S[k, \kappa]$:

$$\begin{aligned} &\| \sum_{\pm} \sum_{\kappa_1,2 \in K_{-k_1}, \text{dist}(\kappa_1,\kappa_2) \sim 1} P_0Q_l[P_{k_1,\kappa_1}Q_{<l-10}^\pm\psi_1\nabla^{-1}P_{k_2,\kappa_2}Q_{<l-10}^\pm\psi_2] \|_{L_t^2L_x^2} \\ &\lesssim \sum_{\pm} \left(\sum_{\kappa_1 \in K_{-k_1}} \|P_{k_1,\kappa_1}Q_{<l-10}^\pm\psi_1\|_{PW[\pm\kappa_1]}^2 \right)^{\frac{1}{2}} \left(\sum_{\kappa_2 \in K_{-k_1}} \|P_{k_2,\kappa_2}Q_{<l-10}^\pm\nabla^{-1}\psi_2\|_{NFA[\pm\kappa_1]}^2 \right)^{\frac{1}{2}} \end{aligned}$$

This in turn can be bounded by

$$\lesssim l^2 2^{-l} \|P_{k_1} \psi_1\|_{A[k_1]} \|P_{k_2} \psi_2\|_{A[k_2]}$$

Multiplying by $2^{(1-\mu)l}$ results in an acceptable estimate. Now assume $P_{k_1} \psi_1$ is of 2nd type, say. In that case, we estimate

$$\begin{aligned} & \|P_0 Q_l [P_{k_1} Q_{<l-10} \psi_1 \nabla^{-1} P_{k_2} Q_{<l-10} \psi_2]\|_{L_t^2 L_x^2} \\ & \lesssim \|P_{k_1} Q_{<l-10} \psi_1\|_{L_t^2 L_x^{2+}} \|P_{k_2} \nabla^{-1} \psi_2\|_{L_t^\infty L_x^2} \end{aligned}$$

Multiplication with $2^{(1-\mu)l}$ again yields an acceptable bound. One estimates the output with respect to $\|\cdot\|_{L_t^\infty L_x^1}$ as well as $\|\cdot\|_{L_t^{1+} L_x^\infty}$ as before, showing that this contribution is of 2nd type as well. The 2nd term of the trichotomy is treated like the first. Now consider the case $k_1 \in [-10, 10]$. We decompose

$$\begin{aligned} & P_0 Q_{>100} \partial_t [P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2] \\ & = P_0 Q_{>100} [P_{k_1} \partial_t \psi_1 \nabla^{-1} P_{k_2} \psi_2] + P_0 Q_{>100} [P_{k_1} \psi_1 R_0 P_{k_2} \psi_2] \end{aligned}$$

We claim that if $P_{k_1} \psi_1$ is of first type, so is the output. This is immediate when $P_{k_2} \psi_2$ is of 1st type. Now assume that $P_{k_2} \psi_2$ is of 2nd type. If $P_{k_1} \psi_1$ has modulation $> 2^{10}$, this is again immediate. In the opposite case, one calculates

$$\begin{aligned} & \|P_0 Q_{>100} [P_{k_1} Q_{<10} \partial_t \psi_1 \nabla^{-1} P_{k_2} \psi_2]\|_{L_t^2 L_x^2} \\ & \lesssim \|P_{k_1} Q_{<10} \partial_t \psi_1\|_{L_t^\infty L_x^2} \|\nabla^{-1} P_{k_2} Q_{>10} \psi_2\|_{L_t^2 L_x^\infty} \end{aligned}$$

Using the definition of $B[k_2]$, one checks that the summation over k_2 can be carried out. Next, we use the first bilinear property in theorem 3.10 (which will be proved later independently of this) to calculate

$$\begin{aligned} & \|P_0 Q_{>100} [P_{k_1} \psi_1 R_0 P_{k_2} \psi_2]\|_{L_t^2 L_x^2} \\ & \lesssim \left(\sum_{c \in C_{k_1, k_2-k_1}} \|P_0 Q_{>100} [P_c Q_{<10} \psi_1 R_0 Q_{>90} \psi_2]\|_{L_t^2 L_x^2}^2 \right)^{\frac{1}{2}} \\ & \lesssim 2^{\delta(k_2-k_1)} \|P_{k_1} \psi_1\|_{A[k_1]} \end{aligned}$$

One can sum over $k_2 < 15$, getting the desired bound. Control over $\|\cdot\|_L$ again follows via the Sobolev embedding. If $P_{k_1} \psi_1$ is of 2nd type, so is the output. This is a simple repetition of arguments before. The case $k_1 < -10$ is more of the same, which finishes the large modulation case. Now consider $P_0 Q_{<100} [P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2]$. First assume $k_1 > 10$, and both $P_{k_1,2} \psi_{1,2}$ of first type. In that case, we claim that the output will be of 2nd type. We need to check that it is controlled with respect to both $\|\cdot\|_{L_t^\infty L_x^{1+}}$ as well as $\|\cdot\|_{L_t^{1+} L_x^\infty}$. For the first norm, this is immediate from Bernstein's inequality. For the 2nd norm, one interpolates between a crude estimate for $L_t^{1+} L_x^\infty$ gotten by placing both $P_{k_1,2} \psi_{1,2}$ into $L_t^{2+} L_x^\infty$ and an estimate for $\|\cdot\|_{L_t^2 L_x^\infty}$ gotten by using improved Strichartz type norms for the inputs. The remaining improved Strichartz type norms constituting $\|\cdot\|_L$ are controlled from Bernstein's inequality. Now assume $P_{k_1} \psi_1$ is of 2nd type. Then the output will again be of 2nd type, as is easily verified by placing $P_{k_2} \psi_2$ into $L_t^\infty L_x^2$. This concludes the case $k_1 > 10$. Now assume $k_1 \in [-10, 10]$, in which case $k_2 < 15$. First assume $P_{k_1} \psi_1$ is of first type. We claim that then the output will be of first type, irrespective of the type of $P_{k_2} \psi_2$. Commence with the case when $P_{k_2} \psi_2$ is

of first type. We need to check the various parts constituting $||\cdot||_{A[0]}$. First, we consider $||\cdot||_{\dot{X}_0^{0, \frac{1}{2}, \infty}}$. Freeze the modulation of the output and decompose

$$\begin{aligned} P_0 Q_j [P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2] &= P_0 Q_j [P_{k_1} Q_{\geq j-10} \psi_1 \nabla^{-1} P_{k_2} \psi_2] \\ &+ P_0 Q_j [P_{k_1} Q_{< j-10} \psi_1 \nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2] + P_0 Q_j [P_{k_1} Q_{< j-10} \psi_1 \nabla^{-1} P_{k_2} Q_{< j-10} \psi_2] \end{aligned}$$

We start with the first term on the righthand side: estimate

$$\begin{aligned} ||P_0 Q_j [P_{k_1} Q_{\geq j-10} \psi_1 \nabla^{-1} P_{k_2} \psi_2]||_{\dot{X}_0^{0, \frac{1}{2}, \infty}} \\ \lesssim 2^{\frac{j}{2}} ||P_{k_1} Q_{\geq j-10} \psi_1||_{L_t^2 L_x^2} ||\nabla^{-1} P_{k_2} \psi_2||_{L_t^\infty L_x^\infty}, \end{aligned}$$

and one easily bounds this by

$$||P_{k_1} \psi_1||_{A[k_1]} ||P_{k_2} \psi_2||_{A[k_2]}$$

Now consider the 2nd term in the above trichotomy. We can estimate this as follows:

$$\begin{aligned} ||P_0 Q_j [P_{k_1} Q_{< j-10} \psi_1 \nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2]||_{\dot{X}_0^{0, \frac{1}{2}, \infty}} \\ \lesssim ||P_{k_1} Q_{< j-10} \psi_1||_{L_t^\infty L_x^2} ||\nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2||_{L_t^2 L_x^\infty} \\ \lesssim 2^{\min\{(\frac{1}{2}-\mu)(k_2-j), 0\}} 2^{\min\{\frac{j-k_2}{4}, 0\}} ||P_{k_1} \psi_1||_{A[k_1]} ||P_{k_2} \psi_2||_{A[k_2]} \end{aligned}$$

We have invoked the improved Bernstein's inequality [37], [23]. Note that one may sum here over $k_2 < 15$ for fixed $j < 100$. Now assume that $P_{k_2} \psi_2$ is of 2nd type. Then we don't use the preceding trichotomy, but divide into the cases $k_2 \geq j - 100$ and $k_2 < j - 100$. In the first case, we estimate

$$\begin{aligned} ||P_0 Q_j [P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2]||_{\dot{X}_0^{0, \frac{1}{2}, \infty}} &\lesssim ||P_{k_1} \psi_1||_{L_t^\infty L_x^2} ||\nabla^{-1} P_{k_2} \psi_2||_{L_t^2 L_x^\infty} \\ &\lesssim 2^{\frac{j-k_2}{2}} ||P_{k_1} \psi_1||_{A[k_1]} ||P_{k_2} \psi_2||_{B[k_2]}. \end{aligned}$$

In the case $k_2 < j - 100$, we split

$$\begin{aligned} P_0 Q_j [P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2] &= P_0 Q_j [P_{k_1} Q_{\geq j-10} \psi_1 \nabla^{-1} P_{k_2} \psi_2] \\ &+ P_0 Q_j [P_{k_1} Q_{< j-10} \psi_1 \nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2] \end{aligned}$$

The first summand on the right is estimated just as before:

$$\begin{aligned} ||P_0 Q_j [P_{k_1} Q_{\geq j-10} \psi_1 \nabla^{-1} P_{k_2} \psi_2]||_{\dot{X}_0^{0, \frac{1}{2}, \infty}} \\ \lesssim 2^{\frac{j}{2}} ||P_{k_1} Q_{\geq j-10} \psi_1||_{L_t^2 L_x^2} ||\nabla^{-1} P_{k_2} \psi_2||_{L_t^\infty L_x^\infty} \end{aligned}$$

as well as

$$\begin{aligned} ||\sum_{k_2 < j-100} P_0 Q_j [P_{k_1} Q_{\geq j-10} \psi_1 \nabla^{-1} P_{k_2} \psi_2]||_{\dot{X}_0^{0, \frac{1}{2}, \infty}} \\ \lesssim 2^{\frac{j}{2}} ||P_{k_1} Q_{\geq j-10} \psi_1||_{L_t^2 L_x^2} ||\nabla^{-1} \psi_2||_{L_t^\infty L_x^\infty} \end{aligned}$$

As for the 2nd summand, we have

$$\begin{aligned} ||P_0 Q_j [P_{k_1} Q_{< j-10} \psi_1 \nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2]||_{\dot{X}_0^{0, \frac{1}{2}, \infty}} \\ \lesssim 2^{\frac{j}{2}} ||P_{k_1} Q_{< j-10} \psi_1||_{L_t^\infty L_x^2} 2^{(\frac{1}{2}-\mu)k_2} 2^{-j(1-\mu)} ||\nabla^{-1} P_{k_2} Q_{\geq j-10} \psi_2||_{\dot{X}_{k_2}^{-(\frac{1}{2}-\mu), 1-\mu, 1}} \end{aligned}$$

Thus we can estimate this contribution by

$$\lesssim 2^{(\frac{1}{2}-\mu)(k_2-j)} ||P_{k_1} \psi_1||_{A[k_1]} ||P_{k_2} \psi_2||_{B[k_2]}$$

Observe that we have obtained the gain $2^{\frac{\min\{k_2-j, 0\}}{2}} 2^{(\frac{1}{2}-\mu)(j-k_2)}$, which allows us to sum over $k_2 < 15$. This concludes the estimate for $\|\cdot\|_{\dot{X}_0^{0, \frac{1}{2}, \infty}}$. We still need to control the complicated null-frame part, as well as $\|\cdot\|_L$. Thus fix $l < -10$, and $-10 \geq \lambda \geq l$, and consider an expression

$$|\lambda|^{-1} \left(\sum_{\kappa \in K_l} \sum_{R \in C_{0, \kappa, \lambda}} \|\tilde{P}_R Q_{<2l}^\pm [P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2]\|_{S[0, \pm \kappa]}^2 \right)^{\frac{1}{2}}$$

The following estimates are irrespective of the type of $P_{k_2} \psi_2$. Then we split this into a bunch of contributions: Observe the identity

$$\tilde{P}_R Q_{<2l}^\pm [P_{k_1} \psi_1 \nabla^{-1} P_{<2l} Q_{<2l} \psi_2] = \tilde{P}_R Q_{<2l}^\pm [\tilde{P}_{R_1} Q_{<2l+O(1)}^\pm \psi_1 \nabla^{-1} P_{<2l} Q_{<2l} \psi_2],$$

where $R_1 = (1 + \frac{1}{1000})R$. Next, note that from the definition of $S[k, \kappa]$, we have

$$\begin{aligned} \|\tilde{P}_{R_1} Q_{<2l+O(1)}^\pm \psi_1\|_{S[k, \pm \kappa]} &\lesssim \sum_{\kappa_1 \in K_{l-100}, \kappa_1 \subset (1 + \frac{1}{100})\kappa} \|\tilde{P}_{R_1} P_{k_1, \kappa_1} Q_{<2l+O(1)}^\pm \psi_1\|_{S[k, \pm \kappa]} \\ &\lesssim \sum_{\kappa_1 \in K_{l-100}, \kappa_1 \subset (1 + \frac{1}{100})\kappa} \|\tilde{P}_{R_1} P_{k_1, \kappa_1} Q_{<2l+O(1)}^\pm \psi_1\|_{S[k, \pm \kappa_1]}, \end{aligned}$$

Note that $\frac{11}{10}\kappa_1 \subset \frac{11}{10}\kappa = \tilde{\kappa}$, the latter as in the definition of $PW[\kappa]$. Hence the above inequality. Putting these observations together, we get

$$\begin{aligned} |\lambda|^{-1} \left(\sum_{\kappa \in K_l} \sum_{R \in C_{0, \kappa, \lambda}} \|\tilde{P}_R Q_{<2l}^\pm [P_{k_1} \psi_1 \nabla^{-1} P_{<2l} Q_{<2l} \psi_2]\|_{S[0, \pm \kappa]}^2 \right)^{\frac{1}{2}} \\ \lesssim |\lambda|^{-1} \left(\sum_{\kappa \in K_{l+O(1)}} \sum_{R \in C_{0, \kappa, \lambda+O(1)}} \|\tilde{P}_R Q_{<2l}^\pm \psi_1\|_{S[k, \pm \kappa]}^2 \right)^{\frac{1}{2}} \|\nabla^{-1} \psi_2\|_{L_t^\infty L_x^\infty}, \end{aligned}$$

which leads to the desired estimate. Further, using

$$\begin{aligned} \|P_0 Q_{<2l} [P_{k_1} \psi_1 \nabla^{-1} P_{<2l} Q_{\geq 2l} \psi_2]\|_{\dot{X}_0^{0, \frac{1}{2}, 1}} \\ + \|P_0 Q_{<2l} [P_{k_1} \psi_1 \nabla^{-1} P_{\geq 2l} \psi_2]\|_{\dot{X}_0^{0, \frac{1}{2}, 1}} \lesssim \|P_{k_1} \psi_1\|_{A[k_1]} \end{aligned}$$

We can estimate

$$\begin{aligned} |\lambda|^{-1} \left(\sum_{\kappa \in K_l} \sum_{R \in C_{0, \kappa, \lambda}} \|\tilde{P}_R Q_{<2l}^\pm [P_{k_1} \psi_1 \nabla^{-1} P_{\geq 2l} \psi_2]\|_{S[0, \pm \kappa]}^2 \right)^{\frac{1}{2}} \\ \lesssim \|P_0 Q_{<2l} [P_{k_1} \psi_1 \nabla^{-1} P_{\geq 2l} \psi_2]\|_{\dot{X}_0^{0, \frac{1}{2}, 1}}, \\ |\lambda|^{-1} \left(\sum_{\kappa \in K_l} \sum_{R \in C_{0, \kappa, \lambda}} \|\tilde{P}_R Q_{<2l}^\pm [P_{k_1} \psi_1 \nabla^{-1} P_{<2l} Q_{\geq 2l} \psi_2]\|_{S[0, \pm \kappa]}^2 \right)^{\frac{1}{2}} \\ \lesssim \|P_0 Q_{<2l} [P_{k_1} \psi_1 \nabla^{-1} P_{<2l} Q_{\geq 2l} \psi_2]\|_{\dot{X}_0^{0, \frac{1}{2}, 1}}, \end{aligned}$$

and so the desired estimate follows easily for the contributions of $P_{k_1} \psi_1 \nabla^{-1} P_{\geq 2l} \psi_2$, $P_{k_1} \psi_1 \nabla^{-1} P_{<2l} Q_{\geq 2l} \psi_2$. The estimate for $\|\cdot\|_L$ is quite similar. This concludes the estimates when $P_{k_1} \psi_1$ is of first type. If it is of 2nd type, so will be the output. This is straightforward to check. The case $k_1 < -10$ is a tedious reiteration of similar estimates and hence omitted. This concludes the proof for the assertions of theorem 3.10 as far as the estimates concerning $\|\cdot\|_S$ are concerned. We now

proceed to the assertions concerning the bilinear estimates, as well as the estimates for

$$\|R_0[\psi_1 A(\nabla^{-1}\psi_2)]\|_{L_t^\infty L_x^2}, \|P_k Q_{<k+O(1)}[\psi_1 A(\nabla^{-1}\psi_2)]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}}$$

We start with the latter, which follows from the refined assertion for the function β in the decomposition

$$P_k Q_{<k+O(1)}[\psi_1 A(\nabla^{-1}\psi_2)] = \alpha + \beta$$

To understand this decomposition, one expands $A(\nabla^{-1}\psi_2)$ into a Taylor series (using real analyticity). One winds up with schematic expressions of the form $\psi \nabla^{-1}\psi$, $\psi \nabla^{-1}\psi \nabla^{-1}\psi$ etc. For the first type of expression, $P_k[\psi \nabla^{-1}\psi]$, the previous proof revealed that the only contribution of 2nd type arises from high-high interactions. But for these it is straightforward to verify that they are controlled with respect to $\|\cdot\|_{\dot{X}_k^{0, \frac{1}{2}, 1}}$, see 3.4(a). Now one proceeds inductively, assuming the assertion to be true for both $P_{k_{1,2}}\psi_{1,2} \forall k_{1,2} \in \mathbf{Z}$, and considering $P_k[P_{k_1}\psi_1 \nabla^{-1}P_{k_2}\psi_2]$. For example, considering high-high interactions, when $P_{k_1}\psi_1$ is of 2nd type, one estimates

$$\begin{aligned} \|P_k Q_{<k+O(1)}[P_{k_1}\psi_1 \nabla^{-1}P_{k_2}\psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} &\lesssim 2^{(\frac{3}{2}-\epsilon)k} \|P_{k_1}\psi_1\|_{L_t^2 L_x^{2+}} \|\nabla^{-1}P_{k_2}\psi_2\|_{L_t^\infty L_x^2} \\ &\lesssim 2^{(\frac{3}{2}-\epsilon)(k-k_2)} \|P_{k_1}\psi_1\|_{A[k_1]} \|P_{k_2}\psi_2\|_{S[k_2]} \end{aligned}$$

One can sum here over $k_1 = k_2 + O(1) \geq k + O(1)$, obtaining the desired estimate. In case of high-low interactions, one reasons as follows: assume $P_{k_1}\psi_1$ is of 2nd type. Then for $j < k + O(1)$, $k = k_1 + O(1)$, we have

$$\begin{aligned} P_k Q_j[P_{k_1}\psi_1 \nabla^{-1}\psi_2] &= P_k Q_j[P_{k_1} Q_{\geq j-10}\psi_1 \nabla^{-1}P_{<j-10}Q_{<j-10}\psi_2] \\ &\quad + P_k Q_j[P_{k_1}\psi_1 \nabla^{-1}P_{<j-10}Q_{\geq j-10}\psi_2] + P_k Q_j[P_{k_1}\psi_1 \nabla^{-1}P_{\geq j-10}\psi_2] \end{aligned}$$

We estimate each of the terms on the right: for the first, we have

$$\begin{aligned} &\left\| \sum_{j < k+O(1)} P_k Q_j[P_{k_1} Q_{\geq j-10}\psi_1 \nabla^{-1}P_{<j-10}Q_{<j-10}\psi_2] \right\|_{\dot{X}_k^{0, \frac{1}{2}, 1}} \\ &\lesssim \sum_{j < k+O(1)} 2^{\frac{j}{2}} \sum_{a \geq j-10} \|P_{k_1} Q_a \psi_1\|_{L_t^2 L_x^{2+}} \|\nabla^{-1}\psi_2\|_{L_t^\infty L_x^\infty} \\ &\lesssim \|P_{k_1}\psi_1\|_{\dot{X}_{k_1}^{0, \frac{1}{2}, 1}} \|\nabla^{-1}\psi_2\|_{L_t^\infty L_x^\infty} \end{aligned}$$

Now consider the 2nd term on the right. We exploit the improved range of Strichartz type norms available for $P_{k_1}\psi_1$:

$$\begin{aligned} &\|P_k Q_j[P_{k_1}\psi_1 \nabla^{-1}P_{<j-10}Q_{\geq j-10}\psi_2]\|_{\dot{X}_k^{0, \frac{1}{2}, \infty}} \\ &\lesssim 2^{\frac{j}{2}} \|P_{k_1}\psi_1\|_{L_t^2 L_x^{2+}} \|\nabla^{-1}P_{<j-10}Q_{\geq j-10}\psi_2\|_{L_t^\infty L_x^M} \\ &\lesssim 2^{\frac{j-k_1}{2+}} \|P_{k_1}\psi_1\|_{A[k_1]} \end{aligned}$$

One can sum over $j < k + O(1)$, resulting in the desired bound. Finally, the last term in the above trichotomy is handled similarly. One handles low-high interactions analogously. This shows that the desired property for β is inherited from one stage of the expansion to the next.

Next, consider $P_k R_0[\psi_1 A(\nabla^{-1}\psi_2)]$. One expands $A(\nabla^{-1}\psi_2)$ into a Taylor series,

and proceeds inductively. Assume one has $\sup_{k \in \mathbf{Z}} \|P_k R_0 \psi_{1,2}\|_{L_t^\infty L_x^2} \leq C$, and consider $P_k R_0[\psi_1 \nabla^{-1} \psi_2]$. One rewrites this as

$$\begin{aligned} P_k R_0[\psi_1 \nabla^{-1} \psi_2] &= \sum_{k_1 > k+10, k_1 = k_2 + O(1)} P_k R_0[P_{k_1} \psi_1 \nabla^{-1} P_{k_2} \psi_2] \\ &\quad + P_k R_0[P_{[k-10, k+10]} \psi_1 \nabla^{-1} \psi_2] + P_k R_0[P_{< k-10} \psi_1 \nabla^{-1} P_{[k-10, k+10]} \psi_2] \end{aligned}$$

One treats each of these terms separately: for the first, let the derivative ∂_t fall inside, replacing this by

$$\begin{aligned} &\| \sum_{k_1 > k+10, k_1 = k_2 + O(1)} P_k \nabla^{-1} [P_{k_1} \partial_t \psi_1 \nabla^{-1} P_{k_2} \psi_2 + P_{k_1} \psi_1 R_0 \psi_2] \|_{L_t^\infty L_x^2} \\ &\lesssim \sum_{k_1 = k_2 + O(1)} [\|P_{k_1} R_0 \psi_1\|_{L_t^\infty L_x^2} \|P_{k_2} \psi_2\|_{L_t^\infty L_x^2} + \|P_{k_1} \psi_1\|_{L_t^\infty L_x^2} \|R_0 P_{k_2} \psi_2\|_{L_t^\infty L_x^2}] \leq C \end{aligned}$$

Next one estimates

$$\begin{aligned} &\|P_k R_0[P_{[k-10, k+10]} \psi_1 \nabla^{-1} \psi_2]\|_{L_t^\infty L_x^2} \\ &\leq \|P_k \nabla^{-1} [P_{[k-10, k+10]} \partial_t \psi_1 \nabla^{-1} \psi_2 + P_{[k-10, k+10]} \psi_1 R_0 \psi_2]\|_{L_t^\infty L_x^2} \\ &\lesssim \|P_{[k-10, k+10]} R_0 \psi_1\|_{L_t^\infty L_x^2} \|\nabla^{-1} \psi_2\|_{L_t^\infty L_x^\infty} \\ &\quad + \|P_{[k-10, k+10]} \nabla^{-1} \psi_1\|_{L_t^\infty L_x^2} \|R_0 P_{< k+15} \psi_2\|_{L_t^\infty L_x^\infty} \end{aligned}$$

Again one checks that this can be bounded by C . The third term above is more of the same. The assertion follows from this. We proceed to the bilinear estimates. For this we expand the expression

$P_k[\psi_1 A(\nabla^{-1} \psi_2)]$ as a Taylor series, obtaining terms of the schematic form ψ , $\psi \nabla^{-1} \psi$, $\psi \nabla^{-1} \psi \nabla^{-1} \psi$ etc. Consider a typical such term of the form

$$P_0[\psi_1 \nabla^{-1} \psi_2 \nabla^{-1} \psi_3 \dots \nabla^{-1} \psi_a]$$

This term will have a coefficient decaying like $(a!)^{-1}$. Freeze the modulation of the output to dyadic size $\sim 2^l$. If $\nabla^{-1} \psi_a$ has modulation $\geq 2^{l+10}$, we can estimate

$$\begin{aligned} &\|P_0 Q_l R_0[\psi_1 \nabla^{-1} \psi_2 \nabla^{-1} \psi_3 \dots \nabla^{-1} Q_{> l+10} \psi_a]\|_{L_t^2 L_x^2} \\ &\lesssim 2^l \|Q_{> l+O(1)}[\psi_1 \nabla^{-1} \psi_2 \nabla^{-1} \psi_3 \dots \nabla^{-1} \psi_{a-1}]\|_{L_t^2 L_x^2} \|\nabla^{-1} Q_{> l+10} \psi_a\|_{L_t^2 L_x^2 + L_t^2 L_x^\infty} \\ &\lesssim 2^{-\frac{l}{2+}} \end{aligned}$$

Summing over $l > O(1)$ results in the desired upper bound. Similarly, we have

$$\begin{aligned} &\| \sum_{l > O(1)} P_0 Q_l R_0[\psi_1 \nabla^{-1} \psi_2 \nabla^{-1} \psi_3 \dots \nabla^{-1} Q_{[l-10, l+10]} \psi_a] \|_{L_t^2 L_x^2} \\ &\lesssim \left(\sum_{l > O(1)} \|Q_{[l-10, l+10]} R_0 \psi_a\|_{L_t^2 L_x^2 + L_t^2 L_x^\infty}^2 \right)^{\frac{1}{2}} \lesssim C \end{aligned}$$

which is again acceptable. Thus assume now that $\nabla^{-1} \psi_a$ has modulation $< 2^{l-10}$. If $[\psi_1 \dots \nabla^{-1} \psi_{a-1}]$ has modulation $\geq 2^{l-10}$, repeat the same process with this expression instead of the longer one. Continuing in this fashion, one either

eventually³⁷ forces ψ_1 to be at modulation $> 2^{l-10a}$, or else one arrives at a situation of the following sort:

$$P_0 Q_l [Q_{>l-10(a-m-1)} [Q_{<l-10(a-m)} [\psi_1 \dots \nabla^{-1} \psi_m] \nabla^{-1} \psi_{m+1}] \dots \nabla^{-1} \psi_a]$$

If $\nabla^{-1} \psi_{m+1}$ has modulation $> 2^{l-10(a-m)}$, one argues just as before. One loses exponentially in m , which is counteracted by the small coefficient eventually applied to the expression from the Taylor expansion. Thus we may also apply an operator $Q_{<l-10(a-m)}$ in front of $\nabla^{-1} \psi_{m+1}$. In this case, however, we can write

$$\begin{aligned} & P_0 Q_l [Q_{>l-10(a-m-1)} [Q_{<l-10(a-m)} [\psi_1 \dots \nabla^{-1} \psi_m] \nabla^{-1} Q_{<l-10(a-m)} \psi_{m+1}] \dots \nabla^{-1} \psi_a] \\ &= P_0 Q_l [Q_{>l-10(a-m-1)} [P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} [\psi_1 \dots \nabla^{-1} \psi_m] \\ &\quad \nabla^{-1} P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} \psi_{m+1}] \dots \nabla^{-1} \psi_a] \end{aligned}$$

As to the first bilinear inequality in Theorem 3.10, we decompose

$$P_k \phi = P_k Q_{\geq k+100} \phi + \sum_{\pm} \sum_{\kappa \in K_{-100}} P_{k,\kappa} Q_{<k+100}^{\pm} \phi$$

Then we have

$$\begin{aligned} & \left(\sum_{c \in C_{k,r}} \|P_0 Q_l R_0 [Q_{>l-10(a-m-1)} [P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} [\psi_1 \dots \nabla^{-1} \psi_m] \right. \\ &\quad \left. \nabla^{-1} P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} \psi_{m+1}] \dots \nabla^{-1} \psi_k] P_c Q_{\geq k+100} \phi \|_{L_t^2 L_x^2}^2 \right)^{\frac{1}{2}} \\ &\lesssim 2^{\min\{k+r,0\}} \|P_0 Q_l R_0 [Q_{>l-10(a-m-1)} [P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} [\psi_1 \dots \nabla^{-1} \psi_m] \\ &\quad \nabla^{-1} P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} \psi_{m+1}] \dots \nabla^{-1} \psi_k] \|_{L_t^\infty L_x^2} \|P_k Q_{\geq k+100} \phi \|_{L_t^2 L_x^2} \\ &\lesssim 2^{\min\{\frac{k}{2},0\}} 2^r 2^{10m} \end{aligned}$$

The loss in m will be counteracted by the small Taylor coefficients. Now consider the contribution of

$$\sum_{\pm} \sum_{\kappa \in K_{-100}} P_{k,\kappa} Q_{<k+100}^{\pm} \phi$$

We have the identity

$$\begin{aligned} & P_0 Q_l [Q_{>l-10(a-m-1)} [P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} [\psi_1 \dots \nabla^{-1} \psi_m] \\ &\quad \nabla^{-1} P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} \psi_{m+1}] \dots \nabla^{-1} \psi_k] \\ &= \sum_{\pm} \sum_{\kappa_1, \kappa_2 \in K_{-100}, \text{dist}(\kappa_1, \kappa_2) \sim 1} \\ & P_0 Q_l [Q_{>l-10(a-m-1)} [P_{l-10(a-m)+O(1), \kappa_1} Q_{<l-10(a-m)}^{\pm} [\psi_1 \dots \nabla^{-1} \psi_m] \\ &\quad \nabla^{-1} P_{l-10(a-m)+O(1), \kappa_2} Q_{<l-10(a-m)}^{\pm} \psi_{m+1}] \dots \nabla^{-1} \psi_k] \end{aligned}$$

³⁷We may assume $a \ll l$, since otherwise one gets a large gain in l just from the Taylor coefficient.

We may thus assume that either $\pm\kappa_1$ or $\pm\kappa_2$ has angular separation ~ 1 from $\pm\kappa$. Assume w. l. o. g. that $\text{dist}(\pm\kappa_1, \pm\kappa) \sim 1$. Then estimate

$$\begin{aligned}
& \left(\sum_{c \in C_{k,r}} \|P_0 Q_l R_0 [Q_{>l-10(a-m-1)} [P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} [\psi_1 \dots \nabla^{-1} \psi_m] \right. \\
& \quad \left. \nabla^{-1} P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} \psi_{m+1}] \dots \nabla^{-1} \psi_k] P_c P_{k,\kappa} Q_{<k+100}^\pm \phi \|_{L_t^2 L_x^2}^2 \right)^{\frac{1}{2}} \\
& \lesssim 2^l \| [P_{l-10(a-m)+O(1), \kappa_1} Q_{<l-10(a-m)}^\pm [\psi_1 \dots \nabla^{-1} \psi_m] \|_{NFA^*[\pm\kappa]} \\
& \|\nabla^{-1} P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} \psi_{m+1}\|_{L_t^\infty L_x^2} \left(\sum_{c \in C_{k,r}} \|P_c P_{k,\kappa} Q_{<k+100}^\pm \phi\|_{PW[\pm\kappa]}^2 \right)^{\frac{1}{2}} \\
& \lesssim 2^{\frac{k}{2}} 2^{\frac{r}{2}} 2^l \|\nabla^{-1} P_{l-10(a-m)+O(1)} Q_{<l-10(a-m)} \psi_{m+1}\|_{L_t^\infty L_x^2}
\end{aligned}$$

We have used that

$$\begin{aligned}
& \| [P_{l-10(a-m)+O(1), \kappa_1} Q_{<l-10(a-m)}^\pm [\psi_1 \dots \nabla^{-1} \psi_m] \|_{NFA^*[\pm\kappa]} \\
& \lesssim \| [P_{l-10(a-m)+O(1), \kappa_1} Q_{<l-10(a-m)}^\pm [\psi_1 \dots \nabla^{-1} \psi_m] \|_{\dot{X}_{l-10(a-m)}^{\frac{1}{2}, 1} + A[l-10(a-m)]} \leq C
\end{aligned}$$

One can sum now over l , obtaining the desired estimate with a loss 2^{10m} , which is made up for by the small Taylor coefficient in front. The 2nd bilinear inequality is proved similarly, as is the version when $A(\nabla^{-1} \psi_2)$ is replaced by $A(\nabla^{-1}(\psi_2 \nabla^{-1} \psi_3))$. This completes the proof of theorem 3.10.

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